Numerical Analysis - Part II

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Lecture 6

Partial differential equations of evolution

Solving the diffusion equation

We consider the solution of the diffusion equation

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \qquad 0 \le x \le 1, \quad t \ge 0,$$

with initial conditions $u(x,0) = u_0(x)$ for t = 0 and Dirichlet boundary conditions $u(0,t) = \phi_0(t)$ at x = 0 and $u(1,t) = \phi_1(t)$ at x = 1. By Taylor's expansion

$$\begin{array}{rcl} \frac{\partial u(x,t)}{\partial t} &=& \frac{1}{k} \big[u(x,t+k) - u(x,t) \big] + \mathcal{O}(k), & k = \Delta t \,, \\ \frac{\partial^2 u(x,t)}{\partial x^2} &=& \frac{1}{h^2} \big[u(x-h,t) - 2u(x,t) + u(x+h,t) \big] + \mathcal{O}(h^2), & h = \Delta x \,, \end{array}$$

so that, for the true solution, we obtain

$$u(x,t+k) = u(x,t) + \frac{k}{h^2} \left[u(x-h,t) - 2u(x,t) + u(x+h,t) \right] + \mathcal{O}(k^2 + kh^2).$$
(1)

That motivates the numerical scheme for approximation $u_m^n \approx u(x_m, t_n)$ on the rectangular mesh $(x_m, t_n) = (mh, nk)$:

$$u_m^{n+1} = u_m^n + \mu \left(u_{m-1}^n - 2u_m^n + u_{m+1}^n \right), \qquad m = 1...M.$$
 (2)

Here $h = \frac{1}{M+1}$ and $\mu = \frac{k}{h^2} = \frac{\Delta t}{(\Delta x)^2}$ is the so-called *Courant number*. With μ being fixed, we have $k = \mu h^2$, so that the local truncation error of the scheme is $\mathcal{O}(h^4)$. Substituting whenever necessary initial conditions u_m^0 and boundary conditions u_0^n and u_{M+1}^n , we possess enough information to advance in (2) from $\mathbf{u}^n := [u_1^n, \dots, u_M^n]$ to $\mathbf{u}^{n+1} := [u_1^{n+1}, \dots, u_M^{n+1}]$. Similarly to ODEs or Poisson equation, we say that the method is *convergent* if, for a fixed μ , and for every T > 0, we have

$$\lim_{h\to 0} |u_m^n - u(x_m, t_n)| = 0 \quad \text{uniformly for} \quad (x_m, t_n) \in [0, 1] \times [0, T] \,.$$

In the present case, however, a method has an extra parameter μ , and it is entirely possible for a method to converge for some choice of μ and diverge otherwise.

Definition 1 (Stability in the context of time-stepping methods for PDEs of evolution)

A numerical method for a PDE of evolution is *stable* if (for zero boundary conditions) it produces a uniformly bounded approximation of the solution in any bounded interval of the form $0 \le t \le T$ when $h \to 0$ and the generalized Courant number $\mu = k/h^r$, with r being the maximum degree of the differential operator, is constant.

Theorem 2 (The Lax equivalence theorem)

Suppose that the underlying PDE is well-posed and that it is solved by a numerical method with an error of $\mathcal{O}(h^{p+r})$, $p \ge 1$, where r is the maximum degree of the differential operator. Then stability \Leftrightarrow convergence. Recall the basic norm inequalities:

$$||x||_2 \le ||x||_1 \le \sqrt{n} ||x||_2$$

$$\|x\|_{\infty} \leq \|x\|_2 \leq \sqrt{n} \, \|x\|_{\infty},$$

$$\|x\|_{\infty} \leq \|x\|_{1} \leq n \, \|x\|_{\infty},$$

where $x \in \mathbb{C}^n$.

Proving stability directly

Although we can deduce from the theorem that $\mu \leq \frac{1}{2}$ implies stability, we will prove directly that stability $\Leftrightarrow \mu \leq \frac{1}{2}$. Let $\mathbf{u}^n = [u_1^n, \ldots, u_M^n]^T$. We can express the recurrence (2)

$$u_m^{n+1} = u_m^n + \mu \left(u_{m-1}^n - 2u_m^n + u_{m+1}^n \right), \qquad m = 1...M,$$

in the matrix form

$$\mathbf{u}_{h}^{n+1} = A_{h}\mathbf{u}_{h}^{n}, \qquad A_{h} = I + \mu A_{*}, \qquad A_{*} = \begin{bmatrix} -2 & 1 \\ 1 & \ddots & \ddots \\ & \ddots & \ddots & 1 \\ & & 1 - 2 \end{bmatrix}_{M \times M}$$

