

**Atomic data for non-thermal electron  
excitation of coronal ions:  
application to FeXIII**

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## Introduction:

In solar corona modellings, it is usually assumed free electrons to be Maxwellian, corresponding to some temperature  $T_e$ . To simulate dynamic phenomenae, a refinement is to introduce a **high energy tail** (or a Kappa-distribution) : indeed, fast electrons are more difficult to thermalize than low energy electrons... Unfortunately, the effect of high-energy electrons is often smeared by the low energy electron contribution, because cross-section decreases with energy. But, we saw yesterday, that for a line ratio done between 2 lines issued from 2 levels of different excitation energy, the high energy tail can play an important role... Non-thermal ? but what's about the free electron **angular** distribution: **isotropic** or **anisotropic** ?

Electron distribution is assumed **axial symmetric** around the quantization axis ( $z$ ) used to define atomic sublevels  $|\alpha J M_J\rangle$ . It can be developed on Legendre polynomials:  $(P_0(\cos\Theta) = 1, P_2(\cos\Theta) = \frac{1}{2}(3(\cos\Theta)^2 - 1), \dots)$

$$g(v, \Theta) = \sum_{K=0}^{\infty} g_K(v) P_K(\cos\Theta) \quad (1)$$

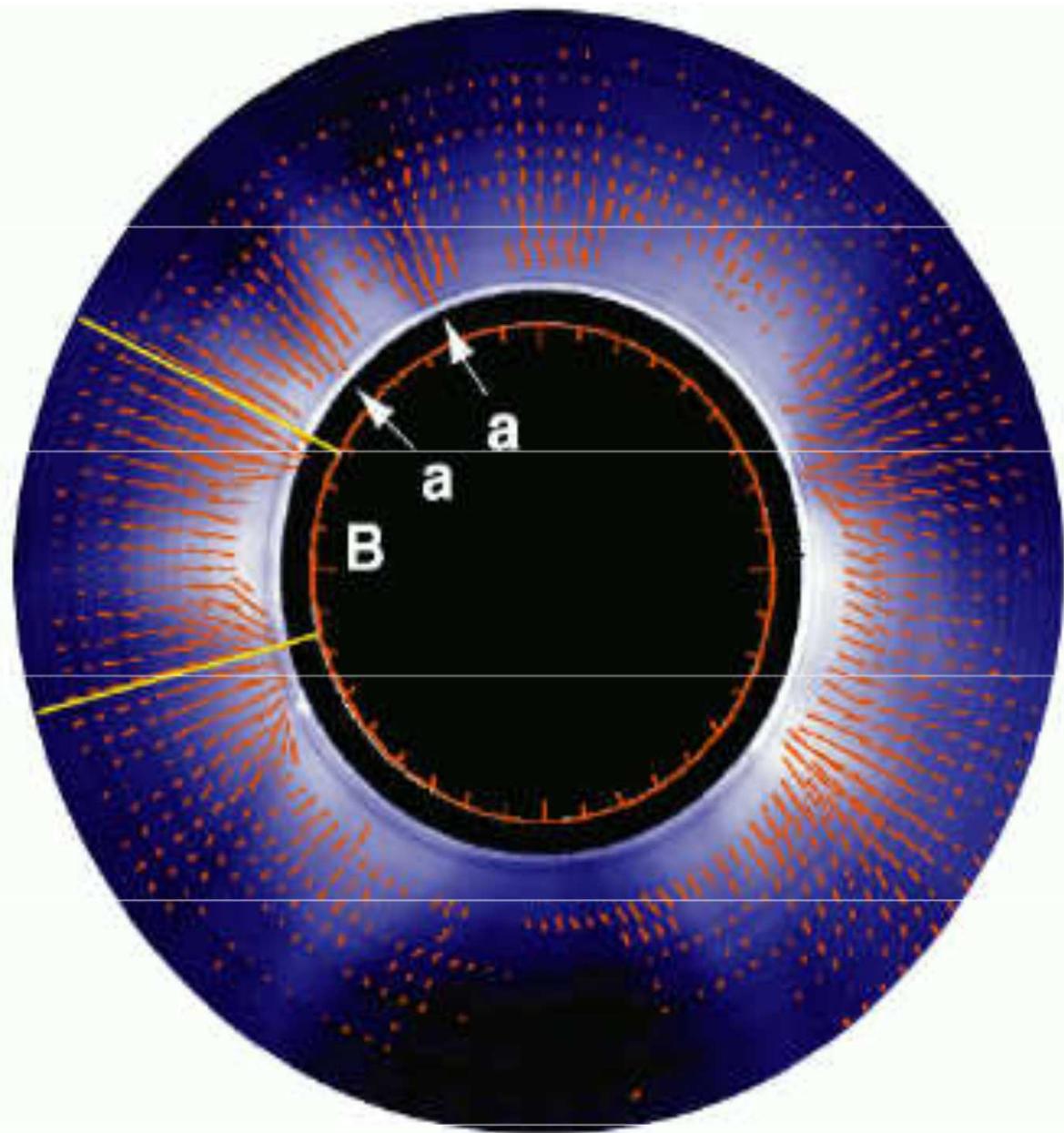
where :  $\int_0^{\infty} dv \int_0^{\pi} d\Theta g(v, \Theta) v^2 \sin\Theta = 1$  (2)

