

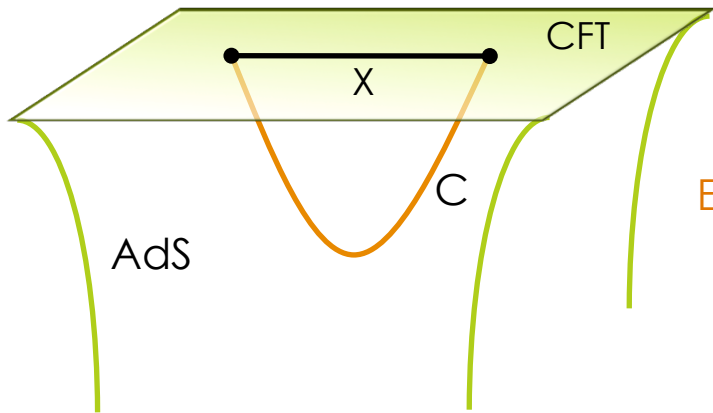
Anomalies and Gaps in HEE

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DAMTP, Cambridge

Holographic Entanglement Entropy



Conformal
Field theory

Anti-de Sitter
Gravity

Entanglement Entropy



Minimal areas

$$S_{EE} = -\text{Tr} \rho_X \log \rho_X$$

Ryu-Takayanagi formula

Long list of features ...

- ▣ Leading contribution yields area law.
- ▣ In a pure state, $S(X)=S(X^c)$.
- ▣ Strong subadditivity and other inequalities.
- ▣ Recover known results.
- ▣ Generalizations: susy, higher spins, ...
- ▣ There is a derivation from first principles in holography.

Long list of applications ...

- ▣ It characterizes d.o.f. of the theory and anomalies.
- ▣ Ties with renormalization group flows.
- ▣ Probe of phase transitions (e.g. confinement/deconfinement).
- ▣ Useful characterization of excited states (quenches).

GOAL for TODAY

- Interplay between **anomalies** and **entanglement entropy**
- Anomalies **probe** other aspects of the dual bulk geometry

FOCUS

- ▣ Gravitational anomalies in $\text{AdS}_3/\text{CFT}_2$
- ▣ Generalization of Ryu-Takayanagi
- ▣ Test in a dynamical setup

1. CFT_2 :
EE & anomalies revisited

with Detournay,
Iqbal & Perlmutter.
arXiv:1405.2792

2. AdS_3 :
Spinning probes

3. Dynamics:
Gapped systems in AdS_3

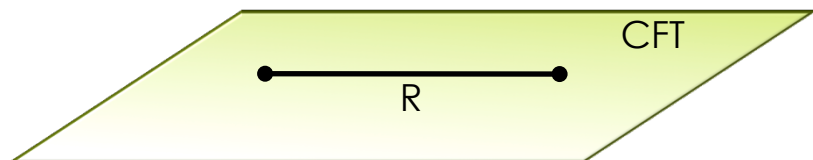
with Belin & Hung.
arXiv: soon!

OUTLINE

CFT₂

Entanglement entropy and anomalies revisited

The relationship between entanglement and anomalies is well known.
For example, single interval on the vacuum state:



$$S_{\text{EE}} = \frac{c}{3} \log \left(\frac{R}{\varepsilon} \right)$$

[Holzhey, Larsen & Wilczek]

Explicit relation between **conformal anomaly**
and **entanglement entropy**

Important assumption here:

$$c = \frac{c_L + c_R}{2} \quad c_L = c_R$$

Is entanglement entropy sensitive to **diffeomorphism anomalies**?

Diffeomorphism anomaly

In the OPE expansion of the stress tensor:

$$T(z)T(0) \sim \frac{c_L/2}{z^4} + \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots$$

$$\bar{T}(\bar{z})\bar{T}(0) \sim \frac{c_R/2}{\bar{z}^4} + \frac{\bar{T}(0)}{\bar{z}^2} + \frac{\bar{\partial}\bar{T}(0)}{\bar{z}} + \dots$$

$$c_L \neq c_R$$

Simple examples: chiral matter or holomorphic CFTs.

