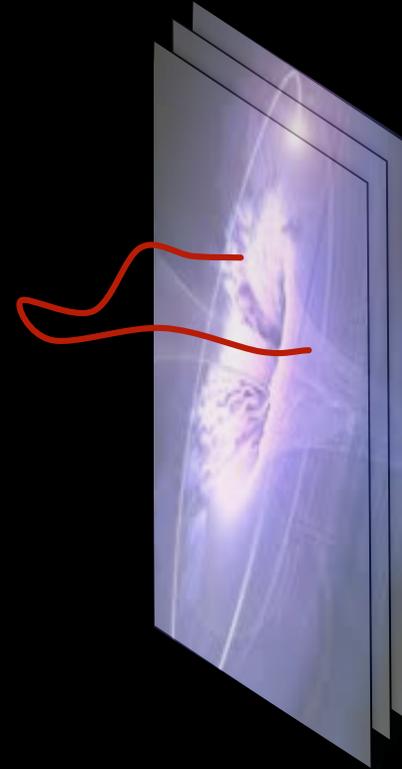
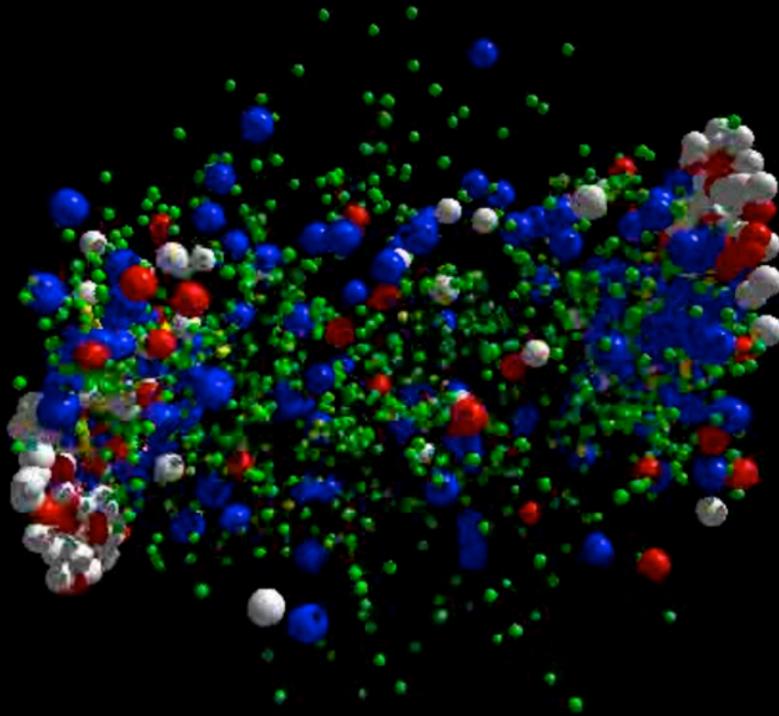


Strong Coupling Isotropization Simplified



David Mateos
ICREA & University of Barcelona

Work with M.Heller, W. van der Schee, M. Spalinski & D. Trancanelli

Dedicated to Paul Townsend

- Paul has been a friend and a mentor throughout my career.
- Met at the beginning of my PhD in 1997 -- hard times!
- Paul was visiting Barcelona for a few months.
- Gave lectures on M-theory and things fell into place.
- Wrote a paper before he left Barcelona (with Gomis & Simon).

Dedicated to Paul Townsend

- In 1998 Paul hosted me as a visiting student at DAMTP...
and at his house!
- During that visit we wrote another paper (with Gauntlett, Koehl and Zamaklar).
- In 2000-2002 I joined DAMTP as a postdoc.
- We wrote several papers on “supertubes”.

Dedicated to Paul Townsend

- One day Paul walked into my office with:

Adding flavor to AdS/CFT

Andreas Karch and Emanuel Katz

Department of Physics, University of Washington, Seattle, WA 98195, USA

(karch@phys.washington.edu)

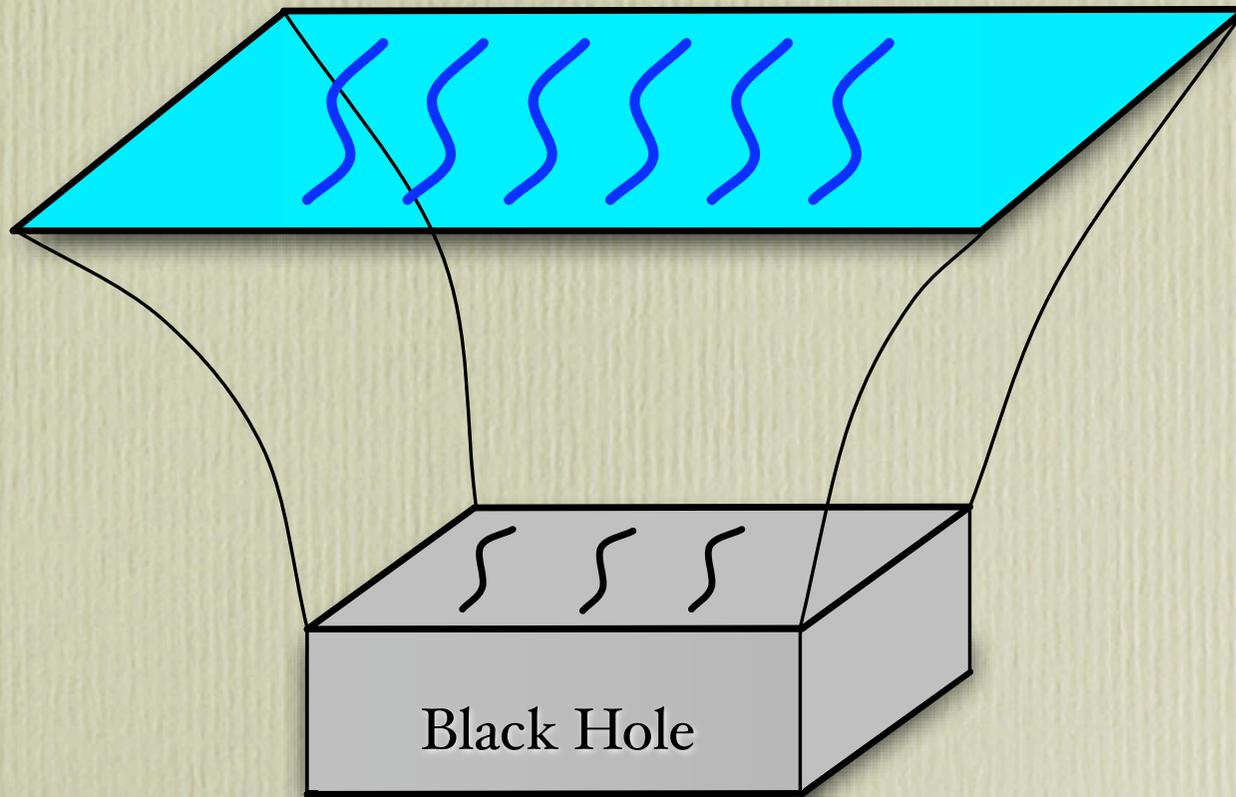
(amikatz@phys.washington.edu)

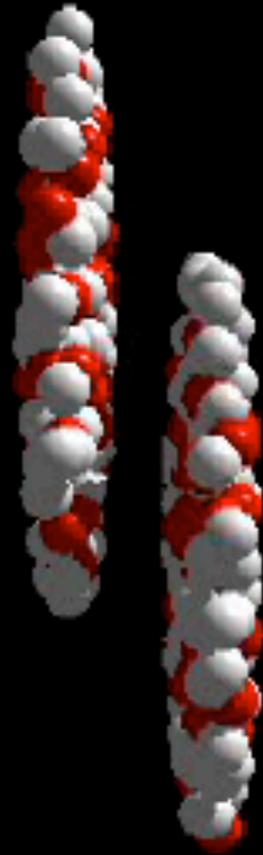
Coupling fundamental quarks to QCD in the dual string representation corresponds to adding the open string sector. Flavors therefore should be represented by space-time filling D-branes in the dual 5d closed string background. This requires several interesting properties of D-branes in AdS. D-branes have to be able to end in thin air in order to account for massive quarks, which only live in the UV region. They must come in distinct sets, representing the chiral global symmetry, with a bifundamental field playing the role of the chiral condensate. We show that these expectations are born out in several supersymmetric examples. To analyze most of these properties it is not necessary to go beyond the probe limit in which one neglects the backreaction of the flavor D-branes.

Dedicated to Paul Townsend

- Beginning of my interest on applications of gauge/string duality to QCD.
- Today: Applications to out-of-equilibrium physics of QGP.

Last decade: Near equilibrium QGP





Out of equilibrium

Out-of-equilibrium QFT



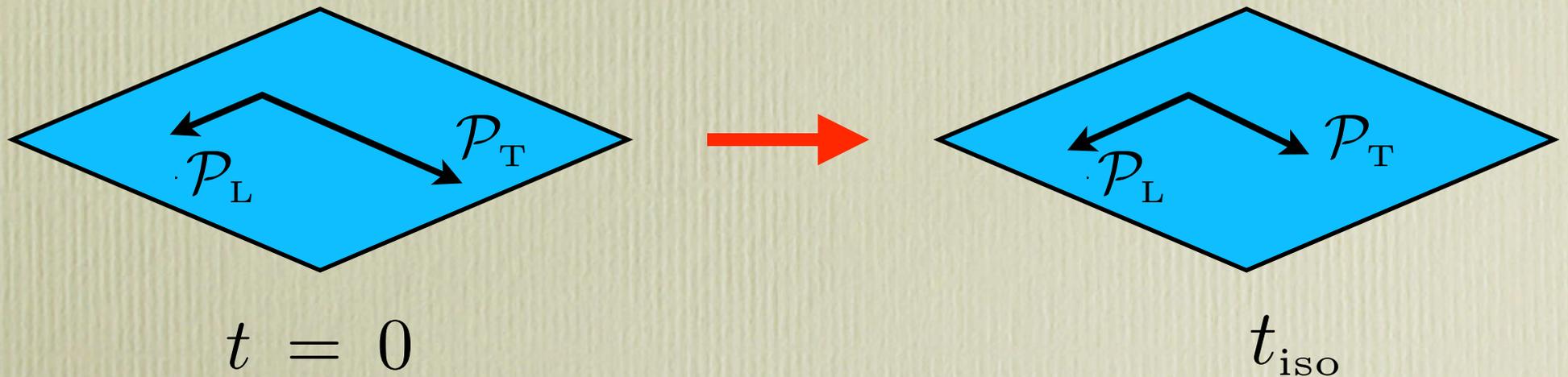
Classical Dynamical GR in AdS

Remember

- QCD dual is beyond supergravity.
- Do not try to do *precision* QCD physics.
- Search for physical insights.

In the context of HIC

- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma
(e.g. $N=4$ SYM plasma).

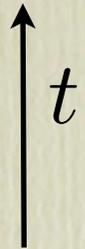


In the context of HIC

- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma (e.g. N=4 SYM plasma).
- Homogeneity: $\partial_\mu T^{\mu\nu} = 0 \rightarrow \partial_t T^{00} = 0 \rightarrow \mathcal{E}_i = \mathcal{E}_f$
- No hydrodynamics: $\omega \rightarrow 0$ as $q \rightarrow 0$
- Only “quasi-normal modes” (QNM).

