Describing M5-branes using Loop Spaces

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Based on:

- CS, arXiv:1007.3301, CMP ...
- C. Papageorgakis and CS, arXiv:1103.6192, JHEP ...
- Sam Palmer and CS arXiv:1105.????

Motivation

Find an algorithm for the construction of self-dual string solutions

- Effective description of M2-branes proposed in 2007.
- This created lots of interest:
  BLG-model: >467 citations, ABJM-model: >605 citations
- Inspired by an idea by Basu-Harvey:
  Propose a lift of the Nahm eqn. describing D1-D3-system:
  Basu-Harvey eqn. describes M2-M5-brane system
- Nahm transform:
  go from Nahm eqn. to Bogomolny monopole eqn.
  switch perspective from D1-brane to D3-brane
- Is there a lift for this Nahm transform?
  go from BH eqn. to self-dual string eqn.
  switch perspective from M2-brane to M5-brane
- Such a transform would open up interesting possibilities:
  eff. description of M5-branes, new integrable structures, . . .
- Alternative Approach: Gustavsson, 0802.3456

Outline

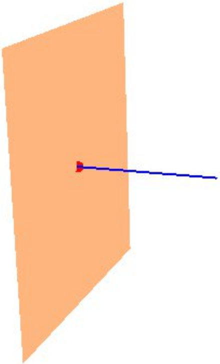
We will discuss the construction of monopoles and lift each ingredient to M-theory.

- Basu-Harvey lift of the Nahm equation and 3-Lie algebras
- Monopoles and self-dual strings
- Principal $\mathsf{U}(1)$-bundles, abelian gerbes and loop space
- ADHMN construction and its lift
- Examples of self-dual string solutions
- Non-abelian tensor multiplet on loop space

D1-D3-Branes and the Nahm Equation

D1-branes ending on D3-branes can be described by the Nahm equation.

dim	0	1	2	3	...	6
$D1$	$\times$					$\times$
$D3$	$\times$	$\times$	$\times$	$\times$		



N D1-branes ending on D3-branes:

A Monopole appears.

$X^i \in \mathfrak{u}(N)$: transverse fluctuations

Nahm equation: $(s = x^6)$

$$\frac{\mathrm{d}}{\mathrm{d}s} X^i + \varepsilon^{ijk}[X^j, X^k] = 0$$

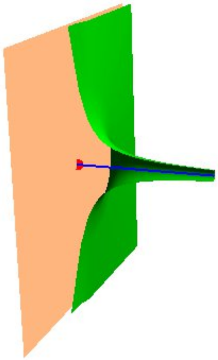
Note $\mathsf{SO}(3)$-invariance.

Solution: $X^i = r(s) G^i$ with

$$r(s) = \frac{1}{s}\ , \quad G^i = \varepsilon^{ijk}[G^j, G^k]$$

Nahm, Diaconescu, Tsimpis

dim	0	1	2	3	...	6
$D1$	$\times$					$\times$
$D3$	$\times$	$\times$	$\times$	$\times$		



Solution: $X^i = r(s)G^i$

$$r(s) = \frac{1}{s}\ , \quad G^i = \varepsilon^{ijk}[G^j, G^k]$$

The D1-branes form a fuzzy funnel:

G^i form irrep of $\mathfrak{su}(2)$:
coordinates on fuzzy sphere S^2_F

D1-worldvolume polarizes: $2d \to 4d$

Myers

IIB	0	1	2	3	4	5	6
$D1$	×						×
$D3$	×	×	×	×			

T-dualize along x^5:

IIA	0	1	2	3	4	5	6
$D2$	×					×	×
$D4$	×	×	×	×		×	

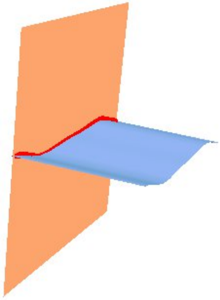
Interpret x^4 as M-theory direction:

M	0	1	2	3	4	5	6
$M2$	×					×	×
$M5$	×	×	×	×	×	×	

The Basu-Harvey lift of the Nahm Equation

M2-branes ending on M5-branes yield a Nahm equation with a cubic term.