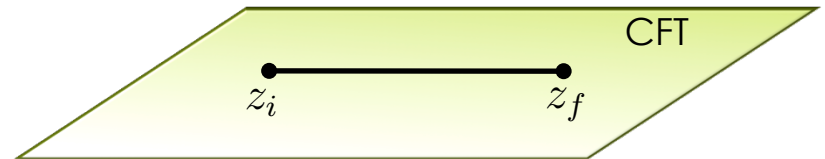
The anomaly can be presented in two ways:

- ▣ Stress tensor is symmetric, but not conserved
- ▣ Stress tensor is conserved, but not symmetric

Twist fields & EE

$$S_{\text{EE}} = \lim_{n \rightarrow 1} S_n = -\text{Tr} \rho \log \rho$$

$$S_n = \frac{1}{1-n} \log \text{Tr} \rho^n$$



$$\text{Tr} \rho^n = \langle \Phi_+(z_i) \Phi_-(z_f) \rangle_{\mathcal{C}^n / \mathbb{Z}_n}$$

Twist field. For simplicity only consider a single interval

[Cardy, Castro-Alvaredo & Doyon; Cardy & Calabrese]

Include diff anomaly: Conformal dimensions of the twist fields are

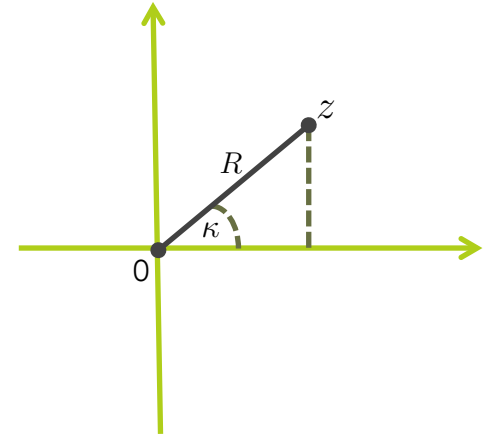
$$\begin{aligned} \Delta &= h + \bar{h} \\ &= \frac{c_L + c_R}{24} \left(n - \frac{1}{n} \right) \end{aligned}$$

$$\begin{aligned} s &= h - \bar{h} \\ &= \frac{c_L - c_R}{24} \left(n - \frac{1}{n} \right) \end{aligned}$$

Results

Vacuum state & boosted interval

$$\begin{aligned} S_{\text{EE}} &= \frac{c_L}{12} \log \left(\frac{z}{\varepsilon} \right) + \frac{c_R}{12} \log \left(\frac{\bar{z}}{\varepsilon} \right) \\ &= \frac{c_L + c_R}{6} \log \left(\frac{R}{\varepsilon} \right) - \frac{c_L - c_R}{6} \kappa \end{aligned}$$



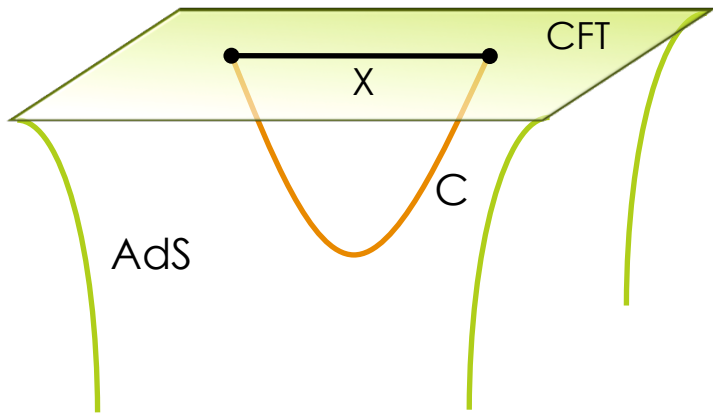
Thermal state: high temperature, finite angular velocity

$$S_{\text{EE}} = \frac{c_L}{6} \log \left(\frac{\beta_L}{\pi \varepsilon} \sinh \frac{\pi R}{\beta_L} \right) + \frac{c_R}{6} \log \left(\frac{\beta_R}{\pi \varepsilon} \sinh \frac{\pi R}{\beta_R} \right)$$

AdS_3

Spinning particles and anomalies

Einstein Hilbert gravity



Ryu-Takayanagi formula

$$S_{\text{EE}} = -\text{Tr}(\rho_X \log \rho_X) = \frac{L_C}{4G_3}$$

It cannot capture the diffeomorphism anomaly: not enough data to have unbalanced left/right central charges.