Posing the problem: Causal Structure

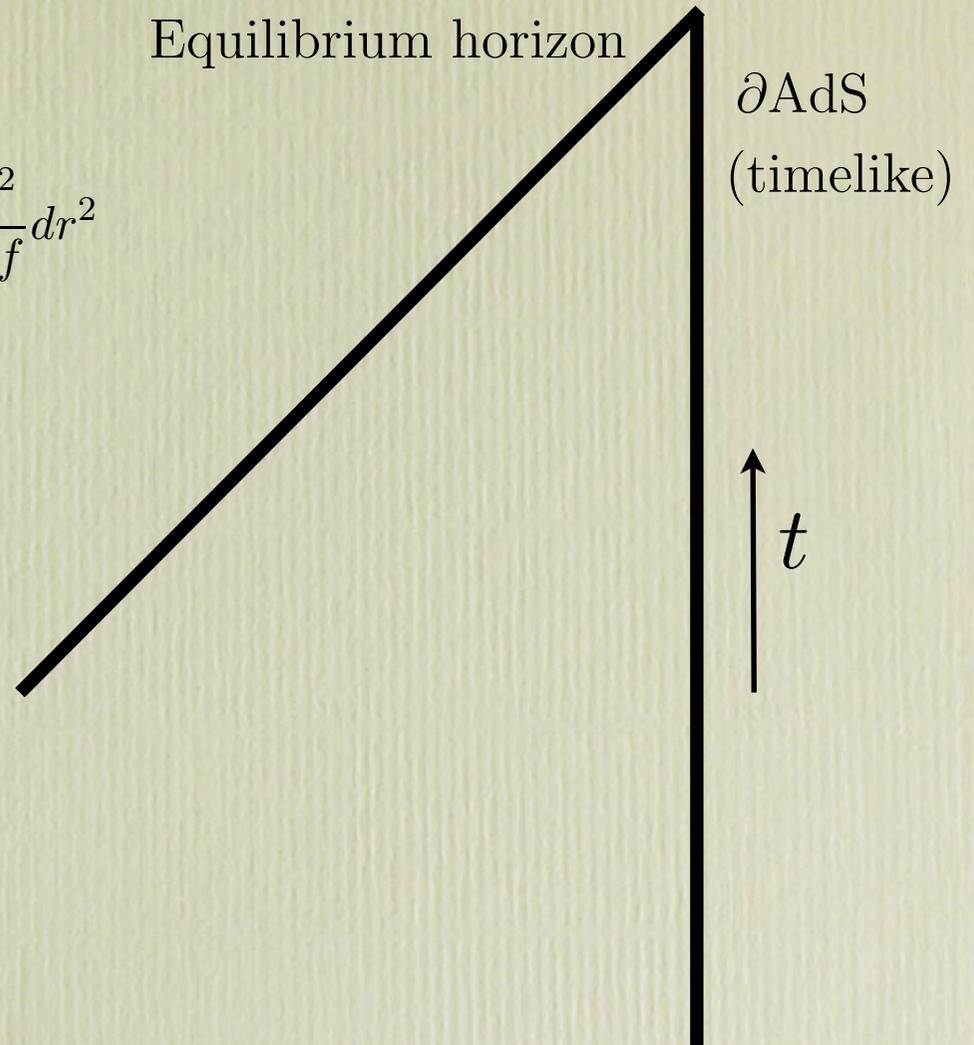
∂AdS
(timelike)



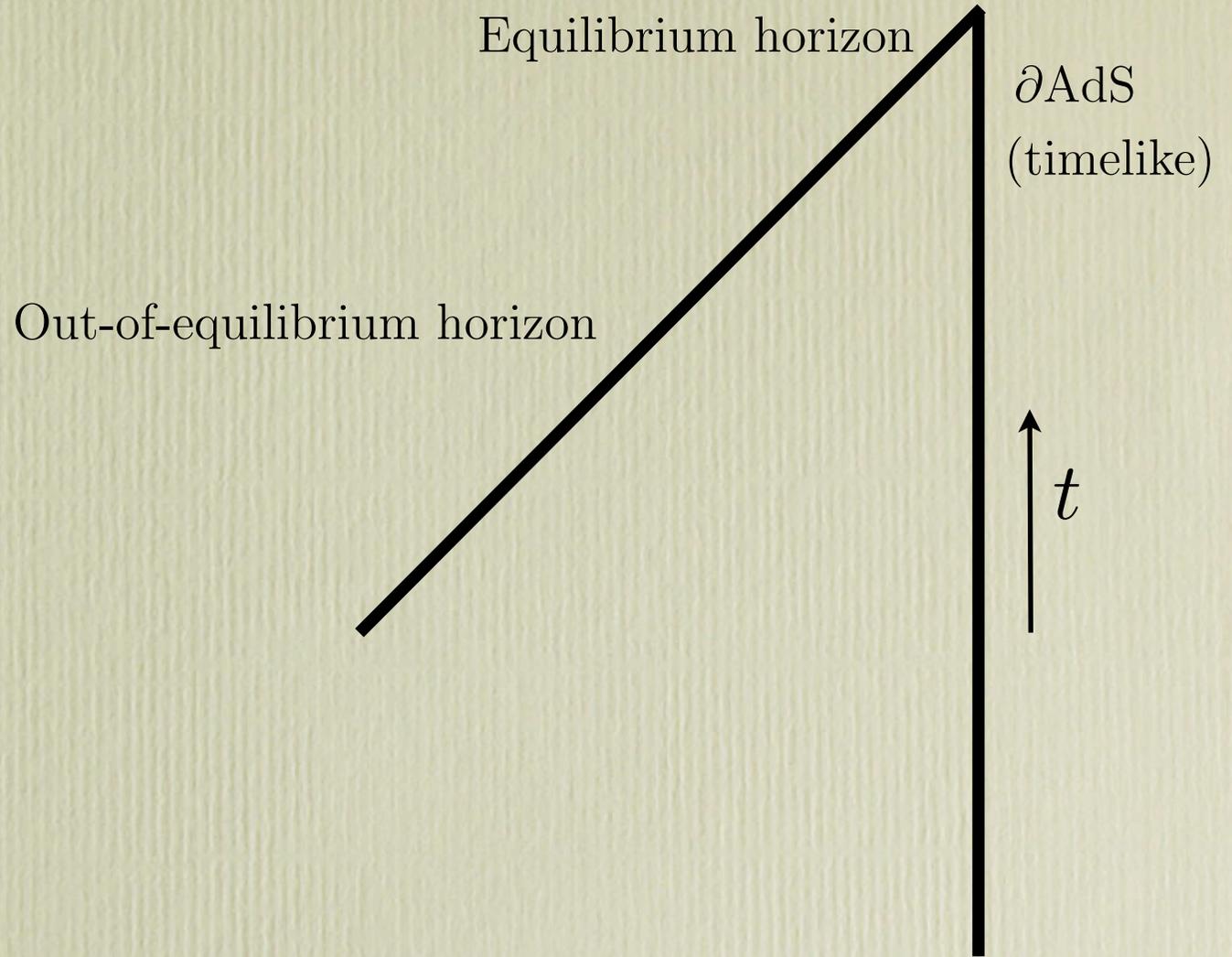
Posing the problem: Causal Structure

$$ds^2 = \frac{r^2}{R^2} (-f dt^2 + dx_1^2 + dx_2^2 + dx_3^2) + \frac{R^2}{r^2 f} dr^2$$