Proving stability directly

Here A_* is TST, with $\lambda_{\ell}(A_*) = -4\sin^2 \frac{\pi\ell h}{2}$, hence $\lambda_{\ell}(A_h) = 1 - 4\mu \sin^2 \frac{\pi\ell h}{2}$, so that its spectrum lies within the interval $[\lambda_M, \lambda_1] = [1 - 4\mu \cos^2 \frac{\pi h}{2}, 1 - 4\mu \sin^2 \frac{\pi h}{2}]$. Since A_h is symmetric, we have

$$\|A_h\|_2 = \rho(A_h) = \begin{cases} |1 - 4\mu \sin^2 \frac{\pi h}{2}| \le 1, & \mu \le \frac{1}{2}, \\ |1 - 4\mu \cos^2 \frac{\pi h}{2}| > 1, & \mu > \frac{1}{2} & (h \le h_\mu) \end{cases}$$

We distinguish between two cases.

1) $\mu \leq \frac{1}{2}$: $\|\mathbf{u}^n\| \leq \|A\| \cdot \|\mathbf{u}^{n-1}\| \leq \cdots \leq \|A\|^n \|\mathbf{u}^0\| \leq \|\mathbf{u}^0\|$ as $n \to \infty$, for every \mathbf{u}^0 .

2) $\mu > \frac{1}{2}$: Choose \mathbf{u}^0 as the eigenvector corresponding to the largest (in modulus) eigenvalue, $|\lambda| > 1$. Then $\mathbf{u}^n = \lambda^n \mathbf{u}^0$, becoming unbounded as $n \to \infty$. Suppose that we want to solve the differential equation

$$y' = f(t, y), \qquad y(t_0) = y_0.$$

Euler's method is given by

$$y_{n+1} = y_n + kf(t_n, y_n),$$

where $k = t_{n+1} - t_n$ is the step size.

Let $u_m(t) = u(mh, t)$, m = 1...M, $t \ge 0$. Approximating $\partial^2/\partial x^2$ as before, we deduce from the PDE that the *semidiscretization*

$$\frac{du_m}{dt} = \frac{1}{h^2}(u_{m-1} - 2u_m + u_{m+1}), \qquad m = 1...M$$
(3)

carries an error of $\mathcal{O}(h^2)$. This is an ODE system, and we can solve it by any ODE solver. Thus, Euler's method yields (2), while backward Euler results in

$$u_m^{n+1} - \mu(u_{m-1}^{n+1} - 2u_m^{n+1} + u_{m+1}^{n+1}) = u_m^n.$$

This approach is commonly known as *the method of lines*. Much (although not all!) of the theory of finite-difference methods for PDEs of evolution can be presented as a two-stage task: first semidiscretize, getting rid of space variables, then use an ODE solver.

Typically, each stage is conceptually easier than the process of discretizing in unison in both time and in space (so-called *full discretization*).

Suppose that we want to solve the differential equation

$$y' = f(t, y), \qquad y(t_0) = y_0.$$

The trapezoidal rule is given by the formula

$$y_{n+1} = y_n + \frac{1}{2}k\Big(f(t_n, y_n) + f(t_{n+1}, y_{n+1})\Big),$$

where $k = t_{n+1} - t_n$ is the step size.

Discretizing the ODE (3) with the trapezoidal rule, we obtain

$$u_m^{n+1} - \frac{1}{2}\mu(u_{m-1}^{n+1} - 2u_m^{n+1} + u_{m+1}^{n+1}) = u_m^n + \frac{1}{2}\mu(u_{m-1}^n - 2u_m^n + u_{m+1}^n),$$
(4)

where m = 1...M. Thus, each step requires the solution of an $M \times M$ TST system. The error of the scheme is $\mathcal{O}(k^3 + kh^2)$, so basically the same as with Euler's method. However, as we will see, Crank–Nicolson enjoys superior stability features, as compared with the method (2).

Suppose that a numerical method (with zero boundary conditions) can be written in the form

$$\mathbf{u}_h^{n+1} = A_h \mathbf{u}_h^n,$$

where $\mathbf{u}_h^n \in \mathbb{R}^M$ are vectors, $A_h \in \mathbb{R}^{M \times M}$ is a matrix, and $h = \frac{1}{M+1}$. Then $\mathbf{u}_h^n = (A_h)^n \mathbf{u}_h^0$, and

$$\|\mathbf{u}_{h}^{n}\| = \|(A_{h})^{n}\mathbf{u}_{h}^{0}\| \le \|(A_{h})^{n}\| \cdot \|\mathbf{u}_{h}^{0}\| \le \|A_{h}\|^{n} \cdot \|\mathbf{u}_{h}^{0}\|,$$

for any vector norm $\|\cdot\|$ and the induced matrix norm $\|A\| = \sup \frac{\|A\mathbf{x}\|}{\|\mathbf{x}\|}$. If we define stability as preserving the boundedness of \mathbf{u}_h^n with respect to the norm $\|\cdot\|$, then, from the inequality above,

 $\|A_h\| \leq 1$ as $h \to 0 \implies$ the method is stable.

