

# Phases of holographic matter

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Based on various works with:

Chris Herzog, Diego Hofman, Gary Horowitz,  
Pavel Petrov, Joe Polchinski, Eva Silverstein,  
Alireza Tavanfar, David Tong, David Vegh

A less compressed, but also not too long version of what I'm about to say can be found in

arXiv:1106.4324 [hep-th]

## Very brief condensed matter motivation

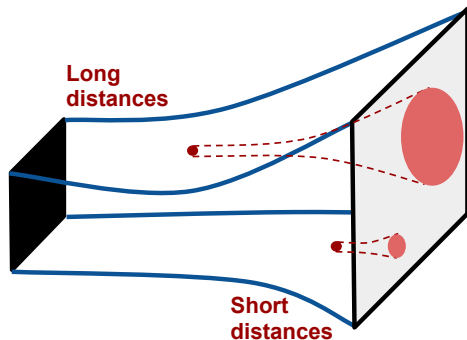
- QED in 2+1 dimensions at finite chemical potential  $\mu$  looks simple:

$$\mathcal{L} \sim F^2 + \bar{\Psi} [\Gamma \cdot (\partial + A) + \Gamma^0 \mu] \Psi.$$

- Chemical potential  $\Rightarrow$  Fermions form a Fermi surface.
- Low energy physics is strongly interacting [e.g. S-S. Lee '09].
- Difficult/impossible to answer basic questions:
  - At what temperature is there a superconducting instability?
  - What do correlators look like at low energies?

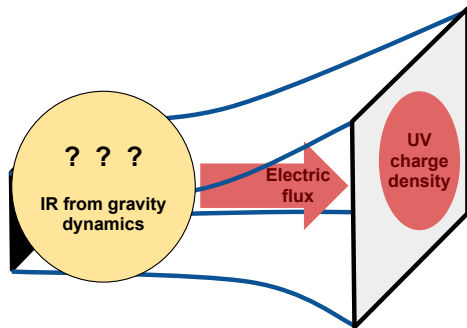
# Holography digested

- In certain quantum field theories the locality of the renormalisation group flow can be usefully geometrised.
- Append a dynamical extra dimension to the field theory.



# Holographic framework for finite density physics

- Change the problem, keep basic ingredients:
  - $SU(N)$  gauge fields coupled to matter with global symmetries.
- Holographically:
  - Finite chemical potential  $\Rightarrow$  charge density  $\Rightarrow$  flux at infinity.
  - Low energy dynamics  $\Rightarrow$  Solve the geometry in the interior.

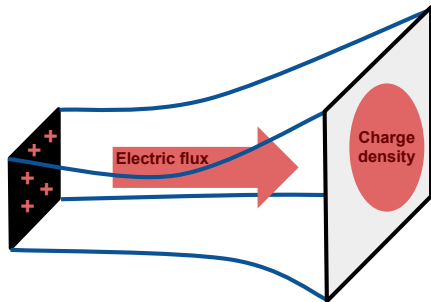


## Zero temperature, nonzero densities: I

- Simplest possible framework: Einstein-Maxwell

$$\mathcal{L} \sim \frac{1}{2\kappa^2} \left( R + \frac{6}{L^2} \right) - \frac{1}{4e^2} F^2.$$

- Answer: Extremal back hole.  $AdS_4 \rightarrow AdS_2 \times \mathbb{R}^2$ . [c.f. Gibbons '85]

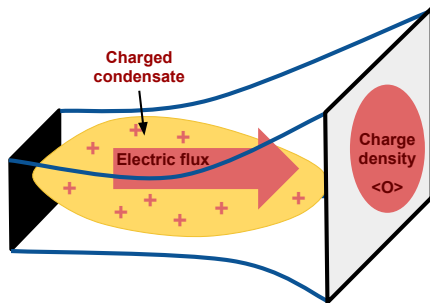
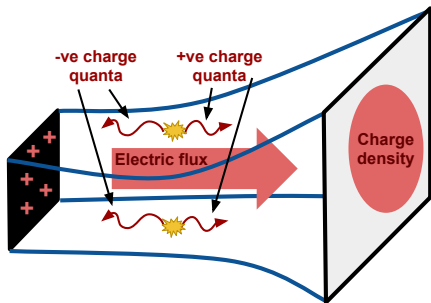


- $AdS_2 \times \mathbb{R}^2$  robust under higher derivative corrections. [e.g. Sen '05]

# Superradiant/Schwinger-like instabilities: charged bosons

[Gubser; Hartnoll-Herzog-Horowitz '08]

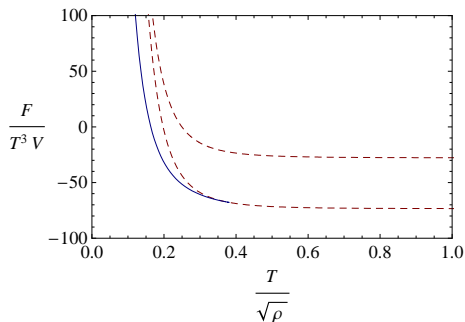
- Electric fields can destabilize scalar fields:  $m_{\text{eff.}}^2 \sim m^2 + g^{tt} A_t A_t$ .
- Pair production + charge neutralisation of black hole.