$$g_K(v) = \frac{2K+1}{2} \int_0^{\pi} P_K(\cos\Theta) g(v, \Theta) \sin\Theta d\Theta \quad (3)$$

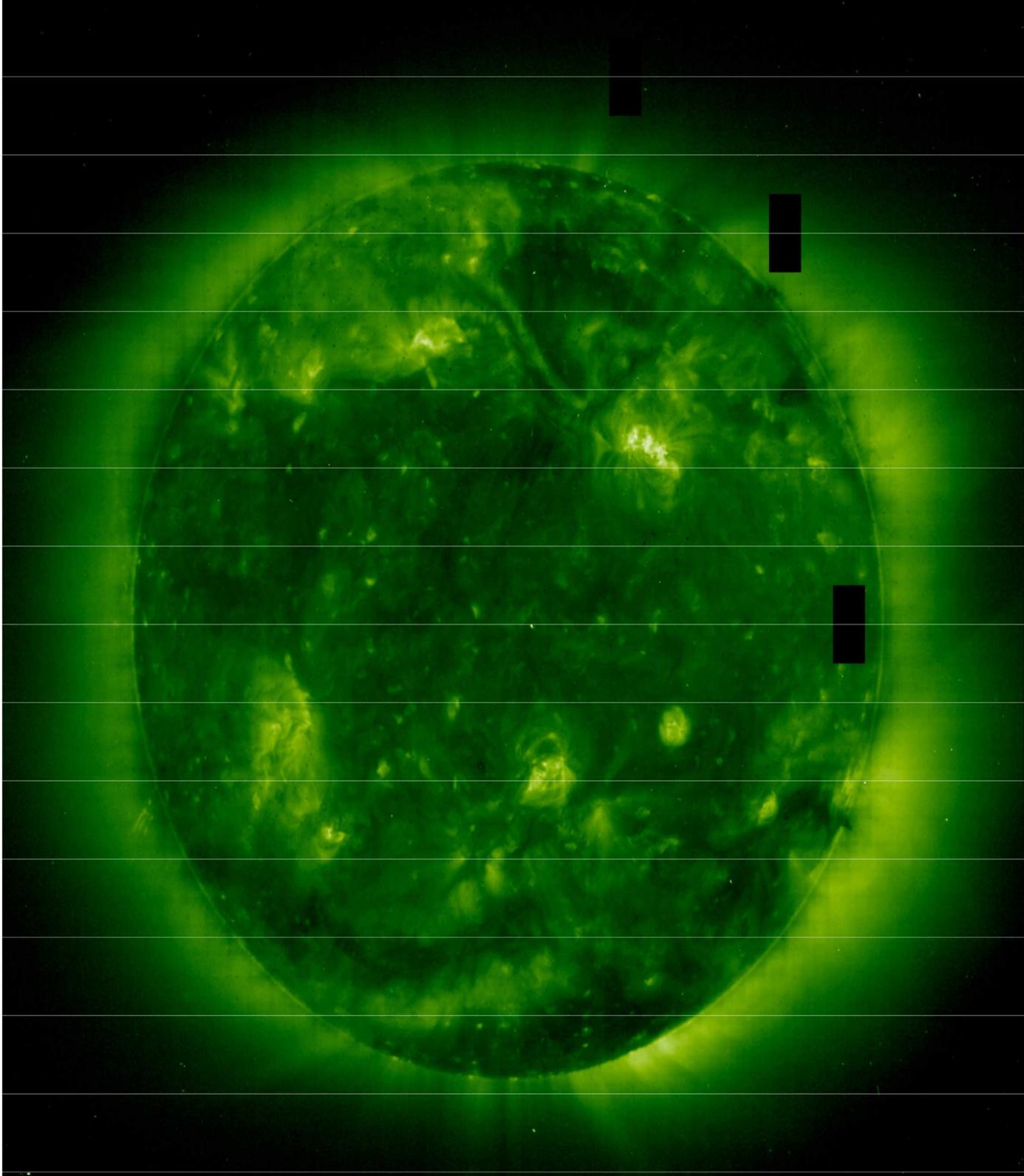
e.g.  $g_K(v) = \frac{1}{2}(2K+1) \delta(v - v_0)$  for **an electron beam**  
 $(v = v_0, \Theta = 0) \Rightarrow g(v, \Theta) = \delta(v - v_0) \delta(\Theta) / \sin\Theta$