M	0	1	2	3	4	5	6
$M2$	$\times$					$\times$	$\times$
$M5$	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	



Basu, Harvey, hep-th/0412310

A Self-Dual String appears.

Substitute $\mathsf{SO}(3)$-inv. Nahm eqn.

$$\frac{\mathrm{d}}{\mathrm{d}s}X^i + \varepsilon^{ijk}[X^j, X^k] = 0$$

by the $\mathsf{SO}(4)$-invariant equation

$$\frac{\mathrm{d}}{\mathrm{d}s}X^\mu + \varepsilon^{\mu\nu\rho\sigma}[X^\nu, X^\rho, X^\sigma] = 0$$

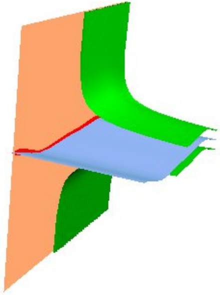
Solution: $X^\mu = r(s)G^\mu$ with

$$r(s) = \frac{1}{\sqrt{s}}\ , \quad G^\mu = \varepsilon^{\mu\nu\rho\sigma}[G^\nu, G^\rho, G^\sigma]$$

The Basu-Harvey lift of the Nahm Equation

M2-branes ending on M5-branes yield a Nahm equation with a cubic term.

M	0	1	2	3	4	5	6
$M2$	$\times$					$\times$	$\times$
$M5$	$\times$	$\times$	$\times$	$\times$	$\times$	$\times$	



Solution: $X^\mu = r(s)G^\mu$

$$r(s) = \frac{1}{\sqrt{s}}\ ,\quad G^\mu = \varepsilon^{\mu\nu\rho\sigma}[G^\nu, G^\rho, G^\sigma]$$

The M2-branes form a fuzzy funnel:

G^μ form a rep of $\mathfrak{so}(4)$:
coordinates on fuzzy sphere S^3_F

M2-worldvolume polarizes: $3d \to 6d$

What is this triple bracket?

What is the algebra behind the triple bracket?

In analogy with Lie algebras, we can introduce 3-Lie algebras.

$$\frac{\mathrm{d}}{\mathrm{d}s}X^\mu + [A_s, X^\mu] + \varepsilon^{\mu\nu\rho\sigma}[X^\nu, X^\rho, X^\sigma] = 0\ , \quad X^\mu \in \mathcal{A}$$

Trivial: $\mathcal{A}$ is a vector space, $[\cdot,\cdot,\cdot]$ trilinear+antisymmetric.

- ▷ Gauge transformations from inner derivations:

  Triple bracket forms a map $D : \mathcal{A}\wedge\mathcal{A} \to \mathrm{Der}(\mathcal{A}) =: \mathfrak{g}_\mathcal{A}$ via

$$D(A,B) \triangleright C := [A,B,C]$$

Demand a "3-Jacobi identity," the fundamental identity:

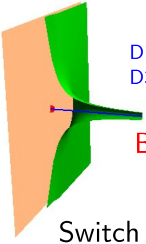
$$\begin{aligned} D(A,B) \triangleright (D(C,D) \triangleright E) &:= [A,B,[C,D,E]] \\ &= [[A,B,C],D,E] + [C,[A,B,D],E] + [C,D,[A,B,E]] \end{aligned}$$

The inner derivations form indeed a Lie algebra:

$$[D(A,B),D(C,D)] \triangleright E := D(A,B) \triangleright (D(C,D) \triangleright E) - D(C,D) \triangleright (D(A,B) \triangleright E)$$

Bracket closes due to fundamental identity.

Lifting monopoles to M-theory yields self-dual strings.



	0	1	2	3	4	5	6
D1	×						×
D3	×	×	×	×			

BPS configuration!

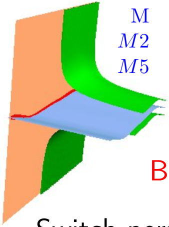
Switch perspective: D1→ D3:
Bogomolny monopole eqn.:

$$F_{ij} = \varepsilon_{ijk}\nabla_k\Phi \;\Rightarrow\; \nabla^2\Phi = 0$$

Single D3: Dirac monopole

$$\Phi = \frac{1}{r} \;\Rightarrow\; r(s) = \frac{1}{s}$$

⇒ matching profile!