# Heating up holographic superconductors

[Hartnoll-Herzog-Horowitz '08]

- Heating up the black hole (away from extremality)
  - ⇒ Make the AdS box smaller.
  - ⇒ Eventually stabilise the tachyon.
- Condensate goes to zero at  $T_C$ . Second order transition.





# Superradiant/Schwinger-like instabilities: charged fermions

[Hartnoll-Polchinski-Silverstein-Tong '09; Hartnoll-Tavanfar '10]

- Fermions obey Pauli exclusion.
  - ⇒ Do not coherently occupy ground state.
  - ⇒ Einstein-Maxwell-Dirac.
- With a background  $A_t$ , fermions try to build up a Fermi sea.
- Difficult problem: self-consistently
  - Occupy all Dirac eigenstates in spacetime with energy  $E < \mu$ .
  - Backreact the energy-momentum tensor of these eigenstates.
- Simplifies in a large-charge/large mass limit.  
[Thomas-Fermi '27, Oppenheimer-Volkov '39]

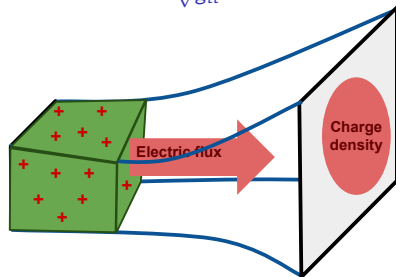
# Electron stars

[Hartnoll-Tavanfar '10; Hartnoll-Hofman-Tavanfar '11]

- Fluid description valid when

$$e^2 \sim \frac{\kappa}{L} \ll 1, \quad L^2 m^2 \sim \frac{L^2 e^2}{\kappa^2} \gg 1.$$

- Solve Einstein-Maxwell-charged ideal fluid equations.
- Local chemical potential:  $\mu_{\text{loc.}} = \frac{A_t}{\sqrt{g_{tt}}}$ .



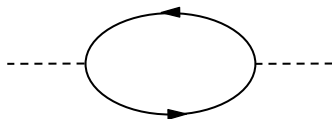
# Electron stars: IR scaling

[Hartnoll-Tavanfar '10; Hartnoll-Hofman-Tavanfar '11]

- The IR geometry of the star is Lifshitz [c.f. Kachru, Liu, Mulligan]

$$ds^2 = -\frac{dt^2}{r^{2z}} + \frac{dr^2}{r^2} + \frac{dx^2 + dy^2}{r^2}.$$

- This is good: Field theory computations also suggest the emergence of an IR dynamical critical exponent  $z > 1$ . 'Landau damping'.

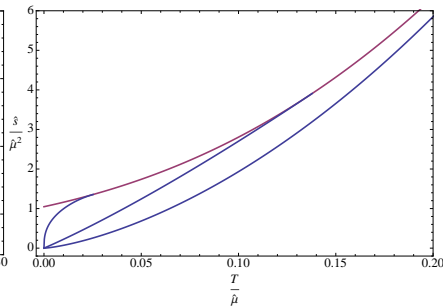
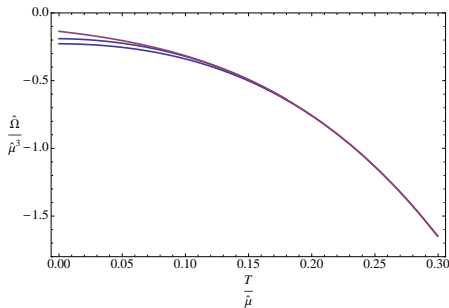


$$\langle A_x A_x \rangle(\omega, k) \sim \frac{1}{k^2 + \frac{|\omega|}{|k|}}, \quad (z = 3)$$

# Electron star birth: a third order transition

[Hartnoll-Petrov '11 (also Thorlacius et al. '11)]

- Free energy and entropy:



- Why?

$$\Delta\Omega \sim p(r)\Delta r \sim (\delta\mu)^{5/2}\Delta r \sim (\delta\mu)^3 \sim (T - T_C)^3.$$

(There is a cancellation in  $-p = \epsilon - \mu\sigma \sim (\delta\mu)^{5/2}$ .)

# The Luttinger count, fractionalisation and horizons

[Hartnoll-Hofman-Vegh '11 (also Iqbal, Liu, Mezei '11)]

- In the fluid limit – many closely spaced Fermi surfaces for fermion  $\psi$ .
- These satisfy a ‘Luttinger count’

$$\text{Charge in fermion fluid} = \sum \text{Fermi surface volumes of } \psi .$$

- Therefore

$$\text{Flux from horizon} = \langle J^t \rangle - \sum \text{Fermi surface volumes of } \psi .$$

- Flux from horizon is ‘missing charge’ in field theory Luttinger count.
- Characteristic of ‘Fractionalised Fermi liquids’. c.f.  $\psi = \text{Tr}\Psi\Phi$ .

# Conclusions

- Difficult questions in condensed matter theory motivate finite charge density holography.
- Extremal AdS black holes not especially robust in the presence of charged matter.
- Bosons  $\Rightarrow$  Holographic superconductors.
- Fermions  $\Rightarrow$  Electron stars.
- Emergent Lifshitz geometry in IR at zero temperature.
  - Landau damped gauge fields?
- Dichotomy: Flux from horizon vs. Flux from matter.