Usually, the norm of \mathbf{u}_h is set to be an averaged Euclidean length, namely, $\|\mathbf{u}_h\|_h := [h \sum_{i=1}^M |u_i|^2]^{1/2}$, and that does not change the Euclidean matrix norm. The reason for the factor $h^{1/2}$ is to ensure that, because of the convergence of Riemann sums, we obtain

$$\|\mathbf{u}_h\|_h := \left[h\sum_{i=1}^M |u_i|^2\right]^{1/2} \to \left[\int_0^1 |u(x)|^2 \mathrm{d}x\right]^{1/2} =: \|u\|_{L_2} \qquad (h \to 0),$$

provided that u is a square-integrable function such that $\mathbf{u}_h = u(x)|_{\Omega_h}.$

Definition 3 (Normal matrices)

We say that a matrix A is *normal* if $A = QD\bar{Q}^T = QDQ^*$, where D is a (complex) diagonal matrix and Q is a unitary matrix (such that $Q\bar{Q}^T = I$, where the bar in \bar{Q} means complex conjugation). In other words, a matrix is normal if it has a complete set of orthonormal eigenvectors.

Examples of the real normal matrices, besides the familiar symmetric matrices $(A = A^{T})$, include also the matrices which are skew-symmetric $(A = -A^{T})$, and more generally the matrices with skew-symmetric off-diagonal part.

Norms of normal matrices

Proposition 4

If A is normal, then $||A|| = \rho(A)$.

Proof. Let **u** be any vector (complex-valued as well). We can expand it in the basis of the orthonormal eigenvectors $\mathbf{u} = \sum_{i=1}^{n} a_i \mathbf{q}_i$. Then $A\mathbf{u} = \sum_{i=1}^{n} \lambda_i a_i \mathbf{q}_i$, and since \mathbf{q}_i are orthonormal, we obtain

$$\|A\|_{2} := \sup_{\mathbf{u}} \frac{\|A\mathbf{u}\|_{2}}{\|\mathbf{u}\|_{2}} = \sup_{a_{i}} \frac{\{\sum_{i=1}^{n} |\lambda_{i}a_{i}|^{2}\}^{1/2}}{\{\sum_{i=1}^{n} |a_{i}|^{2}\}^{1/2}} = |\lambda_{\max}|.$$

Remark 5

More generally, one can prove that, for any matrix A, we have $||A||_2 = [\rho(A\bar{A}^T)]^{1/2}$, and the previous result for normal matrices can be deduced from that formula.

Crank–Nicolson method for diffusion equation

Let

$$u_m^{n+1} - \frac{1}{2}\mu(u_{m-1}^{n+1} - 2u_m^{n+1} + u_{m+1}^{n+1}) = u_m^n + \frac{1}{2}\mu(u_{m-1}^n - 2u_m^n + u_{m+1}^n),$$

where m = 1...M. Then $B\mathbf{u}^{n+1} = C\mathbf{u}^n$, where the matrices B and C are Toeplitz symmetric tridiagonal (TST),

$$\mathbf{u}^{n+1} = B^{-1}C\mathbf{u}^{n}, \qquad B = I - \frac{1}{2}\mu A_{*}, \qquad A_{*} = \begin{bmatrix} -2 & 1 \\ 1 & \ddots & \ddots \\ \ddots & \ddots & 1 \\ C = I + \frac{1}{2}\mu A_{*}, \qquad 1 - 2 \end{bmatrix}_{M \times M}$$

All $M \times M$ TST matrices share the same eigenvectors, hence so does $B^{-1}C$. Moreover, these eigenvectors are orthogonal. Therefore, also $A = B^{-1}C$ is normal and its eigenvalues are

$$\lambda_k(A) = \frac{\lambda_k(C)}{\lambda_k(B)} = \frac{1 - 2\mu \sin^2 \frac{1}{2}\pi kh}{1 + 2\mu \sin^2 \frac{1}{2}\pi kh} \quad \Rightarrow \quad |\lambda_k(A)| \le 1, \qquad k = 1...M.$$

Consequently Crank–Nicolson is stable for all $\mu > 0$.

Note: Similarly to the situation with stiff ODEs, this *does not* mean that $k = \Delta t$ may be arbitrarily large, but that the only valid consideration in the choice of $k = \Delta t$ vs $h = \Delta x$ is accuracy.