M	0	1	2	3	4	5	6
$M2$	×					×	×
$M5$	×	×	×	×	×	×	

BPS configuration!

Switch perspective: M2→ M5:
Self-dual string eqn.:

$$H_{\mu\nu\rho} = \varepsilon_{\mu\nu\rho\sigma}\partial_\sigma\Phi \;\Rightarrow\; \partial^2\Phi = 0$$

Only single M5 known:

$$\Phi = \frac{1}{r^2} \;\Rightarrow\; r(s) = \frac{1}{\sqrt{s}}$$

⇒ matching profile!

Dirac Monopoles and Principal U(1)-bundles

Dirac monopoles are described by principal U(1)-bundles over S^2.

Manifold M with cover $(U_i)_i$. Principal U(1)-bundle over M:

$$F \in \Omega^2(M, \mathfrak{u}(1))\ ,$$
$$A_{(i)} \in \Omega^1(U_i, \mathfrak{u}(1)) \text{ with } F = \mathrm{d}A_{(i)}$$
$$g_{ij} \in \Omega^0(U_i \cap U_j, \mathsf{U}(1)) \text{ with } A_{(i)} - A_{(j)} = \mathrm{d}\log g_{ij}$$

Consider monopole in $\mathbb{R}^3$, but describe it on S^2 around monopole:

S^2 with patches U_+, U_-, $U_+ \cap U_- \sim S^1$: $g_{+-} = \mathrm{e}^{-\mathrm{i}n\phi},\ n \in \mathbb{Z}$

$$c_1 = \frac{\mathrm{i}}{2\pi}\int_{S^2} F = \frac{\mathrm{i}}{2\pi}\int_{S^1} A^+ - A^- = \frac{1}{2\pi}\int_0^{2\pi} n\mathrm{d}\phi = n$$

Monopole charge: n

Self-Dual Strings and Abelian Gerbes

Self-dual strings are described by abelian gerbes.

Manifold M with cover $(U_i)_i$. Abelian (local) gerbe over M:

$$H \in \Omega^3(M, \mathfrak{u}(1)) \ ,$$
$$B_{(i)} \in \Omega^2(U_i, \mathfrak{u}(1)) \text{ with } H = \mathrm{d}B_{(i)}$$
$$A_{(ij)} \in \Omega^1(U_i \cap U_j, \mathfrak{u}(1)) \text{ with } B_{(i)} - B_{(j)} = \mathrm{d}A_{ij}$$
$$h_{ijk} \in \Omega^0(U_i \cap U_j \cap U_k, \mathfrak{u}(1)) \text{ with } A_{(ij)} - A_{(ik)} + A_{(jk)} = \mathrm{d}h_{ijk}$$

Note: Local gerbe: principal $\mathsf{U}(1)$-bundles on intersections $U_i \cap U_j$.

Consider S^3, patches U_+, U_-, $U_+ \cap U_- \sim S^2$: bundle over S^2
Reflected in: $H^2(S^2, \mathbb{Z}) \cong H^3(S^3, \mathbb{Z}) \cong \mathbb{Z}$

$$\frac{\mathrm{i}}{2\pi}\int_{S^3} H = \frac{\mathrm{i}}{2\pi}\int_{S^2} B_+ - B_- = \ldots = n$$

Charge of self-dual string: n

Consider the following double fibration:

$$\begin{array}{ccc} & \mathcal{L}M \times S^1 & \\ {}^{ev}\swarrow & & \searrow^{pr} \\ M & & \mathcal{L}M \end{array}$$

Identify $T\mathcal{L}M = \mathcal{L}TM$, then: $x \in \mathcal{L}M \Rightarrow \dot{x}(\tau) \in T\mathcal{L}M$

Transgression

$$\mathcal{T} : \Omega^{k+1}(M) \to \Omega^k(\mathcal{L}M) \ , \quad \mathcal{T} = pr! \circ ev^*$$

$$(\mathcal{T}\omega)_x(v_1(\tau), \ldots, v_k(\tau)) := \int_{S^1} \mathrm{d}\tau \, \omega(v_1(\tau), \ldots, v_k(\tau), \dot{x}(\tau))$$

An abelian local gerbe over M is a principal $\mathsf{U}(1)$-bundle over $\mathcal{L}M$.