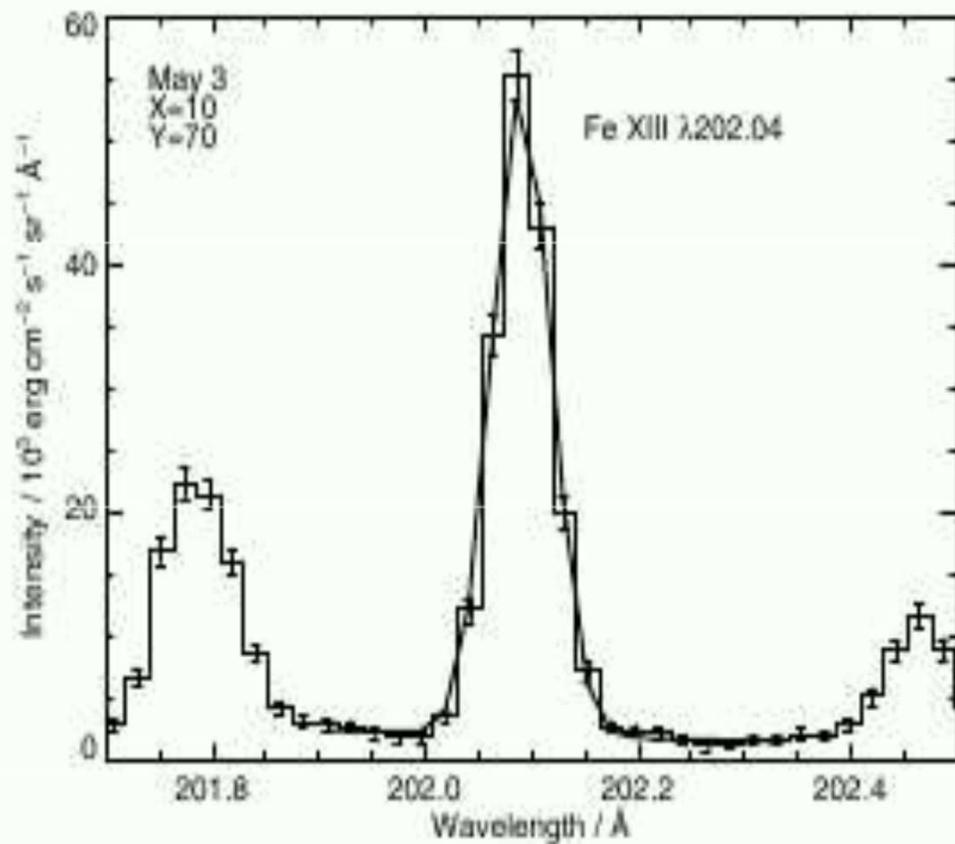
Table I. **FeXIII** energy levels from NIST Tables ( $1s^2 2s^2 2p^6$ )