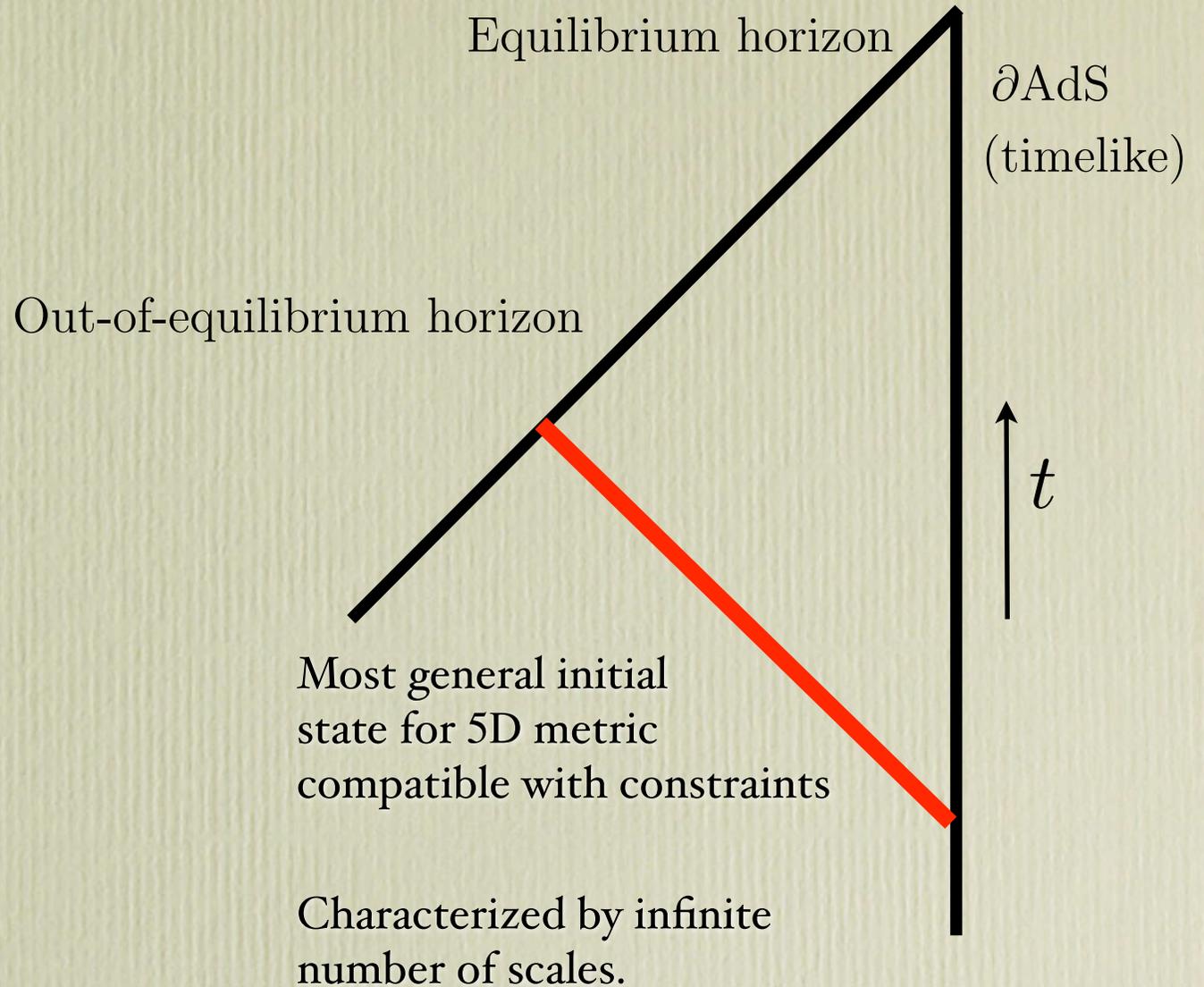
$$f(r) = 1 - \frac{r_0^4}{r^4}$$



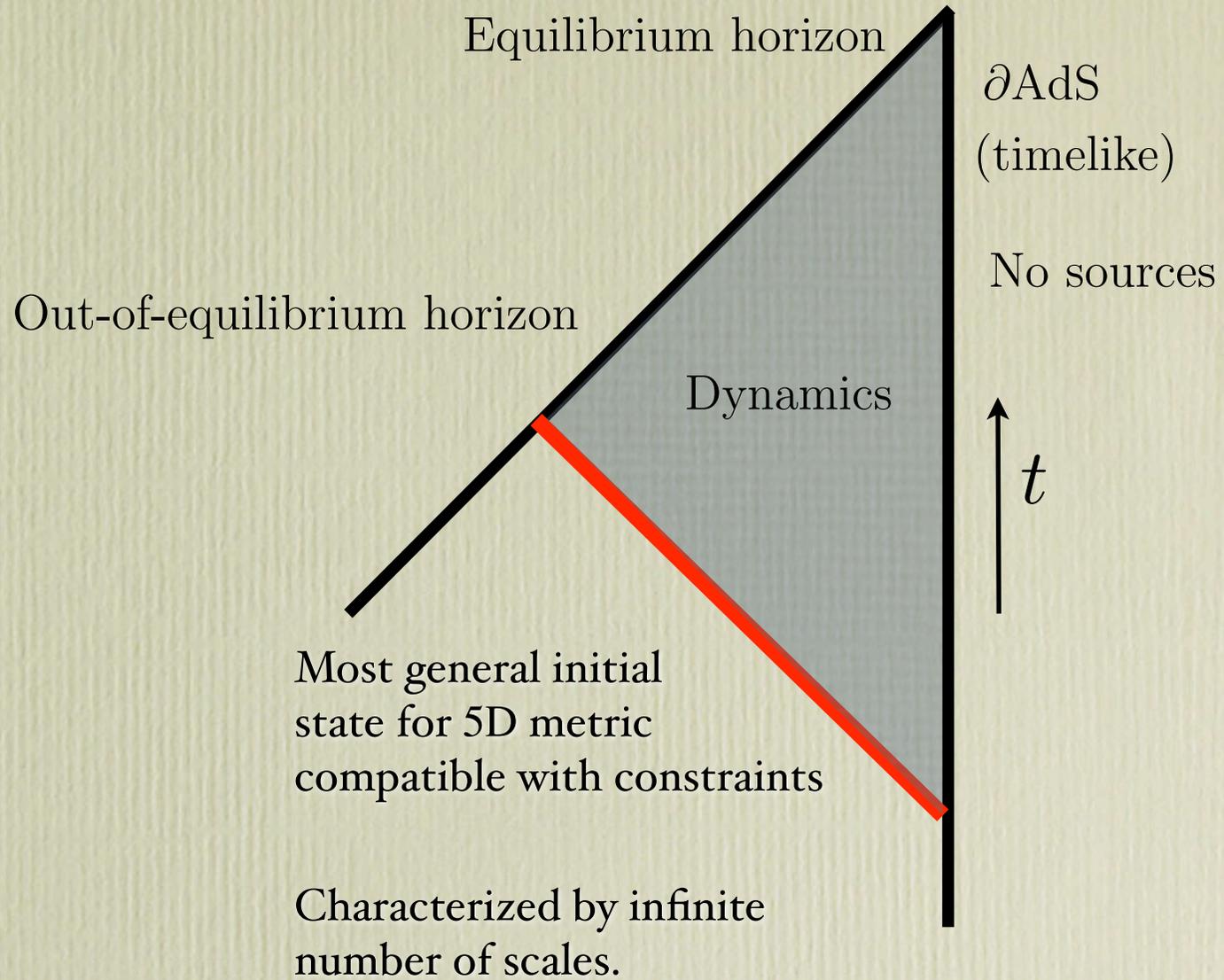
Posing the problem: Causal Structure



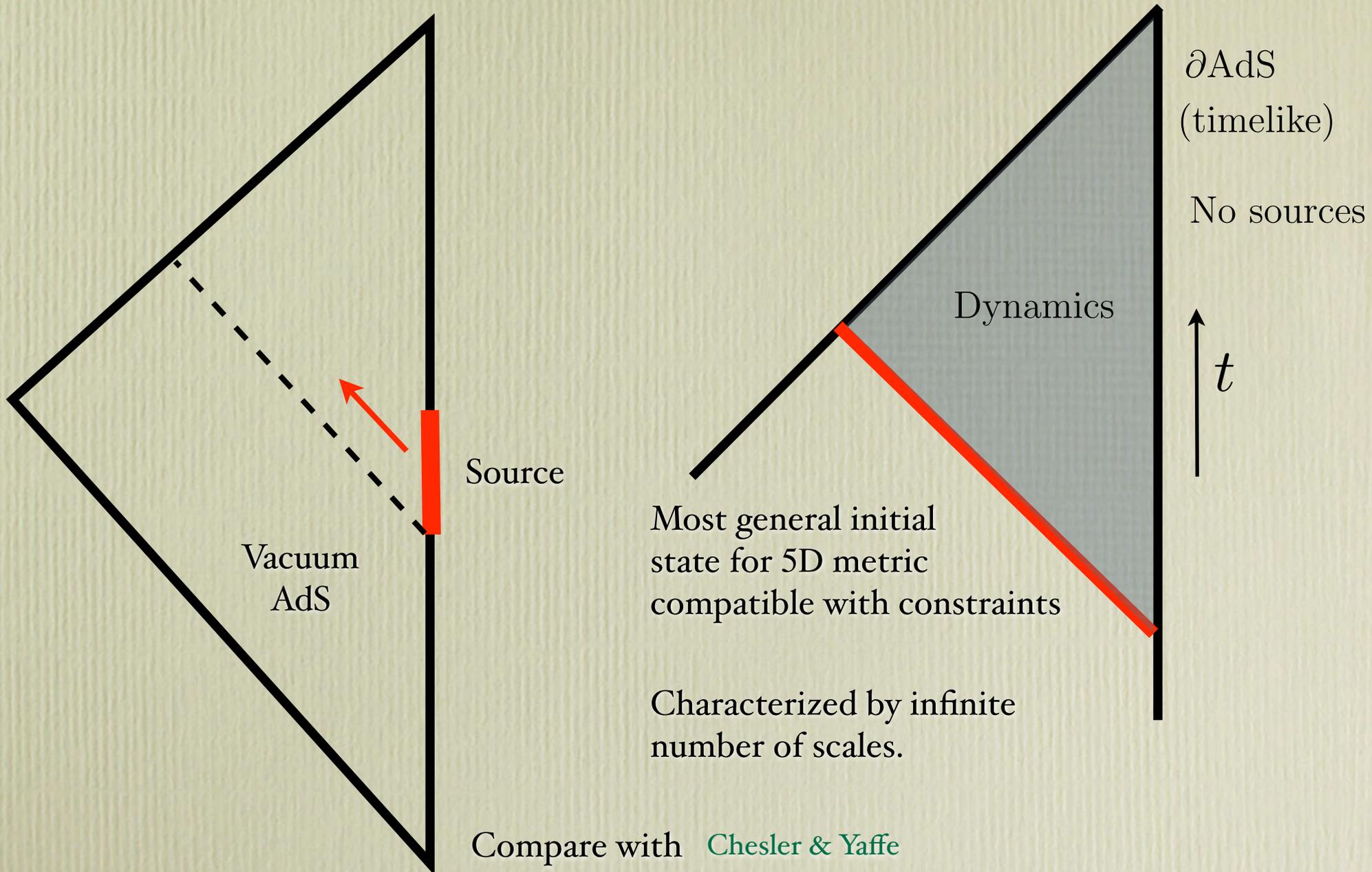
Posing the problem: Causal Structure



Posing the problem: Causal Structure



Posing the problem: Causal Structure



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

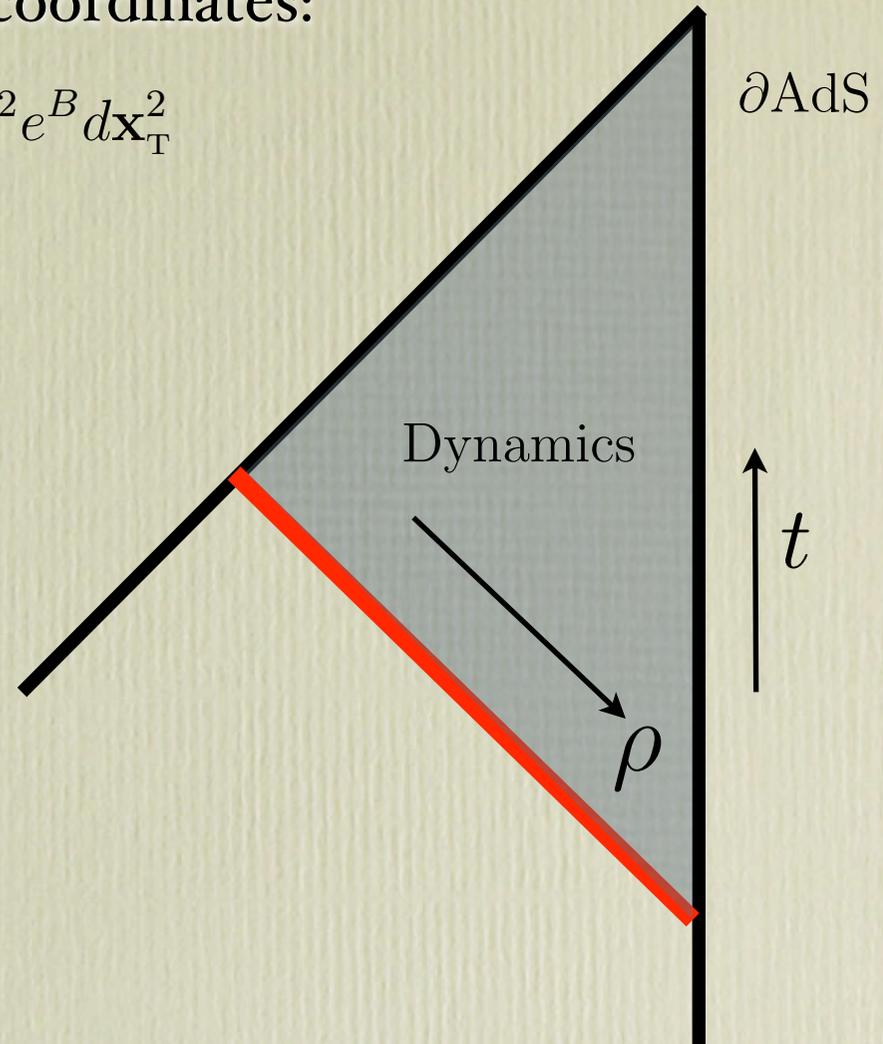
$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B dx_T^2$$

A, Σ, B functions of t, ρ only.