Note: We will work on $\mathcal{L}M \times S^1$, can do $\mathcal{L}M$, too.

The ADHMN construction

There is a map between and solutions to the Nahm equations.

Nahm transform: Instantons on $T^4 \mapsto$ instantons on $(T^4)^*$

Roughly here:

T^4: $\begin{cases} \text{3 rad. } 0 \\ \text{1 rad. } \infty \ : \ \text{D1 WV} \end{cases}$ and $(T^4)^*$: $\begin{cases} \text{3 rad. } \infty \ : \ \text{D3 WV} \\ \text{1 rad. } 0 \end{cases}$

Introduce (twisted) "Dirac operators":

$$\nabla\!\!\!/_{s,x} = -\mathbb{1}\tfrac{\mathrm{d}}{\mathrm{d}s} + \sigma^i \otimes (\mathrm{i}X^i + x^i \mathbb{1}_k) \ , \quad \bar{\nabla\!\!\!/}_{s,x} := \mathbb{1}\tfrac{\mathrm{d}}{\mathrm{d}s} + \sigma^i \otimes (\mathrm{i}X^i + x^i \mathbb{1}_k)$$

Properties:

$\Delta_{s,x} := \bar{\nabla\!\!\!/}_{s,x}\nabla\!\!\!/_{s,x} > 0$, $[\Delta_{s,x}, \sigma^i] = 0 \Leftrightarrow X^i$ satisfy Nahm eqn.

Normalized zero modes: $\bar{\nabla\!\!\!/}_{s,x}\psi_{s,x,\alpha} = 0$, $\mathbb{1} = \int_{\mathcal{I}} \mathrm{d}s\, \bar{\psi}_{s,x}\psi_{s,x}$ yield:

$$A_\mu := \int_{\mathcal{I}} \mathrm{d}s\, \bar{\psi}_{s,x} \frac{\partial}{\partial x^\mu} \psi_{s,x} \quad \text{and} \quad \Phi := -\mathrm{i} \int_{\mathcal{I}} \mathrm{d}s\, \bar{\psi}_{s,x}\, s\, \psi_{s,x}$$

This is a solution to the Bogomolny monopole equations!

Examples: Dirac monopoles

One can easily construct Dirac monopole solutions using the ADHMN construction.

Charge 1: Nahm eqn: $\partial_s X^i = 0$, so put $X^i = 0$. Zero mode:

$$\psi_+ = \mathrm{e}^{-sR}\frac{\sqrt{R+x^3}}{x^1 - \mathrm{i}x^2}\begin{pmatrix} x^1 - \mathrm{i}x^2 \\ R - x^3 \end{pmatrix}$$

Monopole solution:

$$\Phi^+ = -\frac{\mathrm{i}}{2R}\,,\;\; A_i^+ = \frac{\mathrm{i}}{2(x^1+x^2)^2}\left(x^2\left(1-\frac{x^3}{R}\right), -x^1\left(1-\frac{x^3}{R}\right), 0\right)$$

Charge 2: Nahm eqn. nontrivial. Choose:

$$X^i = -\frac{1}{s}T^i \quad \text{with} \quad T^i = \frac{\sigma^i}{2\mathrm{i}} = -\bar{T}^i$$

Resulting solution:

$$\Phi^+ = -\frac{\mathrm{i}}{R}\,,\;\; A_i^+ = \ldots$$

Lift of the "Dirac operator"

There is a natural lift of the Dirac operator to M-theory.

Type IIB (twisted):

$$\nabla^{\rm IIB}_{s,x} = -\mathbb{1}\frac{\rm d}{{\rm d}s} + \sigma^i({\rm i}X^i + x^i\mathbb{1}_k)$$

IIB	0	1	2	3	4	5	6
$D1$	×						×
$D3$	×	×	×	×			

Type IIA (twisted):

$$\nabla^{\rm IIA}_{s,x} = -\gamma_5\mathbb{1}_k\frac{\rm d}{{\rm d}s} + \gamma^4\gamma^i(X^i - {\rm i}x^i)$$