Index		$\alpha$	$J$	Energy (Ryd = 13.606 eV)
1	$3s^2 3p^2$	$^3P$	0	0.0000
2	$3s^2 3p^2$	$^3P$	1	0.0848
3	$3s^2 3p^2$	$^3P$	2	0.1691
4	$3s^2 3p^2$	$^1D$	2	0.4380
5	$3s^2 3p^2$	$^1S$	0	0.8339
20	$3s^2 3p 3d$	$^3P$	1	4.5102
24	$3s^2 3p 3d$	$^3D$	2	4.6406
25	$3s^2 3p 3d$	$^3D$	3	4.6400
26	$3s^2 3p 3d$	$^1F$	3	5.0745

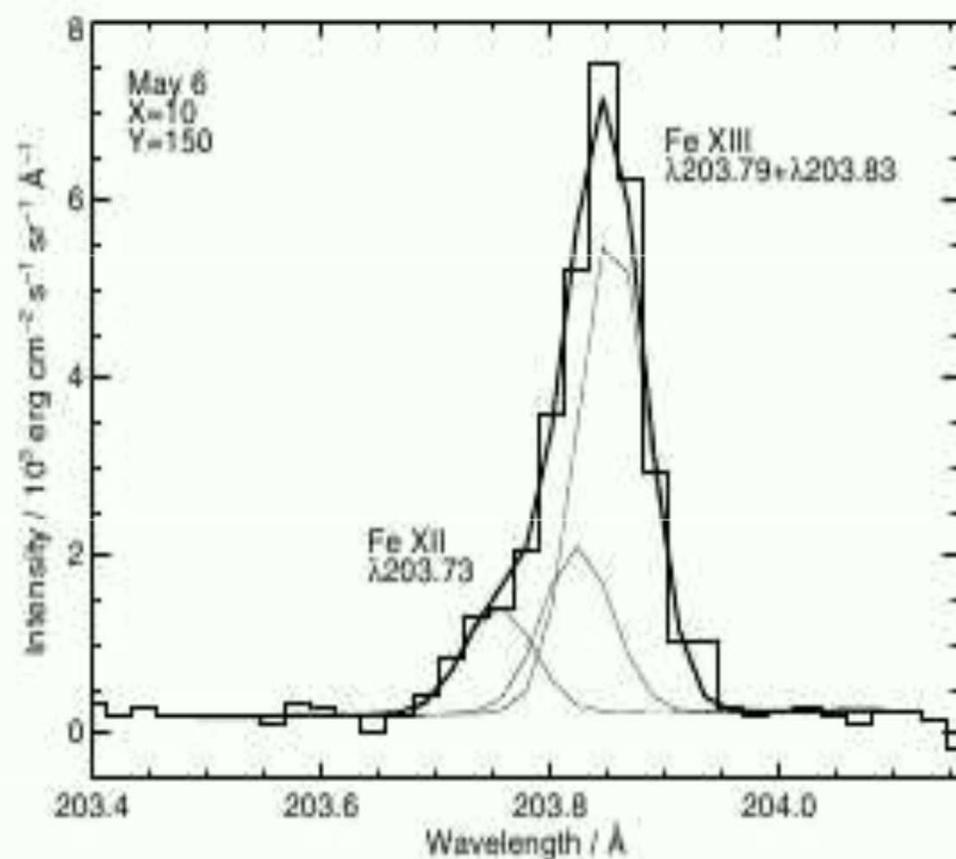


**Fig. 3.** February 23, 1980 KERP polarization map of the 1074.7 nm FeXIII line superimposed on a Mauna Loa K-corona image obtained the same day (from Habbal *et al.* 2001).





**Fig. 8.** An example single Gaussian fit to the Fe XIII  $\lambda$ 202.0 emission line. The May 3 data set is used, and the image pixel indicated.



**Fig. 9.** An example fit to the feature at 203.8  $\text{\AA}$  that comprises Fe XII  $\lambda 203.72$ , and Fe XIII  $\lambda \lambda 203.79, 203.83$  from the May 6 data set. The spatial pixel chosen is indicated. Three thin lines are used to show the three individual Gaussians from the fit.