Topological massive gravity

$$S_{\text{TMG}} = \frac{1}{16\pi G_3} \int d^3x \sqrt{-g} \left(R + \frac{2}{\ell^2} \right) + \frac{1}{32\pi G_3 \mu} \int d^3x \sqrt{-g} \underbrace{\epsilon^{\lambda\mu\nu} \Gamma_{\lambda\sigma}^{\rho} \left(\partial_{\mu} \Gamma_{\rho\nu}^{\sigma} + \frac{2}{3} \Gamma_{\mu\tau}^{\sigma} \Gamma_{\nu\rho}^{\tau} \right)}_{\text{Gravitational Chern-Simons term}}$$

Asymptotically **AdS₃** spacetime:

$$c_L = \frac{3\ell}{2G_3} \left(1 - \frac{1}{\mu\ell} \right) \quad c_R = \frac{3\ell}{2G_3} \left(1 + \frac{1}{\mu\ell} \right)$$

HEE revisited

Task: find the generalization of R-T in the presence of diff anomalies

Two routes (that complement each other):

- ▣ Conical singularity method applied to TMG
[Lewkowycz & Maldacena; Dong; Camps]
- ▣ Design a bulk probe that captures the right CFT physics

Spinning probes

Hint: Recall, twist field has quantum numbers

$$\Delta = \frac{c_L + c_R}{12} (n - 1) + O(n - 1)^2$$

$$s = \frac{c_L - c_R}{12} (n - 1) + O(n - 1)^2$$

$$c_L = \frac{3\ell}{2G_3} \left(1 - \frac{1}{\mu\ell} \right)$$

$$c_R = \frac{3\ell}{2G_3} \left(1 + \frac{1}{\mu\ell} \right)$$

➤ Ryu-Takayanagi formula

Geodesic equation $\rightarrow$ Massive probe $\rightarrow$ $\mathfrak{m} = \frac{\Delta}{\ell} = \frac{1}{4G_3}$

➤ TMG EE formula

Spinning probe $\rightarrow$ $\mathfrak{m} = \frac{\Delta}{\ell} = \frac{1}{4G_3}$ $\mathfrak{s} = s = \frac{1}{4\mu G_3}$

MPD equation

Motion of a **spinning particle** is described by the Mathisson-Papapetrou-Dixon equations:

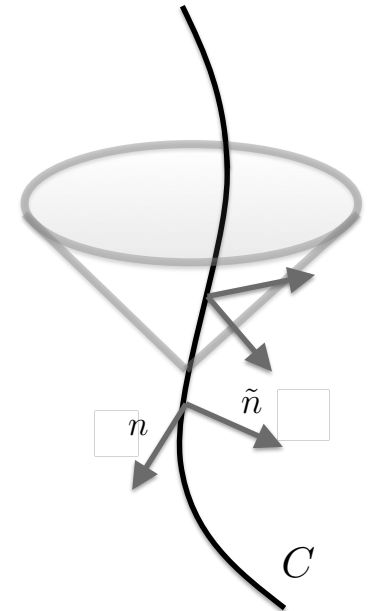
$$\nabla [\mathfrak{m} v^\mu + v_\rho \nabla s^{\mu\rho}] = -\frac{1}{2} v^\nu s^{\rho\sigma} R^\mu_{\nu\rho\sigma}$$

$$s^{\mu\nu} = \mathfrak{s} (n^\mu \tilde{n}^\nu - \tilde{n}^\mu n^\nu)$$

Action principle

$$S_{\text{probe}} = \int_C ds \left(\mathfrak{m} \sqrt{g_{\mu\nu} v^\mu v^\nu} + \mathfrak{s} \tilde{n} \cdot \nabla n \right) + S_{\text{constraints}}$$

Orthogonality of
normal vectors.
Vanishes on-shell.



The Proposal for HEE

$$S_{\text{EE}} = \int_C ds \left(\mathfrak{m} \sqrt{g_{\mu\nu} v^\mu v^\nu} + \mathfrak{s} \tilde{n} \cdot \nabla n \right)$$

$$\mathfrak{m} = \frac{\Delta}{\ell} = \frac{1}{4G_3} \qquad \mathfrak{s} = s = \frac{1}{4\mu G_3}$$

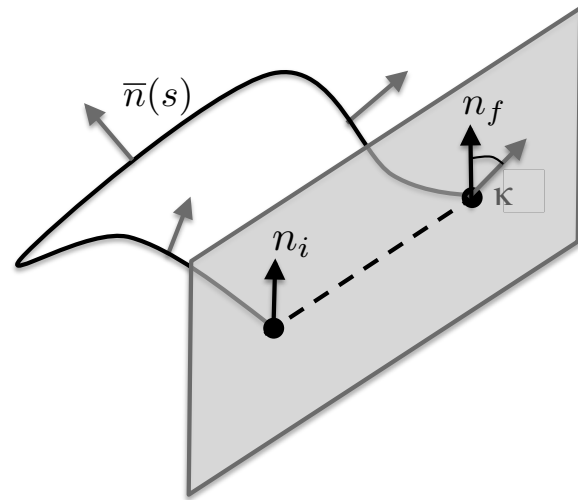
Features:

- We reproduced it from the cone action. However the dynamics is less obvious.
- Spin does not add a new bulk d.o.f. We gave life to a pure gauge mode to make the answer look covariant.
- The action now measures the **torsion** of the curve. It is related the **framing anomaly** in Chern-Simons theory as well.