IIA	0	1	2	3	4	5	6
$D2$	×					×	×
$D4$	×	×	×	×		×	

M-theory (untwisted):

$$\nabla^{\rm M}_{s} = -\gamma_5\frac{\rm d}{{\rm d}s} + \tfrac{1}{2}\gamma^{\mu\nu}D(X^\mu, X^\nu)$$

M	0	1	2	3	4	5	6
$M2$	×					×	×
$M5$	×	×	×	×	×	×	

M-theory (twisted):

$$\nabla^{\rm M}_{s,x(\tau)} = -\gamma_5\frac{\rm d}{{\rm d}s} + \gamma^{\mu\nu}\left(\tfrac{1}{2}D(X^\mu, X^\nu) - {\rm i}x^\mu(\tau)\dot{x}^\nu(\tau)\right)$$

Lifted ADHMN Construction

The lifted ADHMN construction yields solutions to the loop space self-dual string eqns.

Recall: $\Delta^{\rm IIB} := \bar{\nabla\!\!\!/}^{\rm IIB}\nabla\!\!\!/^{\rm IIB}$, $[\Delta^{\rm IIB},\sigma^i]=0 \Leftrightarrow X^i$ satisfy Nahm eqn.

Here: $\Delta^{\rm M} := \bar{\nabla\!\!\!/}^{\rm M}\nabla\!\!\!/^{\rm M}$, $[\Delta,\gamma^{\mu\nu}]=0 \Leftarrow X^\mu$ satisfy BH eqn.

Our Dirac operator involved loop space, so we need to transgress:

$$H=\left(\varepsilon_{\mu\nu\rho\sigma}\frac{\partial}{\partial x^\sigma}\Phi\right)\mathrm{d}x^\mu\wedge\mathrm{d}x^\nu\wedge\mathrm{d}x^\rho$$

is turned into

$$F_{\mu\nu}(x(\tau)):=\frac{\partial}{\partial x^{[\mu}}A_{\nu]}(x(\tau))=\varepsilon_{\mu\nu\rho\sigma}\dot{x}^\rho(\tau)\frac{\partial}{\partial x^\sigma}\Phi(x(\tau))$$

From normalized, $\mathcal{A}$-valued zero modes $\psi_{s,x(\tau)}$ of $\bar{\nabla\!\!\!/}^{\rm M}$ construct

$$A_\mu=\int\mathrm{d}s\,\bar{\psi}_{s,x(\tau)}\frac{\partial}{\partial x^\mu}\psi_{s,x(\tau)}~,\quad \Phi=-\mathrm{i}\int\mathrm{d}s\,\bar{\psi}_{s,x(\tau)}\,s\,\psi_{s,x(\tau)}$$

Verification of the Construction

Verifying the construction is rather straightforward.

$$F_{\mu\nu} = \int \mathrm{d}s\, (\partial_{[\mu}\bar{\psi}_s)\partial_{\nu]}\psi_s$$

$$= \int \mathrm{d}s \int \mathrm{d}t\, (\partial_{[\mu}\bar{\psi}_s)\left(\psi_s\bar{\psi}_t - \nabla^{\mathrm{M}}_s G^{\mathrm{M}}(s,t)\bar{\nabla}^{\mathrm{M}}_t\right)\partial_{\nu]}\psi_t$$

$$= \int \mathrm{d}s \int \mathrm{d}t\, \bar{\psi}_s \left(\gamma^{\mu\kappa}\dot{x}^\kappa G^{\mathrm{M}}(s,t)\gamma^{\nu\lambda}\dot{x}^\lambda - \gamma^{\nu\kappa}\dot{x}^\kappa G^{\mathrm{M}}(s,t)\gamma^{\mu\lambda}\dot{x}^\lambda\right)\psi_t$$

