

Two Stories of Magnetic Catalysis

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Based on work with Stefano Bolognesi

Magnetic Catalysis

Klimenko; Gusynin, Miransky
and Shovkovy, '94

e.g. in $d = 2 + 1$

$$\lim_{m \rightarrow 0^\pm} \langle \bar{\Psi} \Psi \rangle = \mp \frac{B}{4\pi}$$

Story I: Magnetic Catalysis in the bulk

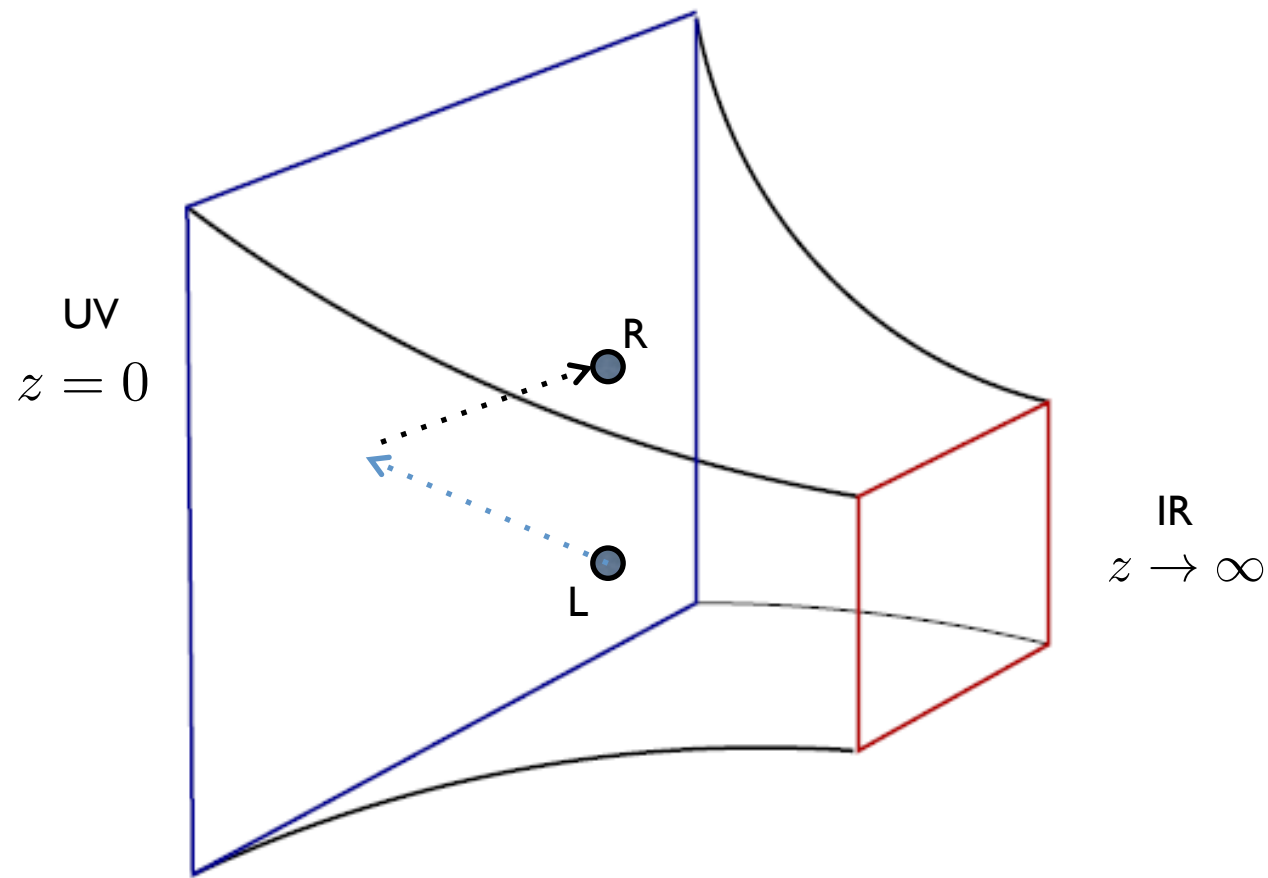
Massless Dirac Fermion in AdS_4

$$S = \frac{i}{2} \int_{\mathcal{M}} \sqrt{-g} \, \bar{\psi} \overleftrightarrow{\not{D}} \psi + \frac{i}{2} \int_{\partial\mathcal{M}} \sqrt{-h} \, \bar{\psi} \psi$$