Sanity checks

Poincare AdS₃ & boosted interval

$$S_{\text{EE}} = \frac{c_L + c_R}{6} \log \left(\frac{R}{\varepsilon} \right) - \frac{c_L - c_R}{6} \kappa$$



BTZ BH: high temperature, finite angular velocity

$$S_{\text{EE}} = \frac{c_L}{6} \log \left(\frac{\beta_L}{\pi \varepsilon} \sinh \frac{\pi R}{\beta_L} \right) + \frac{c_R}{6} \log \left(\frac{\beta_R}{\pi \varepsilon} \sinh \frac{\pi R}{\beta_R} \right)$$

Dynamics

Gapped systems in the presence of anomalies

Holographic gapped systems

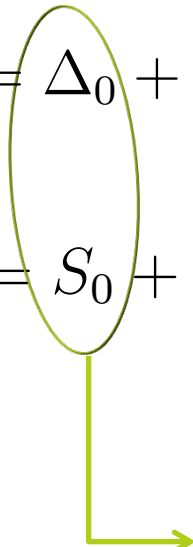
Two basic properties that we seek in a gapped system are:

Spectrum

$$\Delta E = \Delta_0 + \dots$$

Entanglement

$$S_{EE} = S_0 + \dots$$

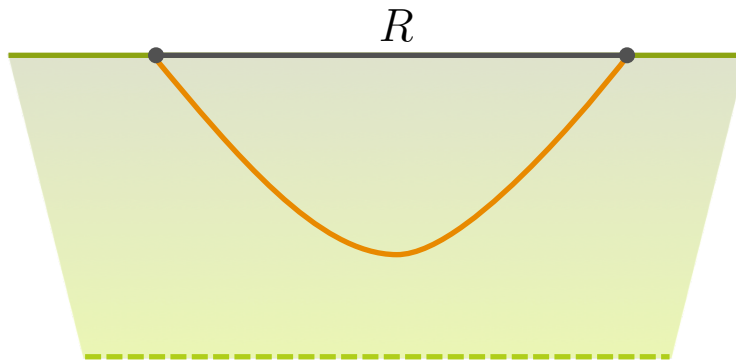


Finite and independent
of the size of the system
as we uncompactify

Holographic gapped systems

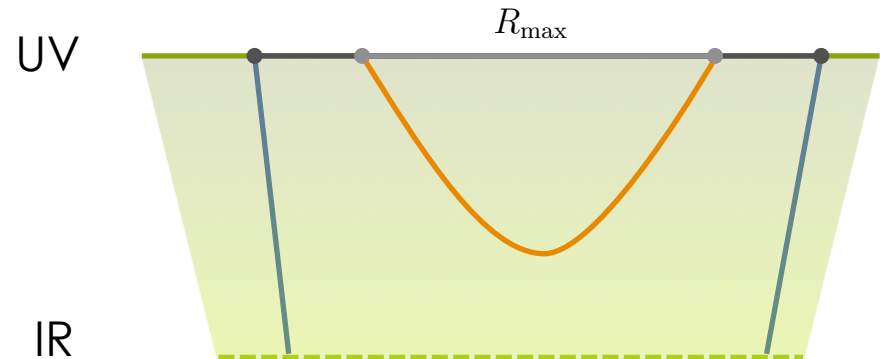
How to tell if a system is gapped in the IR using HEE?

[Klebanov, Kutasov, Murugan; Liu, Mezei]



NOT GAPPED:

Connected solution exists
for any size R and
dominates.