Identity : $[\gamma^{\mu\kappa},\gamma^{\nu\lambda}]\dot{x}^\kappa\dot{x}^\lambda = -2\varepsilon_{\mu\nu\rho\sigma}\gamma^{\sigma\kappa}\gamma_5\dot{x}^\rho\dot{x}^\kappa$

$$F_{\mu\nu} = -\varepsilon_{\mu\nu\rho\sigma}\int \mathrm{d}s \int \mathrm{d}t\, \bar{\psi}_s \left(2\gamma^{\sigma\kappa}\gamma_5 G^{\mathrm{M}}(s,t)\dot{x}^\rho\dot{x}^\kappa\right)\psi_t$$

$$= -\mathrm{i}\varepsilon_{\mu\nu\rho\sigma}\dot{x}^\rho \int \mathrm{d}s \int \mathrm{d}t\, \Big((\partial_\sigma\bar{\psi}_s)\left(\psi_s\bar{\psi}_t - \nabla^{\mathrm{M}}_s G^{\mathrm{M}}(s,t)\bar{\nabla}^{\mathrm{M}}_t\right) t\,\psi_t +$$
$$\bar{\psi}_s\, s\, \left(\psi_s\bar{\psi}_t - \nabla^{\mathrm{M}}_s G^{\mathrm{M}}(s,t)\bar{\nabla}^{\mathrm{M}}_t\right)\partial_\sigma\psi_t\Big)$$

$$= -\mathrm{i}\varepsilon_{\mu\nu\rho\sigma}\dot{x}^\rho \int \mathrm{d}s\, (\partial_\sigma\bar{\psi}_s)\, s\,\psi_s + \bar{\psi}_s\, s\,\partial_\sigma\psi_s$$

Reduction to the ADHMN Construction

The lift reduces in the expected way to the ADHMN construction.

On $\mathcal{L}S^3 \subset \mathcal{L}\mathbb{R}^4$: $x^\mu x^\mu = \dot{x}^\mu \dot{x}^\mu = R^2$, $x^\mu \dot{x}^\mu = 0$.

Reduction (cf. Mukhi/Papageorgakis, 0803.3218):

$$\langle X^4 \rangle = \frac{r}{\ell_p^{3/2}} e_4 = g_{\mathrm{YM}} e_4 \ , \quad \dot{x}^4(\tau_0) = R \Rightarrow \dot{x}^i(\tau_0) = x^4(\tau_0) = 0$$

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma} \dot{x}^\rho \frac{\partial}{\partial x^\sigma} \Phi_{\mathrm{SDS}} \quad \to \quad F_{ij} = \varepsilon_{ijk} \frac{\partial}{\partial x^k} R \Phi_{\mathrm{SDS}} + \dots$$

$$\frac{\mathrm{d}}{\mathrm{d}s} X^\mu = \tfrac{1}{3!} \varepsilon^{\mu\nu\rho\sigma} [X^\nu, X^\rho, X^\sigma] \quad \to \quad \frac{\mathrm{d}}{\mathrm{d}s} X^i = \tfrac{1}{2} \varepsilon^{ijk} R [X^j, X^k] + \dots$$

$$\begin{aligned} \nabla\!\!\!/^{\mathrm{M}} &= -\gamma_5 \frac{\mathrm{d}}{\mathrm{d}s} + \gamma^{\mu\nu} \left(\tfrac{1}{2} D^{(\rho)}(X^\mu, X^\nu) - \mathrm{i} x^\mu(\tau) \dot{x}^\nu(\tau) \right) \\ &\to -\gamma_5 \frac{\mathrm{d}}{\mathrm{d}s} + \gamma^{\mu\nu} \left(\tfrac{1}{2} D^{(\rho)}(X^\mu, X^\nu) - \mathrm{i} x^\mu(\tau_0) \dot{x}^\nu(\tau_0) \right) \\ &= -\gamma_5 \frac{\mathrm{d}}{\mathrm{d}s} + R \gamma^{4i} \left(X^{i\alpha} D^{(\rho)}(e_\alpha, e_4) - \mathrm{i} x^i(\tau_0) \right) + \dots = \nabla\!\!\!/^{\mathrm{IIA}} + \dots \end{aligned}$$

Charge 1: Choose again trivial Nahm data. Zero modes:

$$\psi \sim \mathrm{e}^{-R^2 s} \begin{pmatrix} \mathrm{i}\left(R^2 + x^2\dot{x}^1 - x^1\dot{x}^2 - x^4\dot{x}^3 + x^3\dot{x}^4\right) \\ x^3(\dot{x}^1 + \mathrm{i}\dot{x}^2) + x^4(\dot{x}^2 - \mathrm{i}\dot{x}^1) - (x^1 + \mathrm{i}x^2)(\dot{x}^3 - \mathrm{i}\dot{x}^4) \\ 0 \\ 0 \end{pmatrix}$$