**Intensity ratio** as function of  $\Theta$  ( $\alpha_i J_i \rightarrow \alpha_f J_f$ ) :

$$I(\text{anisotropic: } \eta \neq 0) / I(\text{isotropic : } \eta = 0)$$

$$-0.5 \leq \eta \leq +0.25$$

$$\frac{I(\eta, \Theta)}{I(\eta = 0)} = 1 + (3 \cos^2 \Theta - 1)\eta, \quad \eta = W(\alpha_i J_i \alpha_f J_f) \frac{\rho_0^2(\alpha_i J_i)}{\rho_0^0(\alpha_i J_i)}$$

$$\text{where } W(\alpha_i J_i \alpha_f J_f) = \frac{1}{2\sqrt{2}} \frac{\begin{Bmatrix} J_i & J_i & 2 \\ 1 & 1 & J_f \end{Bmatrix}}{\begin{Bmatrix} J_i & J_i & 0 \\ 1 & 1 & J_f \end{Bmatrix}} \quad (\text{M1, E1})$$

- **Level population**  $N(\alpha_i J_i) = \rho_0^0(\alpha_i J_i) \sqrt{2J_i + 1}$ ,
  - **Level alignment**  $\rho_0^2(\alpha_i J_i)$
- (see Sahal-Bréchet S. 1977, Ap. J, 213, 887)

$$J_i = 0 : \frac{\rho_0^2(\alpha_i J_i)}{\rho_0^0(\alpha_i J_i)} = 0;$$

$$J_i = 1 : \frac{\rho_0^2(\alpha_i J_i)}{\rho_0^0(\alpha_i J_i)} = \frac{\sqrt{2} (N_1 - N_0)}{N(\alpha_i J_i)};$$

$$J_i = 2 : \frac{\rho_0^2(\alpha_i J_i)}{\rho_0^0(\alpha_i J_i)} = \sqrt{\frac{10}{7}} \frac{(2N_2 - N_1 - N_0)}{N(\alpha_i J_i)}; \quad \text{etc.}$$

( $N_i$  is the population of the sublevel  $|M_{J_i}|$ )

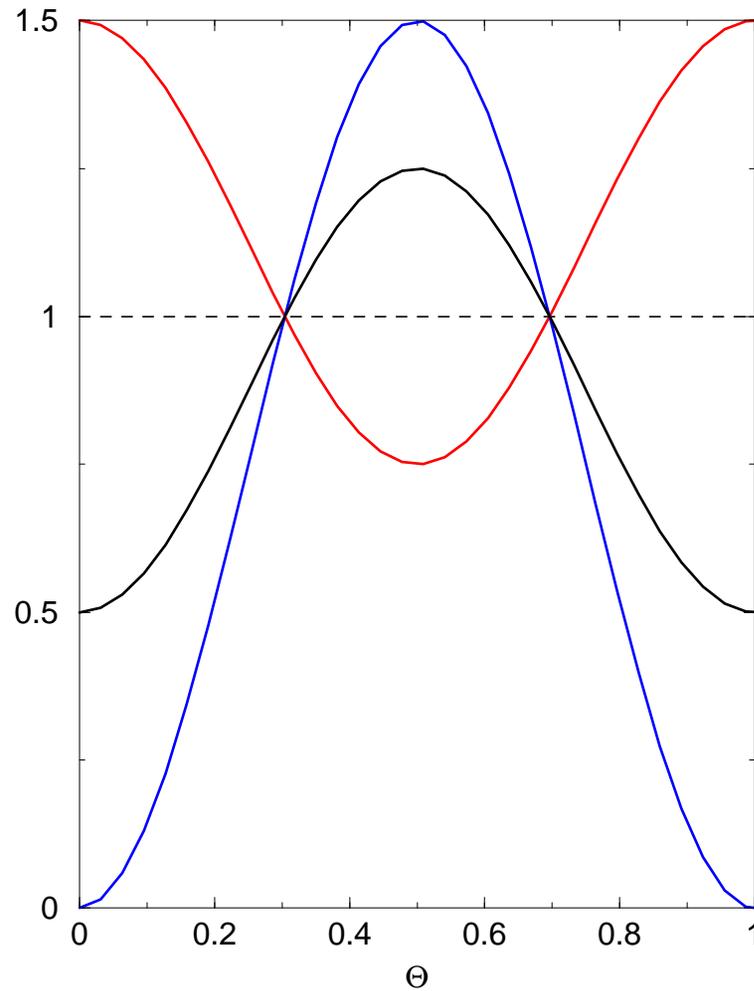


Figure 1:  $I(\eta, \Theta)/I(\eta = 0)$ :  $\eta = -0.50$ ,  $\eta = -0.25$ ,  $\eta = +0.25$