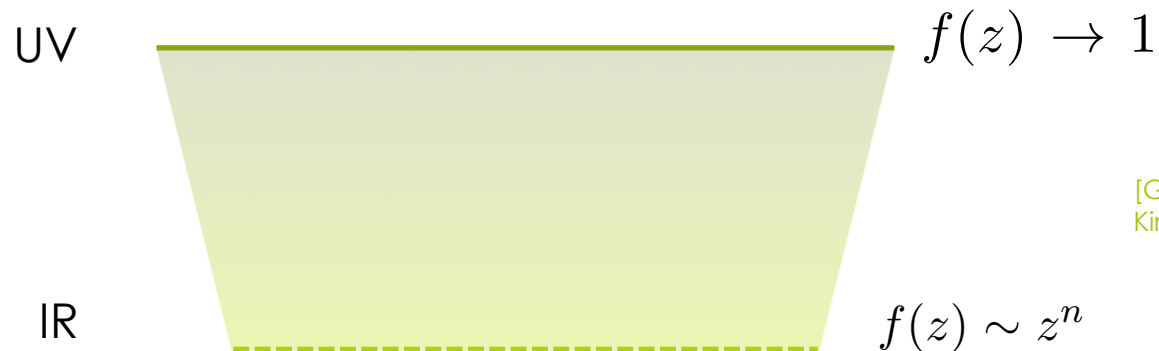
GAPPED:

Connected path ceases to
exist above a max length.
Disconnected solution
takes over.

Domain walls

Geometry of the gapped system:

$$ds^2 = \frac{\eta_{ij} dx^i dx^j}{z^2} + \frac{dz^2}{z^2 f(z)}$$



[Gubser; Freedman et al;
Kiristis et al; ...]

NOTE: These solutions, for arbitrary $f(z)$, have vanishing Cotton tensor

A “gapped” domain wall is allowed in a theory with a diffeomorphism anomaly!

Gaps & Diff anomalies

Anomalies are powerful and wise.

Chiral anomalies are protected by RG.
Chiral matter will not become massive in the IR. No gaps!

2d 't Hooft anomaly matching condition: $c_L - c_R$

How does holography capture this simple and
yet robust statement!?

MPD solutions

This is a challenge to our spinning probe!
Does HEE capture anomaly matching correctly?

Functional: $S_{\text{EE}} = \int_C ds \left(\mathfrak{m} \sqrt{g_{\mu\nu} v^\mu v^\nu} + \mathfrak{s} \tilde{n} \cdot \nabla n \right)$

MPD Dynamics: $\nabla [\mathfrak{m} v^\mu + v_\rho \nabla s^{\mu\rho}] = -\frac{1}{2} v^\nu s^{\rho\sigma} R^\mu_{\nu\rho\sigma}$

$$s^{\mu\nu} = -\mathfrak{s} \epsilon^{\mu\nu\lambda} v_\lambda$$

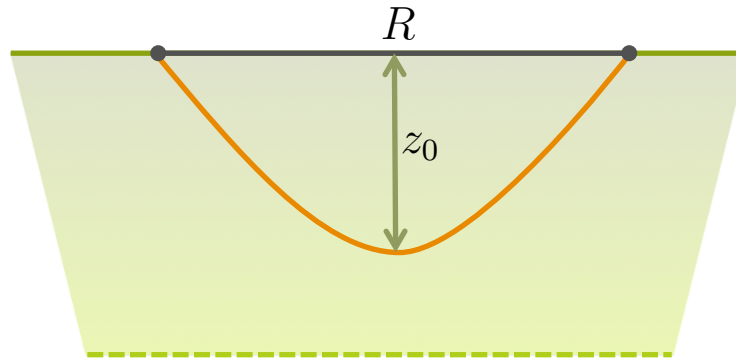
Task: Characterize solutions to MPD in the domain wall background

$$ds^2 = \frac{\eta_{ij} dx^i dx^j}{z^2} + \frac{dz^2}{z^2 f(z)}$$

$$\nabla [\mathfrak{m}v^\mu + v_\rho \nabla s^{\mu\rho}] = -\frac{1}{2}v^\nu s^{\rho\sigma} R^\mu_{\nu\rho\sigma}$$

Comments:

- Equation highly non-linear.
- Exact solutions: either disconnected or only valid in AdS_3 .
- Connected solutions are **not static**. The probe wants to twist.
- Build connected solutions perturbatively in the spin.



$$R \sim z_0 + O(z_0^2) , \quad z_0 \text{ small}$$

$$R \sim z_0^{1-n/2} + s^2 z_0^{1+n/2} + \cdots , \quad z_0 \text{ large}$$

Connected solution exists for any size of the boundary interval.

Outlook

What is next?

Anomalies

- ▣ Generalization to higher dimensions: spinning membranes?
- ▣ Generalization in $4k+2$ QFTs
- ▣ Mixed gauge-gravitational anomalies

Gaps

- ▣ Work in progress: linearized spectrum should corroborate results.
- ▣ Exact connected solutions to MPD equations.
- ▣ CFT interpretation to all MPD solutions?



Thank you!