- In equilibrium:

$$A = \rho^2 \left(1 - \frac{\rho_h^4}{\rho^4}\right), \quad \Sigma = \rho$$

$$B = 0, \quad \rho_h = \pi T$$



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Near-boundary fall-off determines:

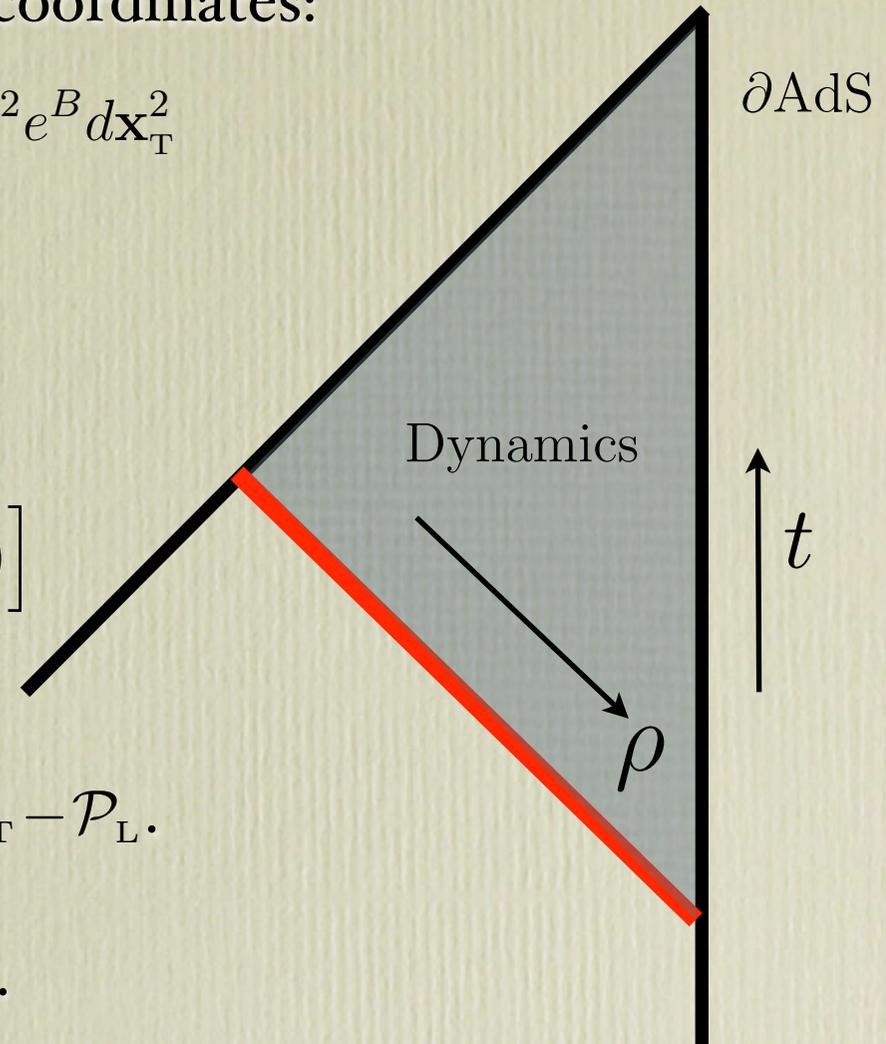
$$\langle T_{\mu\nu} \rangle = \frac{N_c^2}{2\pi^2} \text{diag} \left[\mathcal{E}, \mathcal{P}_L(t), \mathcal{P}_T(t), \mathcal{P}_T(t) \right]$$

- In particular, B determines $\Delta\mathcal{P} = \mathcal{P}_T - \mathcal{P}_L$.

- Initial state specified by $B(t = 0, \rho)$.

- CFT: Must specify anisotropy distribution in modes.

- Infinitely-many-scales problem.



Time evolution

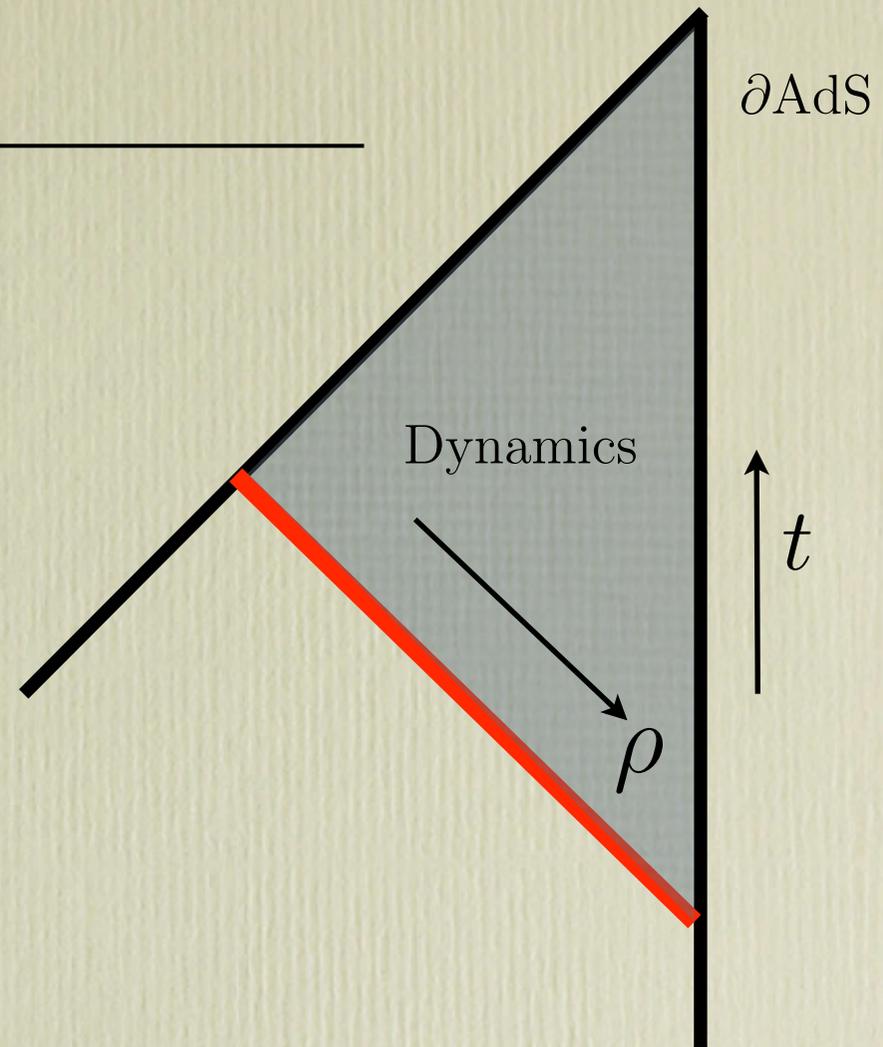
- In equilibrium:

$$A = \rho^2 \left(1 - \frac{\rho_h^4}{\rho^4}\right), \quad \Sigma = \rho$$

$$B = 0, \quad \rho_h = \pi T$$

- Evolve initial state according to:

- ▶ Full, non-linear EEQs.
- ▶ EEQs linearized around final equilibrium state.



Motivation: Close-limit Approximation

Price & Pullin '94

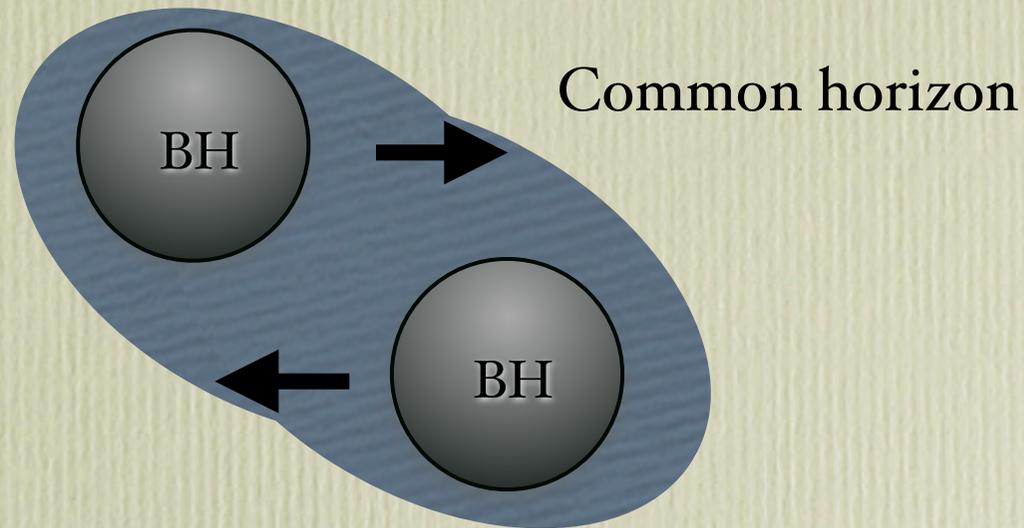
- BH collision in asymptotically flat 4D general relativity:



Motivation: Close-limit Approximation

Price & Pullin '94

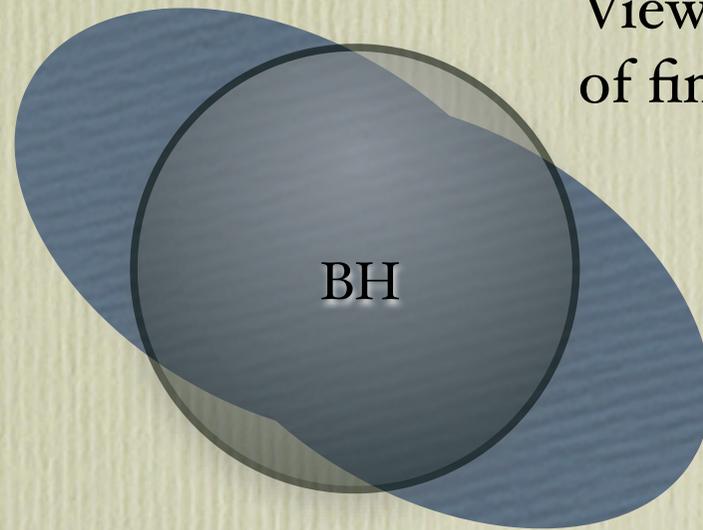
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Motivation: Close-limit Approximation

Price & Pullin '94

- BH collision in asymptotically flat 4D general relativity:

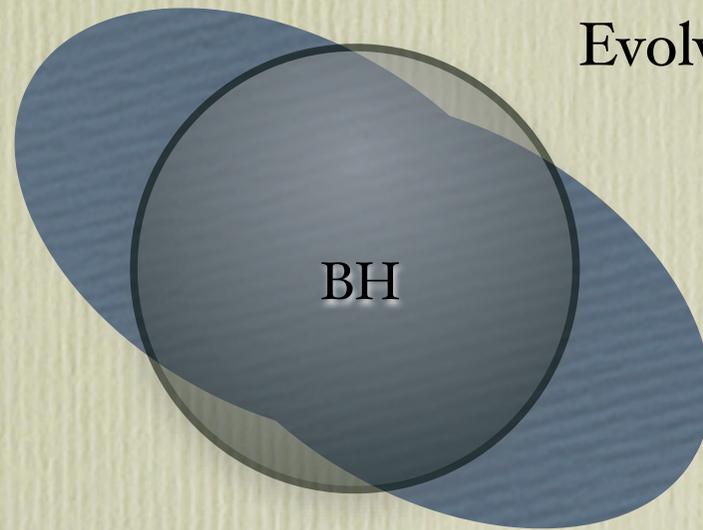


View as perturbation
of final BH

Motivation: Close-limit Approximation

Price & Pullin '94

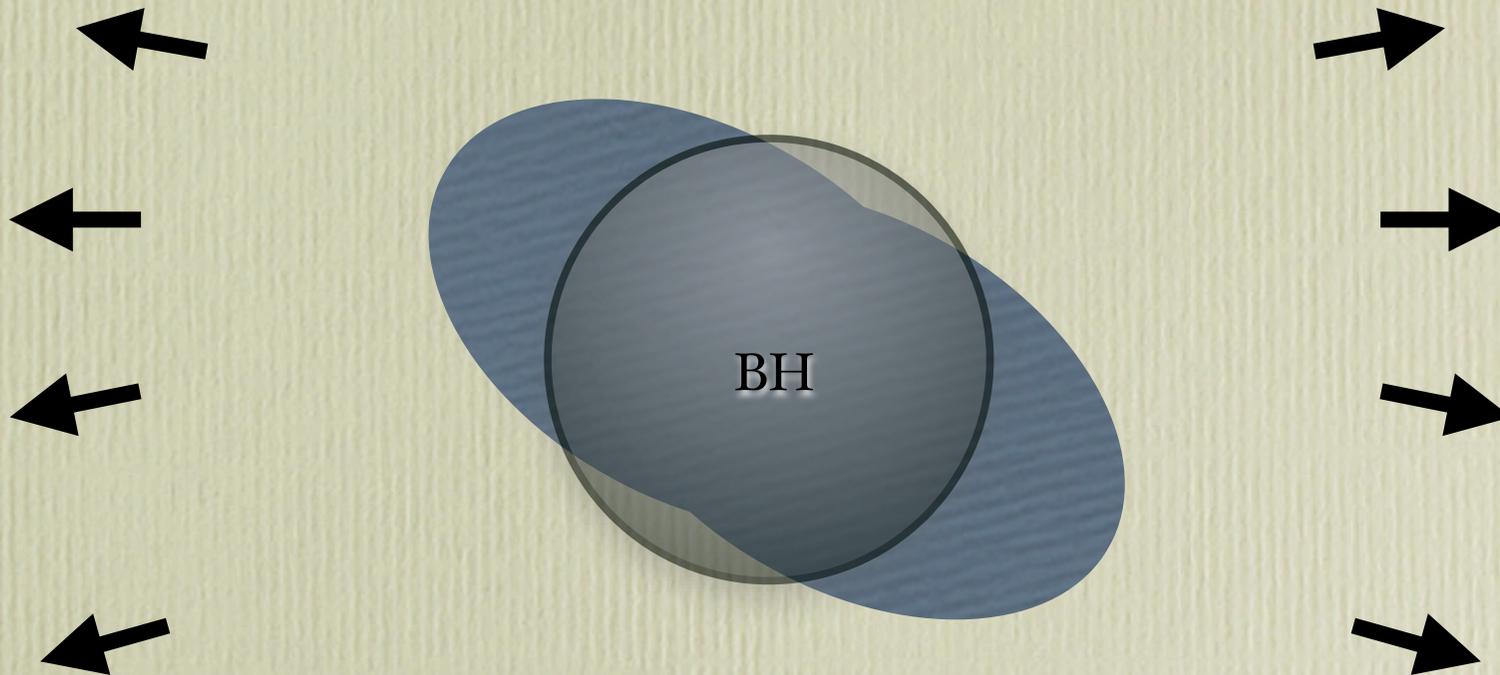
- BH collision in asymptotically flat 4D general relativity:



Motivation: Close-limit Approximation

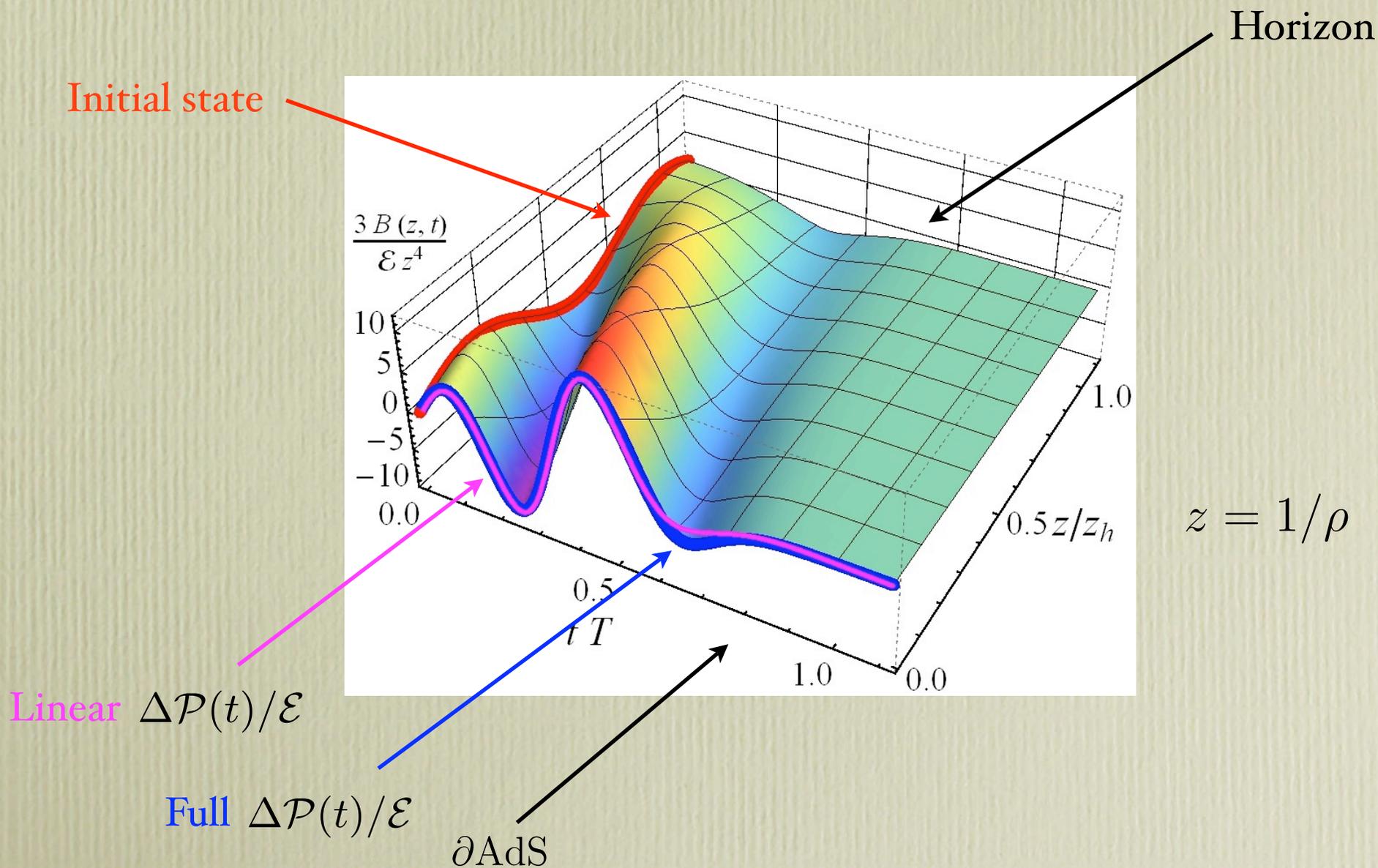
Price & Pullin '94

- BH collision in asymptotically flat 4D general relativity:

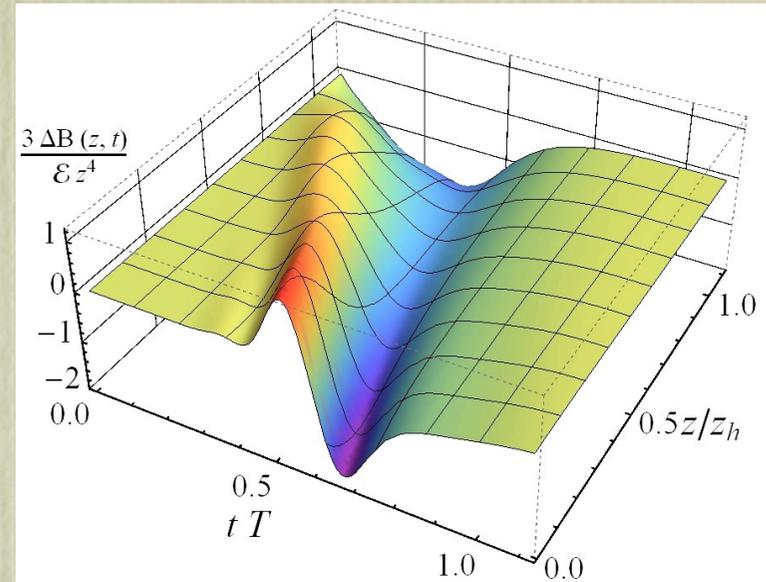
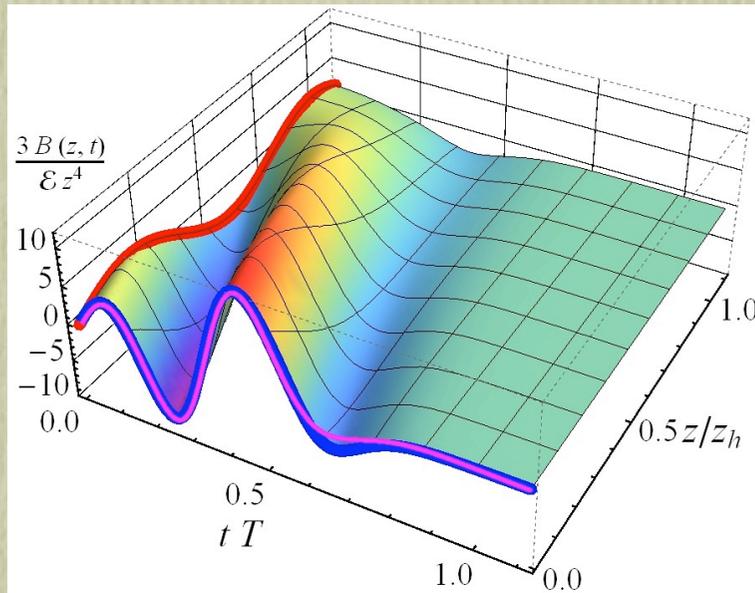


- Wave-form at infinity accurately reproduced (but perhaps non-asymptotic properties would not be).
- Analog in AdS: Boundary stress-tensor.

