The effect of non-thermal distributions on coronal radiative losses

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Outline

Introduction

Π.

 κ -distributions Why κ -distributions? Definition and basic properties Contribution from the continuum Bound-bound radiative losses, element contribution, Fe ions Consequences

n-distributions
Why *n*-distributions?
Definition and basic properties, pseudotemperature
Bound-bound radiative losses, element contribution, Fe ions
Consequences

Introduction

Emissivity ε_{ij} of a spectral line line – transition from the level *i* to level *j* in a *k*-times ionized element *x* is given by:

$$\varepsilon_{ij} = \frac{hc}{\lambda_{ij}} A_{ij} n_i = \frac{hc}{\lambda_{ij}} \frac{A_{ij}}{n_e} \frac{n_i}{n_e^k} \frac{n_k^k}{n^k} \frac{n_k^k}{n^k} A_x \frac{n_H}{n_e} n_e^2$$
$$= G_{x,ij} (T, n_e, A_x, \lambda_{ij}, \kappa) n_e^2$$

FIP effect

> Radiative losses $E_{\rm R}$: $E_{\rm R}(T, n_{\rm e}, \kappa) = n_{\rm e}^2 \int \sum_{x,i,j} G_{x,ij}(A_x, \lambda_{ij}, T, n_{\rm e}, \kappa) d\lambda = n_{\rm e}^2 \Phi(T, \kappa)$ $-E_{\rm R}(T, n_{\rm e}, \kappa) + E_{\rm H} + \vec{\nabla}.\vec{F}_{\rm cond} = 0$





Synthetic spectra

Modification of CHIANTI, spectra for C,N,O,Ne,Mg,AI,Si,S,Ar,Ca,Fe,Ni



II. *k*-distributions: why

- Study the emission need to know about the microphysics
- Suprathermal component ("high-energy tail") present during flares and also in solar wind
- Maksimovic et al. (1997): solar wind velocity distribution is better approximated by one κ-distribution than with one or sum of two Maxwellians
- Collier (2004): if the mean particle energy is not held constant, the entropy is *not* maximalized by a Maxwellian distribution
- > If the order of the mean energy conserved is, entropy is maximalized by the κ -distribution
- > Dudík et al. (2009): considerable influence on the EUV filter responses
- Dzifčáková & Kulinová (2010a,b) κ-distribution diagnosed in transition region spectra (Si III line ratios) and in flares (Fe XII line ratios)
- Their presence in AR corona is still an open question. Nevertheless...



k-distributions: definition

Owocki & Scudder (1983), Dzifčáková (2006a):

$$f(E,\kappa)dE = A_{\kappa} \frac{1}{\sqrt{\pi} (k_{\rm B}T)^{3/2}} \frac{E^{1/2}dE}{\left(1 + \frac{E}{(\kappa - 3/2)k_{\rm B}T}\right)^{\kappa+1}}$$

$$\int_{10^{-7}}^{10^{-7}} \frac{10^{-7}}{10^{-6}} \frac{10^{-7}}{10^{-6}} \frac{10^{-7}}{10^{-6}} \frac{10^{-7}}{10^{-6}} \frac{10^{-7}}{10^{-7}} \frac{100}{100} \frac{1000}{1000} \frac{1000}{10000}$$

$$E \left[eV\right]$$

$$E \left[eV\right]$$

$$E^{1/2}dE$$

$$F^{1/2}dE$$





.....





III. *n*-distributions: why/when

Consistent with X-Ray spectra observed with flares (Dzifčáková et al. 2008, AA 488, 311; Kulinová et al., poster)



n-distributions: definition

Seely, Feldman & Doschek (1987), Dzifčáková & Kulinová (2001):

$$f(E,n)dE = B^{n} \frac{2}{\sqrt{\pi} (k_{\rm B}T)^{n/2+1}} E^{n/2} e^{-E/k_{\rm B}T} dE$$



 \succ n = 1 Maxwellian distribution

$$\langle E \rangle = \frac{3}{2} k_{\rm B} \tau$$

> Pseudo-temperature τ :

$$\tau = \left(\frac{n}{2} + 1\right)\frac{2T}{3}$$











Conclusions

Conditions in solar corona allow for non-thermal distributions. We examined their influence on the radiative-loss function.

- ➤ Total radiative losses slightly lesser (up to 60%) for *κ*-distribution,
 ⇒ lesser requirement for heating input
 higher for *n*-distribution
 ⇒ cooling processes for flares
- The behavior of radiative losses copies the behavior of individual ions/ lines
- Future work: brehmsstrahlung, effect of radiation field on spectral lines