Table 2. Fe XIII wavelengths and  $W(\alpha_i J_i \alpha_f J_f)$   
“algebraic” coefficients

Wavelength	(Å)	i - f	$J_i - J_f$	$W(\alpha_i J_i \alpha_f J_f)$
10747	(IR)	2 - 1	1 - 0	+ 0.3536
10798	(IR)	3 - 2	2 - 1	+ 0.2092
196.540	(UV)	26 - 4	3 - 2	+ 0.1732
202.044	(UV)	20 - 1	1 - 0	+ 0.3536
203.797	(UV)	24 - 3	2 - 2	- 0.2092
203.828	(UV)	25 - 3	3 - 2	+ 0.1732

Table 3. Percentage occupation numbers for FeXIII  
 (Young, Mason and Thomas, Proc. 3rd SOHO Workshop)

Level	$10^8$	$10^9$	$10^{10}$	$10^{11}$	( $\text{cm}^{-3}$ )
$(3s^2 3p^2) \ ^3P_0$	88	41	14	8.9	
$\ ^3P_1$	6.9	27	30	25	
$\ ^3P_2$	5.5	30	45	38	
$\ ^1D_2$	0.07	1.4	11	23	
$\ ^1S_0$	0.001	0.01	0.2	1.5	

Table 4.1 Collision strength for  $K = 2$  angular distribution,  
 $E = 15$  Ryd (204 eV),  $\rho_0^2(i) \leftarrow \rho_0^0(j)$

i	j	$\Omega_{2,0}^2$	i	j	$\Omega_{2,0}^2$
2	1	- 6.250-4	25	1	- 3.822-5
3	1	+1.783-3	25	2	- 1.305-4
3	2	+1.142-3	25	3	- 3.273-2
20	1	- 6.604-2	26	1	- 1.374-4
20	2	- 2.108-4	26	2	- 2.719-4
20	3	- 1.759-3	26	3	- 1.329-3
24	1	- 5.549-5			
24	2	- 1.563-2			
24	3	+1.863-2			

Collision strengths, for  $K = 0$  angular distribution,  $\Omega_{0,0}^0$ ,  $\rho_0^0(\alpha_i J_i) \leftarrow \rho_0^0(\alpha_j J_j)$ . They are related to usual collision strengths  $\Omega_{ij}$ , which are symmetric, i.e.  $\Omega_{ij} = \Omega_{ji}$ :

$$\Omega_{0,0}^0(i j) = \sqrt{\frac{2J_i + 1}{2J_j + 1}} \Omega_{ij}$$

Assuming a coronal model (excitation  $j \rightarrow i$ , radiation  $i \rightarrow f$ ) and an electron beam distribution, we have:

$$\frac{\rho_0^2(i)}{\rho_0^0(i)} = \frac{g_2(v) \Omega_{2,0}^2(i j)}{g_0(v) \Omega_{0,0}^0(i j)} = \frac{5 \Omega_{2,0}^2(i j)}{\Omega_{0,0}^0(i j)}$$

Table 5.1 Collision strengths  $\Omega_{ji}$ , compared to AK [3], FM[4]  
 $E=15$  Ryd. ( $a-b \equiv a \times 10^{-b}$ )

j	i	present	AK	FM	j	i	present	AK	FM
1	2	2.054-2	2.099-2	2.05-2	1	25	6.989-3	8.073-3	6.80-3
1	3	9.933-2	9.349-2	9.80-2	2	25	2.180-2	2.589-2	2.18-2
2	3	2.259-1	2.178-1	2.08-1	3	25	<b>6.404-0</b>	<b>6.063-0</b>	<b>8.48-0</b>
1	20	<b>1.710-0</b>	<b>1.622-0</b>	<b>2.71-0</b>	1	26	3.534-3	3.915-3	3.60-3
2	20	6.592-3	9.454-3	1.37-1	2	26	1.187-2	1.311-2	1.25-2
3	20	4.470-1	3.888-1	2.68-1	3	26	1.807-1	2.170-1	3.38-1
1	24	1.920-3	2.599-3	1.40-3					
2	24	<b>1.484-0</b>	<b>1.379-0</b>	<b>1.51-0</b>					
3	24	<b>2.076-0</b>	<b>1.949-0</b>	<b>3.25-0</b>					

Table 6.  $\frac{\rho_0^2(ij)}{\rho_0^2(ij)}$  and  $\eta(i f)$  for  $E = 15, 30, 45$  Ryd.