Results



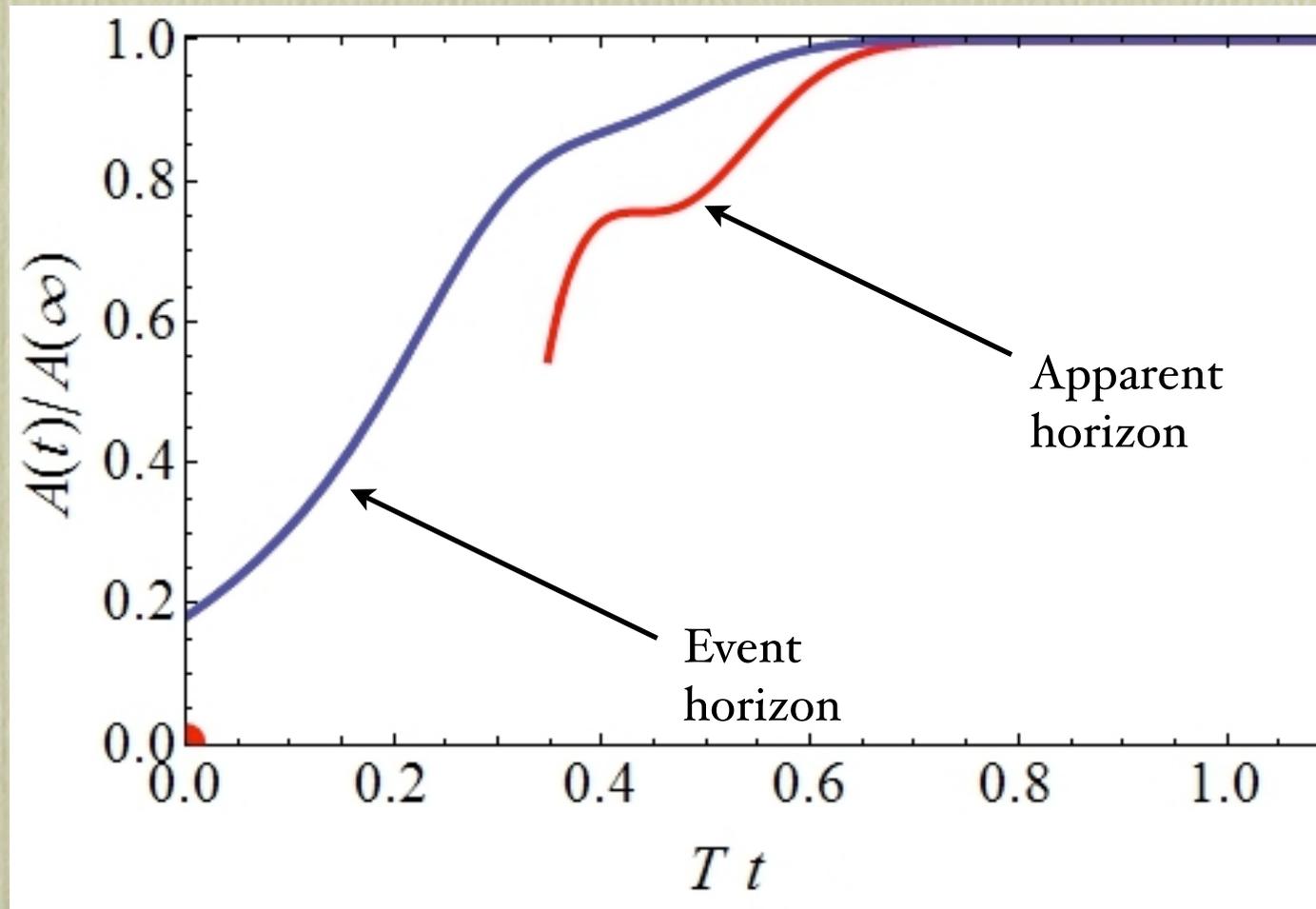
Results



- Over 2000 initial profiles.
- May or may not have AH at $t=0$.
- Ratio of scales gives accuracy: $2/10 \sim 20\%$
- $\Delta\mathcal{P}(t)/\mathcal{E} \sim 10$ implies far from equilibrium.

Results

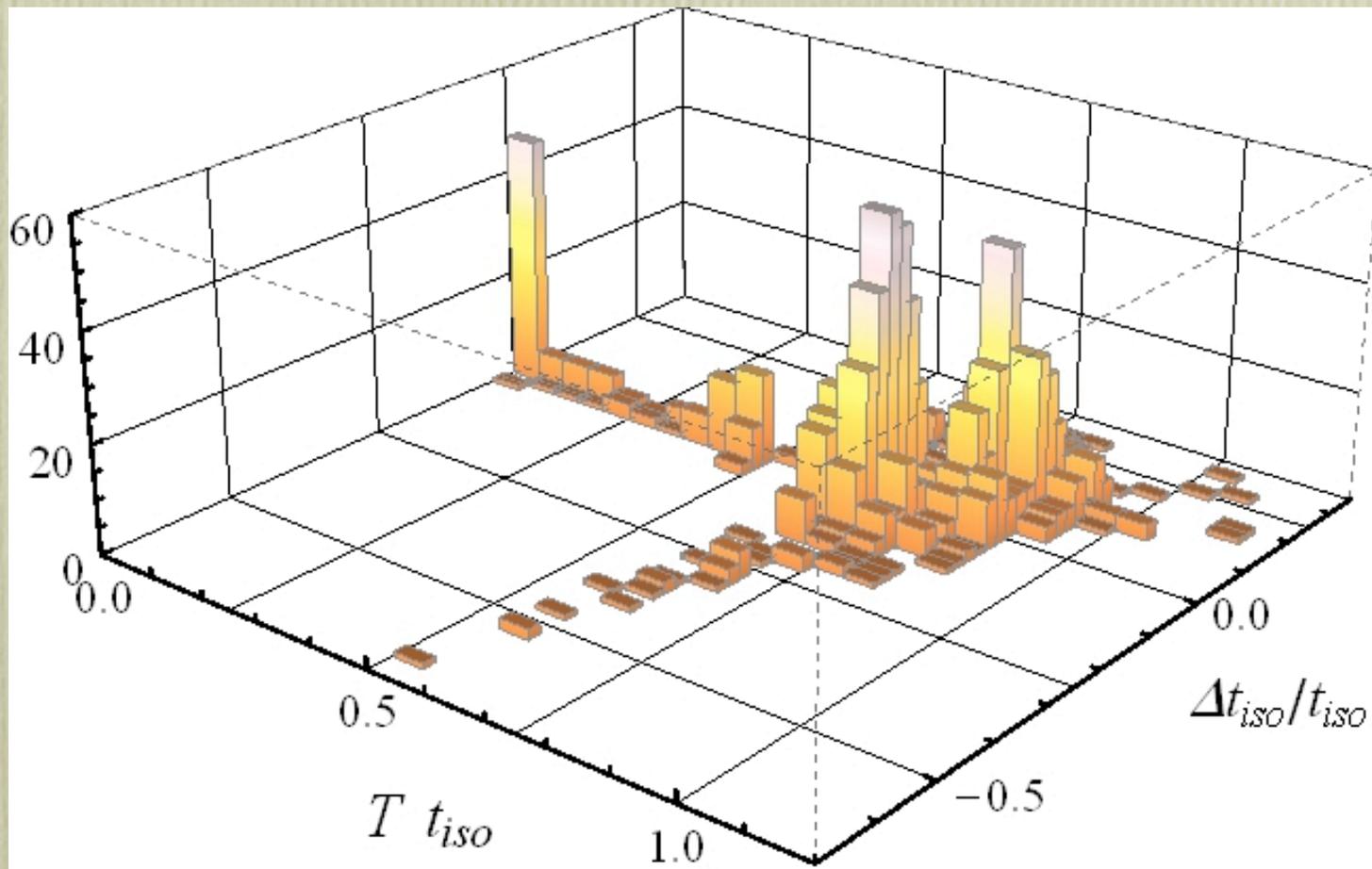
- “Entropy” increases during isotropization.



Results

- Isotropization time of order $1/T$ predicted by LA within 20%.

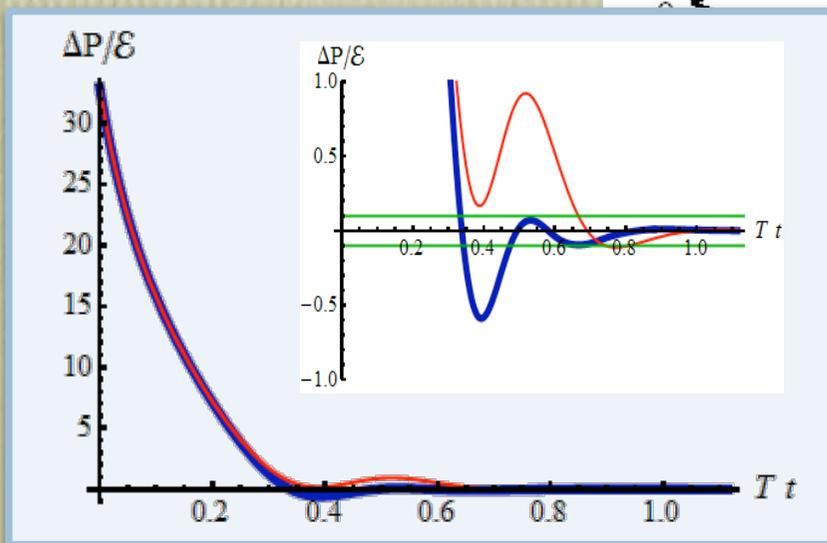
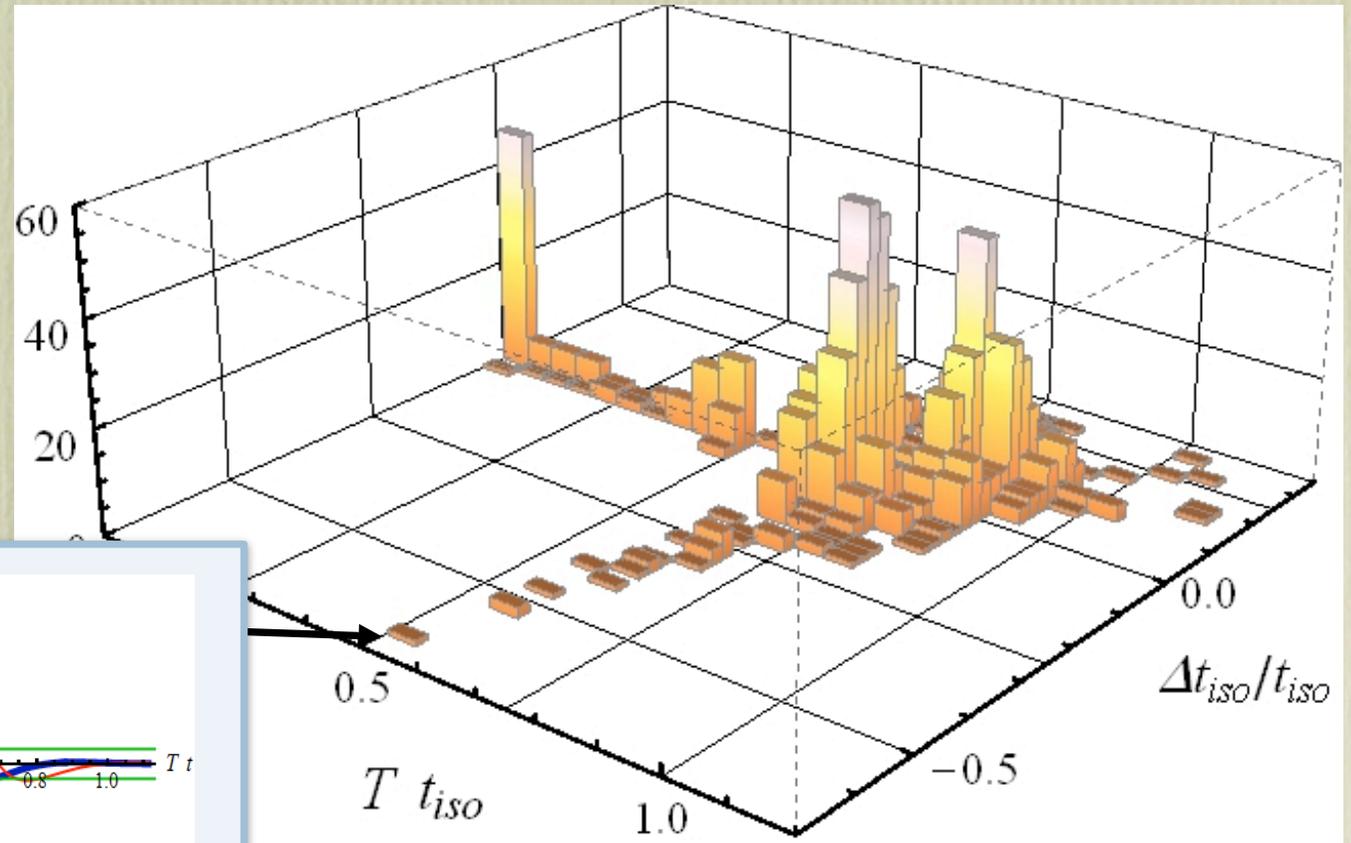
$$\Delta\mathcal{P}(t)/\mathcal{E} \leq 0.1$$