Solution:

$$\Phi = \frac{\mathrm{i}}{2R^2}\ , \quad F = \frac{2\mathrm{i}\sin\theta^1\sin^2\theta^2(\dot{\theta}^2\,\mathrm{d}\phi\wedge\mathrm{d}\theta^1 - \dot{\theta}^1\,\mathrm{d}\phi\wedge\mathrm{d}\theta^2 + \dot{\phi}\,\mathrm{d}\theta^1\wedge\mathrm{d}\theta^2)}{\sqrt{\dot{\phi}^2 + 2(\dot{\theta}^1)^2 + 4(\dot{\theta}^2)^2 - (\dot{\phi}^2 + 2(\dot{\theta}^1)^2)\cos(2\theta^2) - 2\dot{\phi}^2\cos(2\theta^1)\sin^2\theta^2}}$$

This solves the loop-space self-dual string equation.

Regression:

$$\begin{aligned} H &= F|_{\dot{\theta}^1=1,\dot{\theta}^2=0,\dot{\phi}=0} \wedge \sin\theta^2\mathrm{d}\theta^1 - F|_{\dot{\theta}^1=0,\dot{\theta}^2=1,\dot{\phi}=0} \wedge \mathrm{d}\theta^2 \\ &\quad + F|_{\dot{\theta}^1=0,\dot{\theta}^2=0,\dot{\phi}=1} \wedge \sin\theta^1\sin\theta^2\mathrm{d}\phi \\ &= 6\mathrm{i}\sin\theta^1\sin^2\theta^2\,\mathrm{d}\theta^1\wedge\mathrm{d}\theta^2\wedge\mathrm{d}\phi\ , \end{aligned}$$

This is indeed the expected solution.

Charge 2:

Nahm data:

$$X^\mu = \frac{e_\mu}{\sqrt{2s}} ~, \quad e_\mu \text{ generate } \mathcal{A}$$

Solution:

$$\Phi(x) = \frac{\mathrm{i}}{R^2}$$

As expected: twice the charge of the case $k = 1$.

Remarks

Our lift of the ADHMN construction is very natural and rather straightforward.

- The lift of the Dirac operator was natural considering the corresponding brane configurations.
- It is natural to go to loop space to describe self-dual strings.
- The construction nicely involves the Basu-Harvey equation.
- It reduces nicely to the ADHMN construction.
- The construction does produce transgressed self-dual strings.
- A regression can be performed to get original self-dual string.
- It can be trivially rendered nonabelian:

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma}\dot{x}^\rho \nabla^\sigma X^6$$

The non-abelian tensor multiplet

A recently proposed 3-Lie algebra valued tensor-multiplet implies a transgression.

Recall the transgression map:

$$(\mathcal{T}\omega)_x(v_1(\tau),\dots,v_k(\tau)) := \int_{S^1} \mathrm{d}\tau\, \omega(v_1(\tau),\dots,v_k(\tau),\dot{x}(\tau))$$

Equations found by Lambert, Papageorgakis, 1007.2982:

$$\nabla^2 X^I - \tfrac{\mathrm{i}}{2}[\bar{\Psi},\Gamma_\nu\Gamma^I\Psi,C^\nu] - [X^J,C^\nu,[X^J,C_\nu,X^I]] = 0$$

$$\Gamma^\mu\nabla_\mu\Psi - [X^I,C^\nu,\Gamma_\nu\Gamma^I\Psi] = 0$$

$$\nabla_{[\mu}H_{\nu\lambda\rho]} + \tfrac{1}{4}\varepsilon_{\mu\nu\lambda\rho\sigma\tau}[X^I,\nabla^\tau X^I,C^\sigma] + \tfrac{\mathrm{i}}{8}\varepsilon_{\mu\nu\lambda\rho\sigma\tau}[\bar{\Psi},\Gamma^\tau\Psi,C^\sigma] = 0$$

$$F_{\mu\nu} - D(C^\lambda,H_{\mu\nu\lambda}) = 0$$

$$\nabla_\mu C^\nu = D(C^\mu,C^\nu) = 0$$

$$D(C^\rho,\nabla_\rho X^I) = D(C^\rho,\nabla_\rho\Psi) = D(C^\rho,\nabla_\rho H_{\mu\nu\lambda}) = 0$$