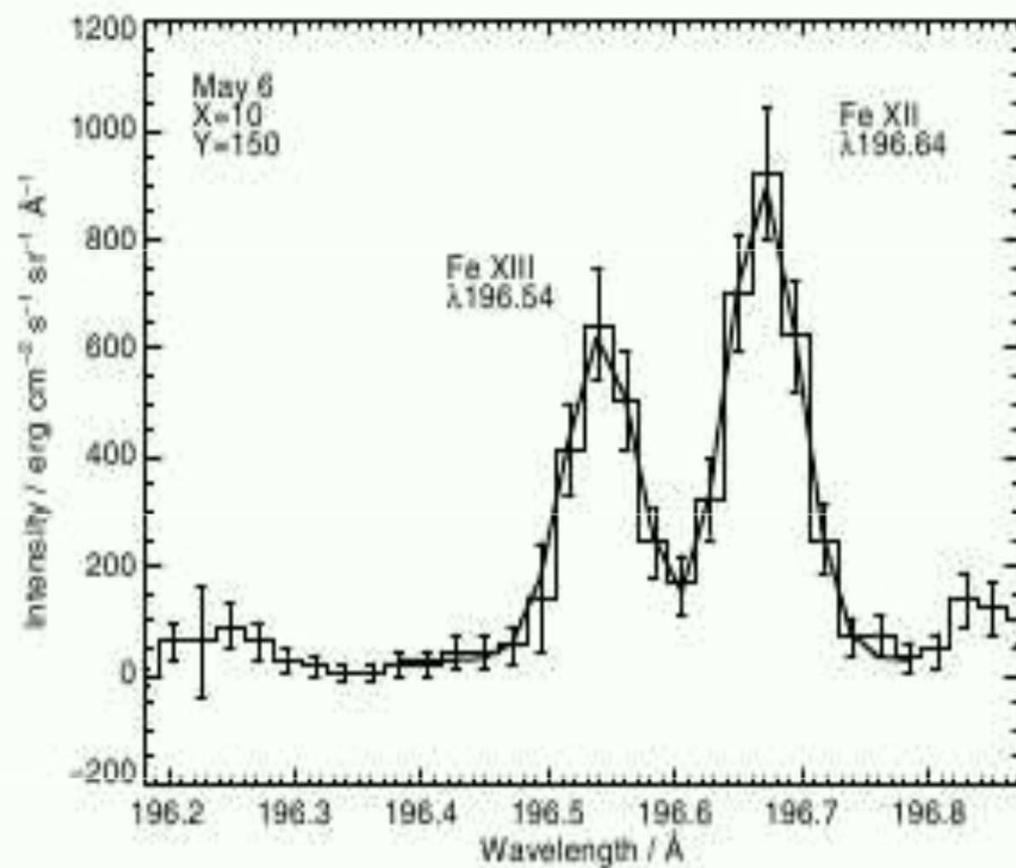
j	i	f	$\frac{\rho_0^2}{\rho_0^2}(15)$	$\eta(15)$	$\frac{\rho_0^2}{\rho_0^2}(30)$	$\eta(30)$	$\frac{\rho_0^2}{\rho_0^2}(45)$	$\eta(45)$
1	20	1	- 0.115	- 0.041	- 0.060	- 0.021	- 0.022	- 0.008
2	24	3	- 0.040	+ 0.008	- 0.022	+ 0.005	- 0.008	+ 0.002
3	24	3	+ 0.045	- 0.009	+ 0.028	- 0.006	+ 0.015	- 0.003
3	25	3	- 0.022	- 0.004	- 0.010	- 0.002	- 0.001	- 0.000

For the first line ( $j = f = 1, i = 20$ ), at  $E = 5.5$  Ryd

$$\frac{\rho_0^2}{\rho_0^2}(5.5) = -0.131 \text{ and } \eta(5.5) = -0.046$$

## References

- [1] Arnaud, J. 2003, EAS Publication Series 9, 209
- [2] Young, P.R., Watanabe, T., Hara, H. & Mariska, J.T. 2009, A & A 495, 587
- [3] Aggarwal, K.M. & Keenan F.P. 2004, A & A 418, 371
- [4] Fawcett, B. C. & Mason, H.E. 1989, ADNDT 43, 245



**Fig. 7.** An example double-Gaussian fit to Fe XIII  $\lambda$ 196.54 and Fe XII  $\lambda$ 196.64 from the May 6 data set. The chosen spatial pixel is indicated. The data are plotted as a thin line and the fit with a thick line. The error bars on the intensity measurements are also shown.