Results

- Isotropization time of order $1/T$ predicted by LA within 20%.

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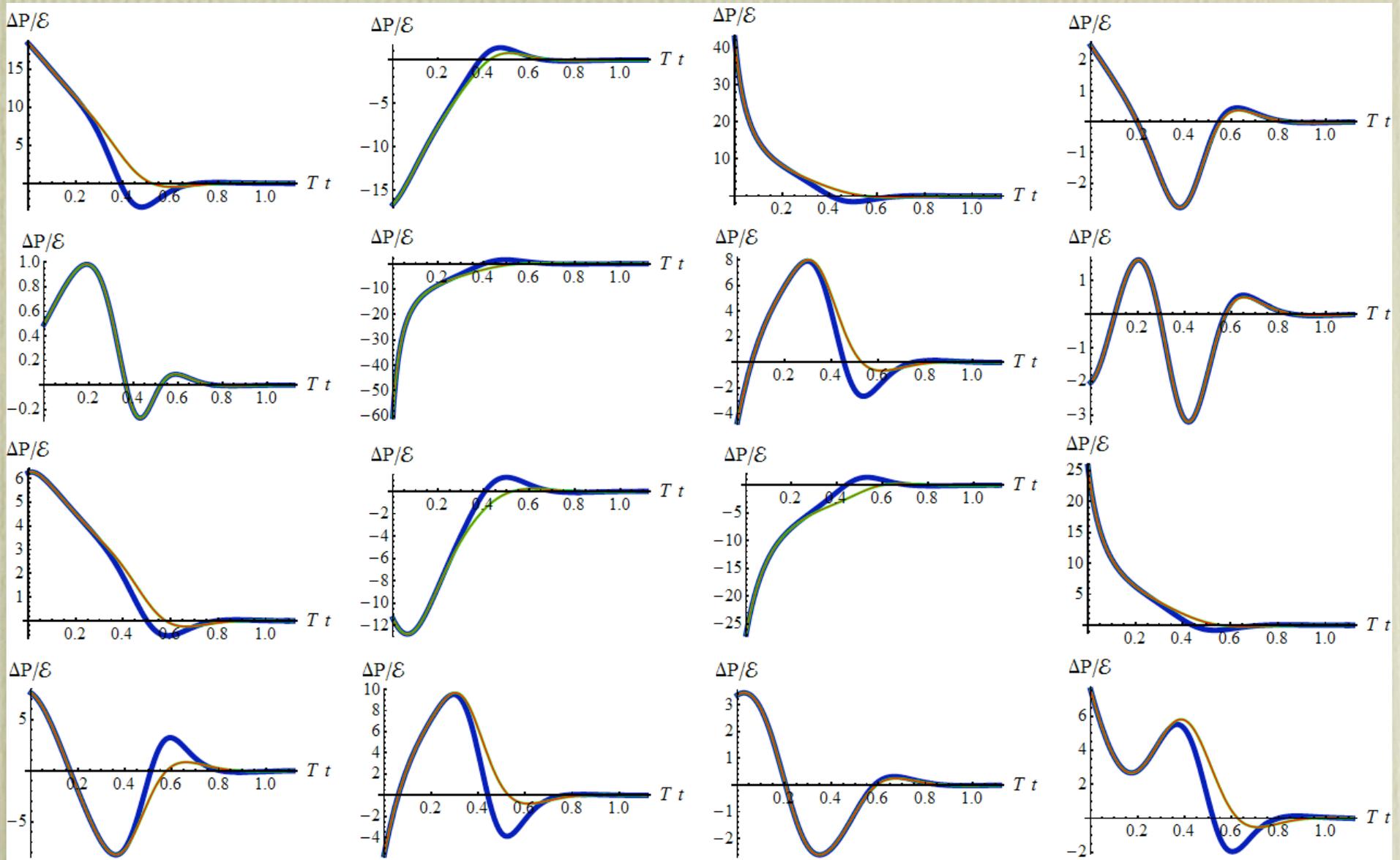


Expand in QNMs

$$\delta B(t, r) = \sum_i c_i b_i(r) e^{-i\omega_i t}$$

Expand in QNMs (Full / Linear / QNMs)

$$\delta B(t, r) = \sum_i c_i b_i(r) e^{-i\omega_i t}$$



Discussion

- Gauge: Small perturbations around equilibrium plasma
→ Linear response theory
- Gravity: Small perturbations around equilibrium black hole
→ Linearized Einstein's equations
- In both cases expect linear approx. if $\Delta\mathcal{P}/\mathcal{E} \ll 1$.
- LA in Fourier space:
 - HDMs: $\omega \rightarrow 0$ as $q \rightarrow 0$
 - QNMs: $\omega(0) \neq 0$
- We have studied far-from-equilibrium dynamics of QNMs.

Discussion

- For small perturbations:

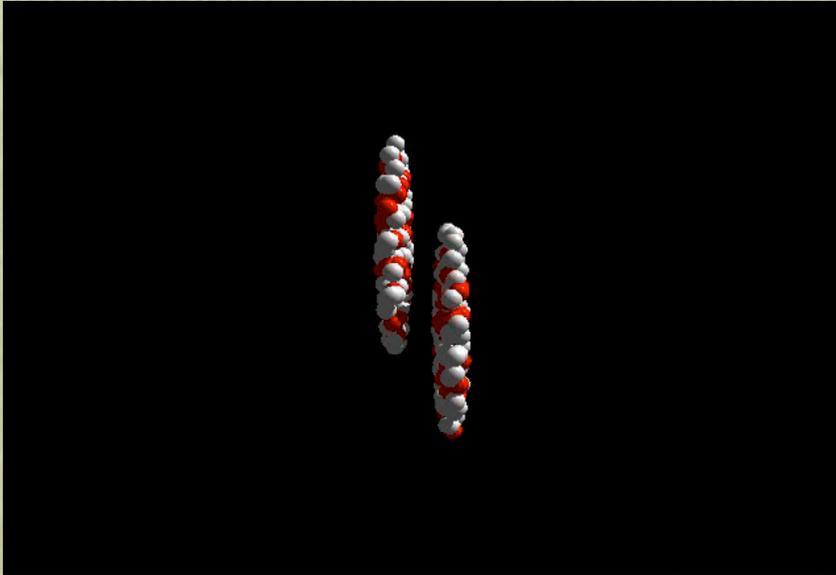
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.

- Extend to not-so-small perturbations by adding interactions.
- Expected to break down for $\Delta\mathcal{P}/\mathcal{E} \sim 1$ but it does not.
- Relaxation still characterized by few frequencies.
- Linear approx. valid for stress tensor 1-point function; other observables probably not well captured.

Discussion

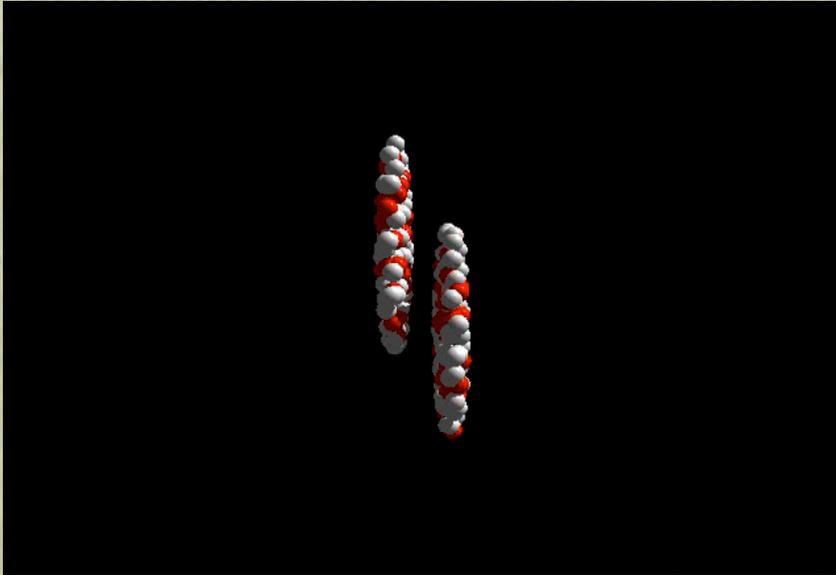
- Next step: Include hydrodynamics (boost-invariant case).
- Preliminary results indicate it works.