Factorization of $C^\rho = C\dot{x}^\rho$. Here, 3-Lie algebra transgression:

$$(\mathcal{T}\omega)_x(v_1(\tau),\dots,v_k(\tau)) := \int_{S^1} \mathrm{d}\tau\, D(\omega(v_1(\tau),\dots,v_k(\tau),\dot{x}(\tau)),C)$$

The non-abelian tensor multiplet on loop space

The corresponding equations can all be rewritten on loop space.

Transgression of matter fields (not in Huang, Huang, 1008.3834)

$$\mathring{X}^I(x(\tau)) := R\, D(C, X^I(x(\tau)))\,, \quad \mathring{\Psi}(x(\tau)) := \Gamma^\rho \dot{x}_\rho D(C, \Psi(x(\tau)))$$

Equations of motion (SYM-like):

$$\nabla^2 \mathring{X}^I + \tfrac{\mathrm{i}}{2}\tfrac{1}{R}\dot{x}^\nu[\bar{\mathring{\Psi}}, \Gamma_\nu \Gamma^I \mathring{\Psi}] + [\mathring{X}^J, [\mathring{X}^J, \mathring{X}^I]] = 0\,,$$
$$\tfrac{1}{R}\Gamma^{\mu\nu}\dot{x}_\nu \nabla_\mu \mathring{\Psi} - \Gamma^I[\mathring{X}^I, \mathring{\Psi}] = 0\,,$$
$$\nabla_\mu \mathring{F}^{\mu\nu} + [\mathring{X}^I, \nabla^\nu \mathring{X}^I] + \tfrac{\mathrm{i}}{2}\left([\bar{\mathring{\Psi}}, \Gamma^\nu \mathring{\Psi}] - \tfrac{2}{R^2}\dot{x}^\sigma \dot{x}^\nu [\bar{\mathring{\Psi}}, \Gamma_\sigma \mathring{\Psi}]\right) = 0\,,$$

Supersymmetry transformations (SYM-like):

$$\delta \mathring{X}^I = \tfrac{1}{R}\mathrm{i}\bar{\varepsilon}\Gamma^I \dot{x}^\rho \Gamma_\rho \mathring{\Psi}\,,$$
$$\delta \mathring{A}_\mu = \tfrac{1}{R^2}\mathrm{i}\bar{\varepsilon}\Gamma_{\mu\lambda}\Gamma_\rho \dot{x}^\lambda \dot{x}^\rho \mathring{\Psi}\,,$$
$$\delta \mathring{\Psi} = \tfrac{1}{R}\Gamma^{\nu\mu}\dot{x}_\nu \Gamma^I \nabla_\mu \mathring{X}^I \varepsilon + \tfrac{1}{2}\Gamma_{\mu\nu}\mathring{F}^{\mu\nu}\varepsilon - \tfrac{1}{2}\Gamma^{IJ}[\mathring{X}^I, \mathring{X}^J]\varepsilon\,,$$
$$\delta C^\mu = 0$$

Remarks

The loop space tensor multiplet fits well into the picture.

- Get SYM theory on loop space from the tensor multiplet
- C-field blocks modes of the theory, need to get rid of it
- Our loop space self-dual string equation extends compatibly:

$$F_{\mu\nu} = \varepsilon_{\mu\nu\rho\sigma}\dot{x}^\rho \nabla^\sigma X^6$$

- ADHMN construction for two M5-branes using this equation
- extension to real and hermitian 3-algebra (ABJM model)
- Right direction, more work in progress to get rid of $\dot{x}$ etc.

Summary:

- ✓ Reformulation of self-dual string equation on loop space
- ✓ Generalized ADHMN construction for self-dual string
- ✓ Explicit constructions in many cases
- ✓ Reformulate non-abelian tensor multiplet eqns. on loop space
- ✓ Partially generalized ADHMN construction

Future directions:

- ▷ Find (2,0)-theory
- ▷ Study classical integrability in more detail
- ▷ Quantization of S^3 via gerbes and groupoids

Describing M5-branes using Loop Spaces

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