Table 4.2 Collision strength for  $K = 2$  angular distribution,  
 $E = 30$  Ryd (408 eV)

i	j	$\Omega_{2,0}^2$	i	j	$\Omega_{2,0}^2$
2	1	-3.418-4	25	1	+7.900-5
3	1	+1.724-3	25	2	+4.449-5
3	2	+1.034-3	25	3	-1.584-2
20	1	-3.813-2	26	1	-4.617-5
20	2	-7.234-5	26	2	-9.533-5
20	3	-1.467-3	26	3	-7.224-4
24	1	-1.815-5			
24	2	-8.986-3			
24	3	+1.242-2			

Table 4.3 Collision strength for  $K = 2$  angular distribution,  
 $E = 45$  Ryd (612 eV)

i	j	$\Omega_{2,0}^2$	i	j	$\Omega_{2,0}^2$
2	2	-2.103-4	25	1	+1.403-4
3	1	+1.725-3	25	2	+1.253-4
3	2	+9.964-4	25	2	<b>-1.267-3</b>
20	1	<b>-1.358-2</b>	26	1	-1.627-5
20	2	-2.674-5	26	2	-3.866-5
20	3	-1.226-3	26	3	-2.951-4
24	1	-5.981-6			
24	2	<b>-3.084-3</b>			
24	3	<b>+6.827-3</b>			

Table 5.2 Collision strengths  $\Omega_{ji}$ , compared to AK [3], FM[4]  
 $E=30$  Ryd. ( $a-b \equiv a \times 10^{-b}$ )

j	i	present	AK	FM	j	i	present	AK	FM
1	2	1.133-2	1.189-2	1.12-2	1	25	6.616-3	7.265-3	5.70-3
1	3	9.888-2	9.320-2	8.42-2	2	25	2.354-2	2.031-2	1.64-2
2	3	2.051-1	2.175-1	1.70-1	3	25	<b>6.842-0</b>	<b>7.274-0</b>	<b>1.00+1</b>
1	20	<b>1.857-0</b>	<b>1.980-0</b>	<b>3.20-0</b>	1	26	1.871-3	2.051-3	1.90-3
2	20	3.340-3	5.408-3	1.58-1	2	26	6.081-3	6.553-3	6.30-3
3	20	4.833-1	4.573-1	3.11-1	3	26	1.849-1	2.457-1	3.86-1
1	24	9.965-4	1.353-3	7.00-4					
2	24	<b>1.603-0</b>	<b>1.644-0</b>	<b>1.79-0</b>					
3	24	<b>2.252-0</b>	<b>2.317-0</b>	<b>3.83-0</b>					

Table 5.3 Collision strengths  $\Omega_{ji}$ , compared to AK [3], FM[4]  
 $E=45$  Ryd. ( $a-b \equiv a \times 10^{-b}$ )

j	i	present	AK	FM	j	i	present	AK	FM
1	2	7.216-3	7.643-3	7.00-3	1	25	6.662-3	7.131-3	5.00-3
1	3	9.644-2	9.608-2	7.47-2	2	25	1.696-2	1.852-2	1.34-2
2	3	1.952-1	2.390-1	1.47-1	3	25	<b>6.634-0</b>	<b>8.103-0</b>	<b>1.12+1</b>
1	20	<b>1.821-0</b>	<b>2.216-0</b>	<b>3.57-0</b>	1	26	1.187-3	1.310-3	1.20-3
2	20	1.977-3	3.915-3	1.74-1	2	26	3.677-3	3.938-3	3.70-3
3	20	4.746-1	5.122-1	3.44-1	3	26	1.771-1	2.683-1	4.26-1
1	24	6.013-4	8.449-4	4.00-4					
2	24	<b>1.567-0</b>	<b>1.837-0</b>	<b>1.99-0</b>					
3	24	<b>2.211-0</b>	<b>2.594-0</b>	<b>4.27-0</b>					