Potential implication



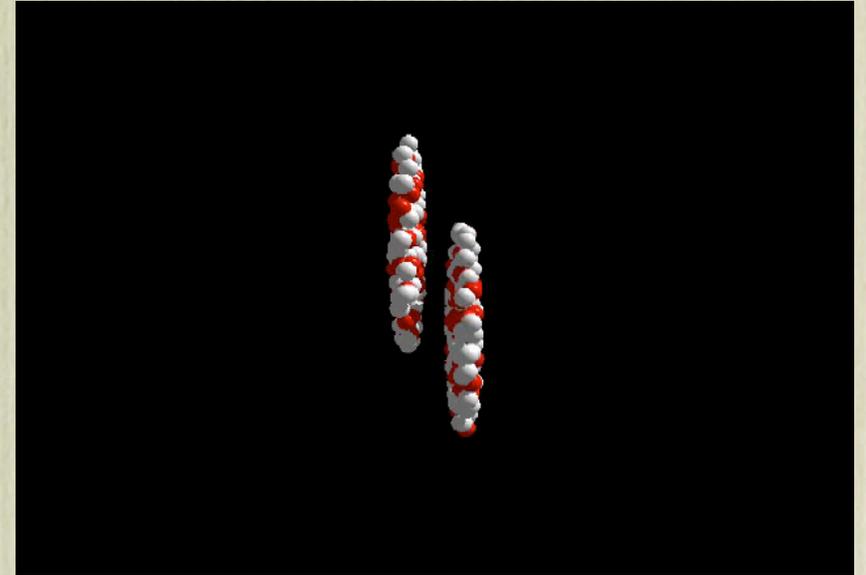
Initial state

Potential implication



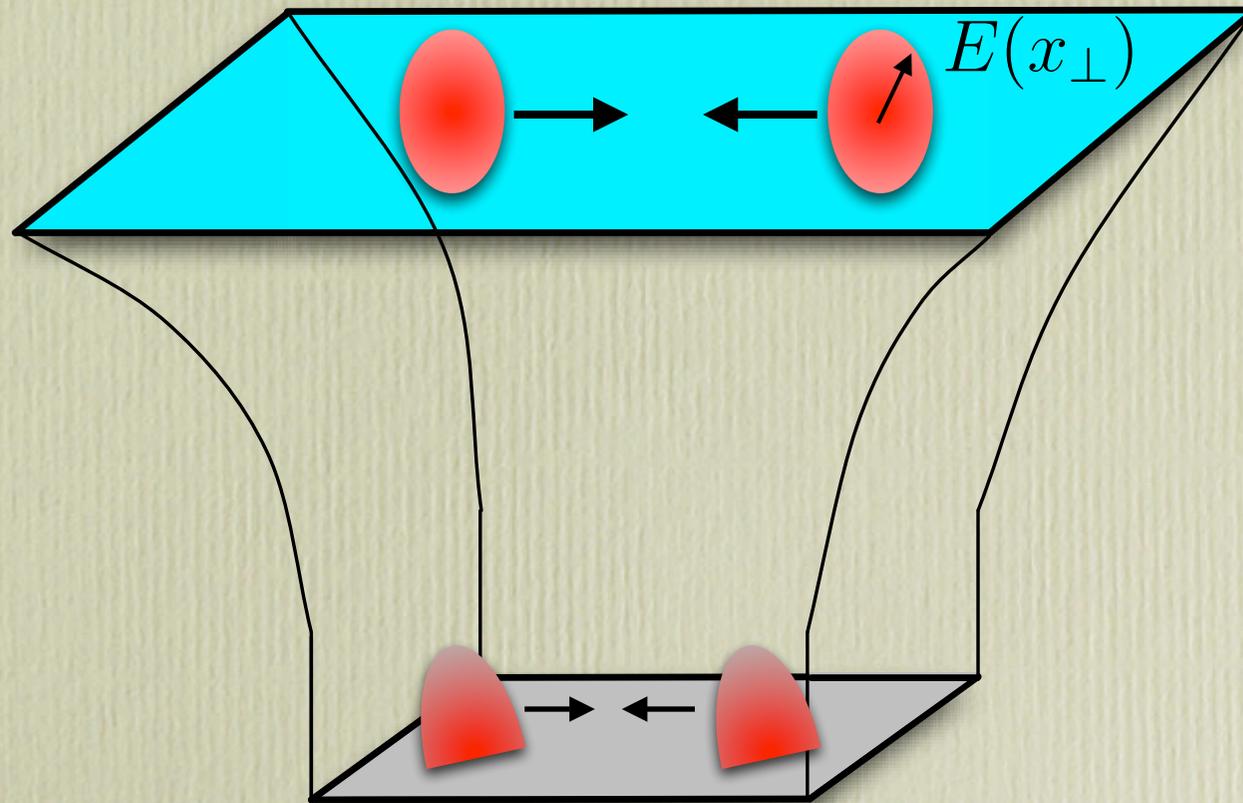
Initial state

Map?
→

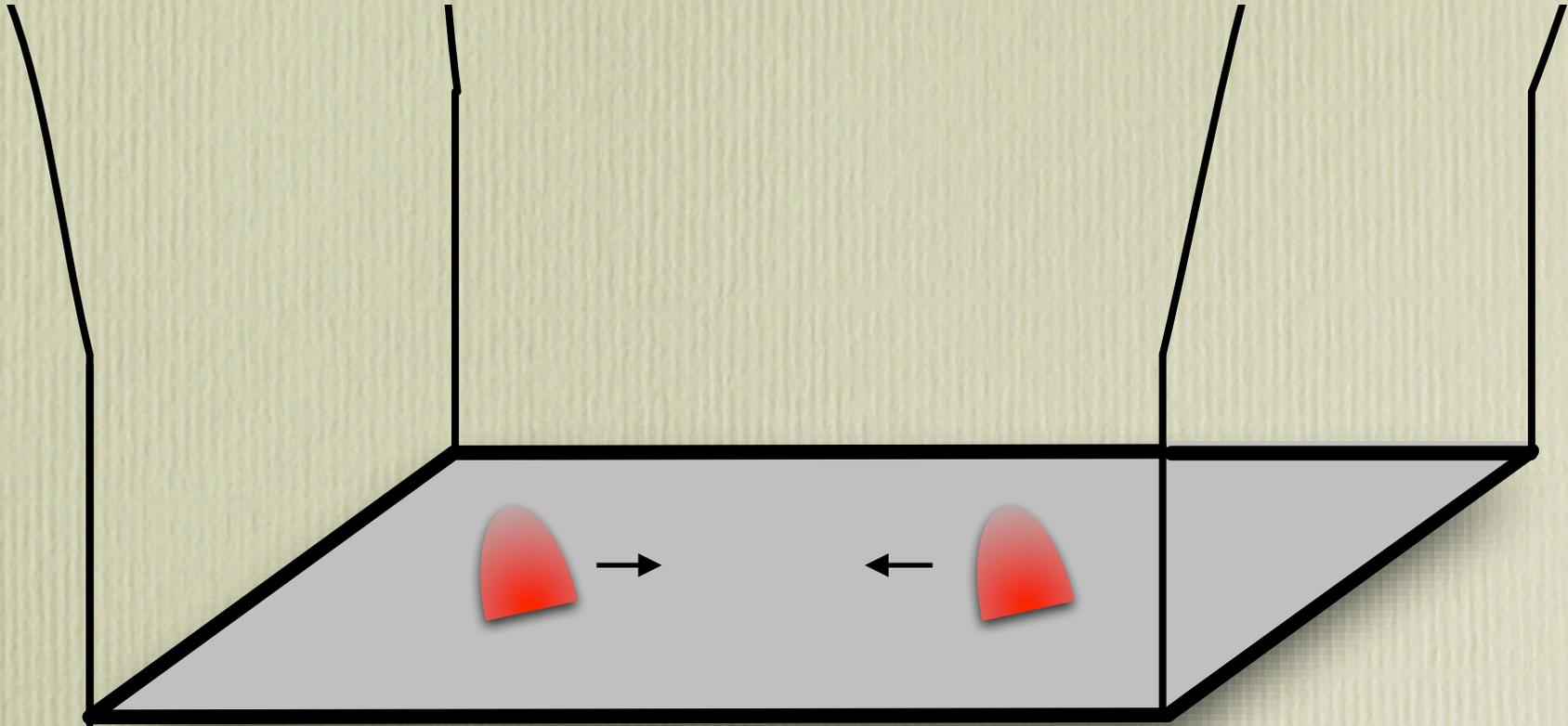


Hydro becomes applicable

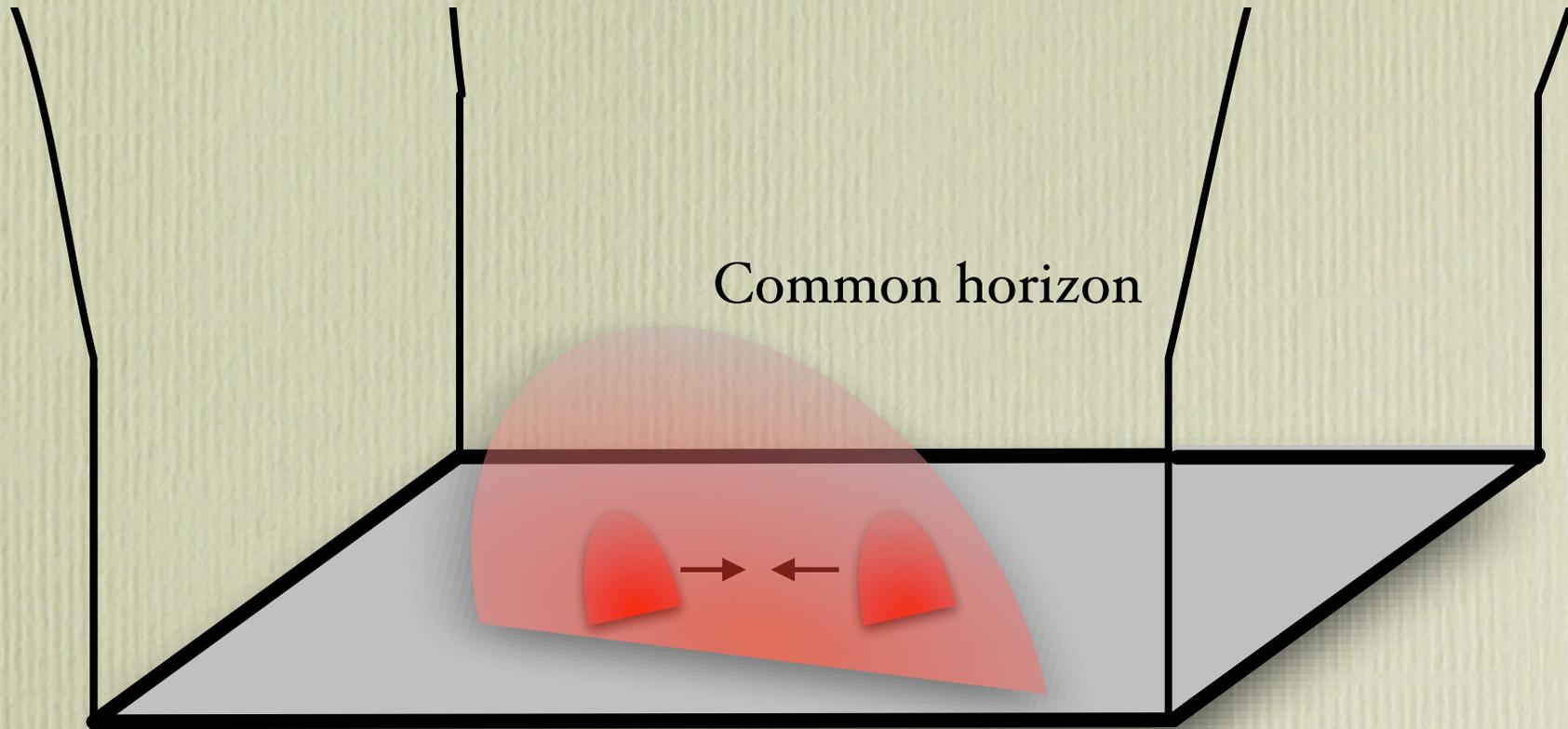
Potential implication



Potential implication



Potential implication



Full non-linearity of gravity encoded in the initial horizon.

Dedicated to Paul Townsend

- Hope it's clear how much respect I have for Paul both as a friend and as a teacher.

- Perhaps only one fault...

Dedicated to Paul Townsend

- you are too serious.



Thank you Paul, and Happy Birthday!