Convergence of the Crank-Nicolson method for diffusion equation

It is not difficult to verify that the local error of the Crank-Nicolson scheme is $\eta_m^n = \mathcal{O}(k^3 + kh^2)$, where $\mathcal{O}(k^3)$ is inherited from the trapezoidal rule (compared to $\mathcal{O}(k^2)$ for the Euler method). We also have

$$\|\boldsymbol{\eta}^n\| = \{h \sum_{m=1}^M |\eta_m^n|^2\}^{1/2} = \mathcal{O}(k^3 + kh^2).$$

Hence, for the error vectors \mathbf{e}^n we have

 $B\mathbf{e}^{n+1} = C\mathbf{e}^n + \eta^n \quad \Rightarrow \quad \|\mathbf{e}^{n+1}\| \le \|B^{-1}C\| \cdot \|\mathbf{e}^n\| + \|B^{-1}\| \cdot \|\eta^n\| \,.$

We have just proved that $||B^{-1}C|| \le 1$, and we also have $||B^{-1}|| \le 1$, because all the eigenvalues of B are greater than 1 (by Gershgorin's theorem). Therefore, $||\mathbf{e}^{n+1}|| \le ||\mathbf{e}^n|| + ||\boldsymbol{\eta}^n||$, and

$$\|\mathbf{e}^n\| \le \|\mathbf{e}^0\| + n\|\eta\| = n\|\eta\| \le \frac{cT}{k}(k^3 + kh^2) = cT(k^2 + h^2).$$

Thus, taking $k = \alpha h$ will result in $\mathcal{O}(h^2)$ error of approximation.

We consider the solution of the advection equation

$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial x}, \qquad 0 \le x \le 1, \quad t \ge 0,$$

with initial conditions $u(x,0) = u_0(x)$ for t = 0 and Dirichlet boundary conditions $u(0,t) = \phi_0(t)$ at x = 0 and $u(1,t) = \phi_1(t)$ at x = 1.

Crank–Nicolson for advection equation

Let

$$u_m^{n+1} - u_m^n = \frac{1}{4}\mu(u_{m+1}^{n+1} - u_{m-1}^{n+1}) + \frac{1}{4}\mu(u_{m+1}^n - u_{m-1}^n), \qquad m = 1...M.$$

(This is the trapezoidal rule applied to the semidiscretization of advection equation $\frac{\partial u}{\partial t} = \frac{\partial u}{\partial x}$). In this case, $\mathbf{u}^{n+1} = B^{-1}C\mathbf{u}^n$, where the matrices B and C are Toeplitz antisymmetric tridiagonal,

$$B = \begin{bmatrix} 1 & -\frac{1}{4}\mu & & \\ \frac{1}{4}\mu & 1 & \ddots & \\ & \ddots & \ddots & -\frac{1}{4}\mu \\ & & \frac{1}{4}\mu & 1 \end{bmatrix}, \qquad C = \begin{bmatrix} 1 & \frac{1}{4}\mu & & \\ -\frac{1}{4}\mu & 1 & \ddots & \\ & \ddots & \ddots & \frac{1}{4}\mu \\ & & -\frac{1}{4}\mu & 1 \end{bmatrix}$$

Crank–Nicolson for advection equation

Similarly to Exercise 4, the eigenvalues and eigenvectors of the matrix

$$S = \begin{bmatrix} \alpha & \beta \\ -\beta & \alpha & \ddots \\ & \ddots & \ddots & \beta \\ & & -\beta & \alpha \end{bmatrix},$$

are given by $\lambda_k = \alpha + 2i\beta \cos kx$, and $\mathbf{w}_k = (i^m \sin kmx)_{m=1}^M$, where $x = \pi h = \frac{\pi}{M+1}$. So, all such *S* are normal and share the same eigenvectors, hence so does $A = B^{-1}C$, hence *A* is normal and

$$\lambda_k(A) = \frac{\lambda_k(C)}{\lambda_k(B)} = \frac{1 + \frac{1}{2} \operatorname{i} \mu \cos kx}{1 - \frac{1}{2} \operatorname{i} \mu \cos kx} \quad \Rightarrow \quad |\lambda_k(A)| = 1, \qquad k = 1...M.$$

So, Crank–Nicolson is again stable for all $\mu > 0$.

Euler for advection equation

Finally, consider the Euler method for advection equation

$$u_m^{n+1} - u_m^n = \mu(u_{m+1}^n - u_m^n), \qquad m = 1...M.$$

We have $\mathbf{u}^{n+1} = A\mathbf{u}^n$, where

$${\sf A} = \left[egin{array}{cccc} 1-\mu & \mu & & \ & 1-\mu & \ddots & \ & & \ddots & \mu & \ & & 1-\mu \end{array}
ight],$$

but A is not normal, and although its eigenvalues are bounded by 1 for $\mu \leq 2$, it is the spectral radius of AA^T that matters, and we have $\rho(AA^T) \approx (|1 - \mu| + |\mu|)^2$, so that the method is stable only if $\mu \leq 1$.