4 component Dirac spinor



Flipping Helicity at the Boundary

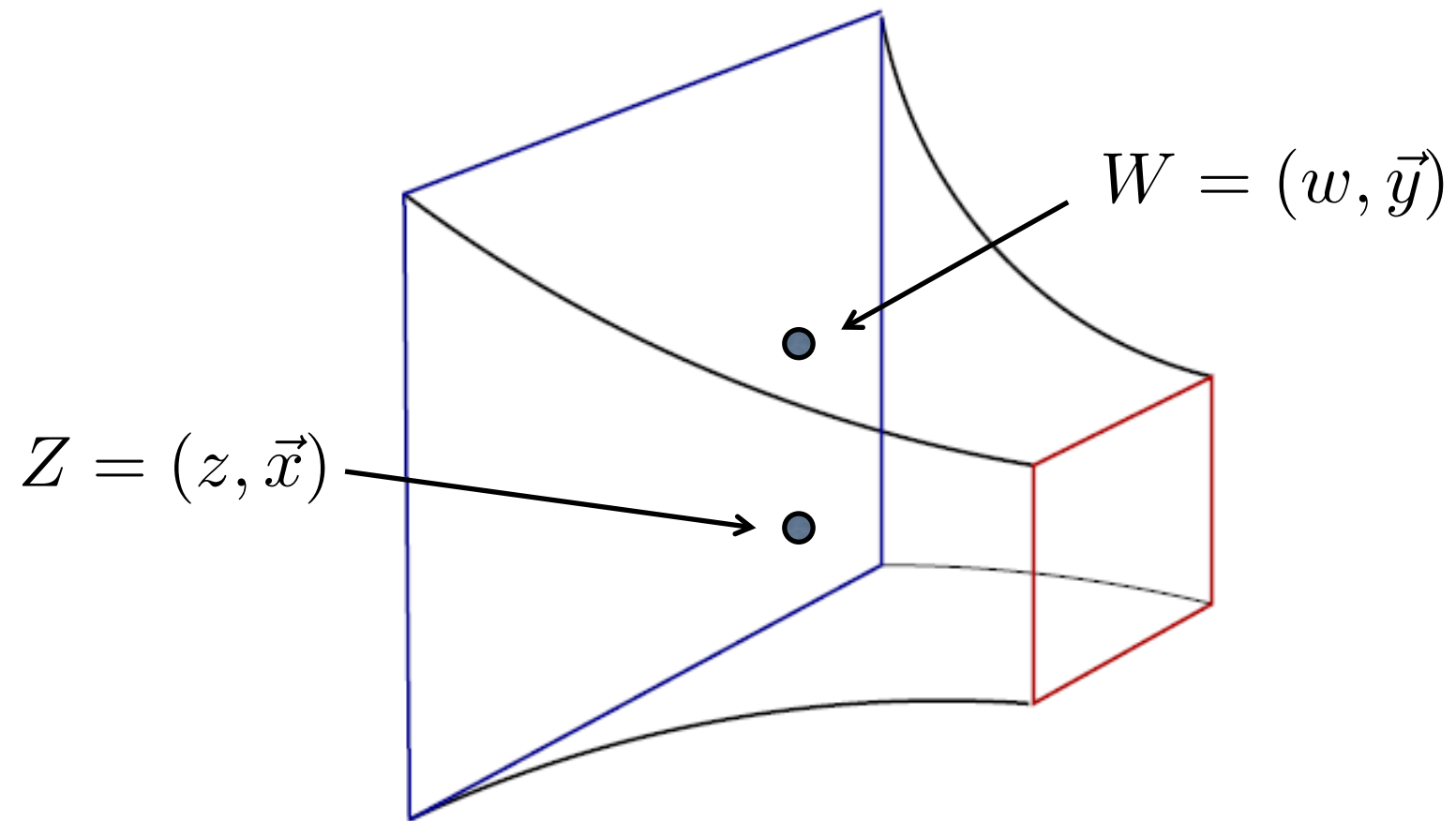


$$\begin{aligned}\psi_L &= \frac{1}{2}(1 - \gamma^5)\psi \longrightarrow \psi_R \\ \psi_R &= \frac{1}{2}(1 + \gamma^5)\psi \longrightarrow \psi_L\end{aligned}$$

$$ds^2 = L^2 \frac{dz^2 + \eta_{\mu\nu} dx^\mu dx^\nu}{z^2}$$

Chiral Symmetry Breaking in the Bulk

Allen and Lutken, '86

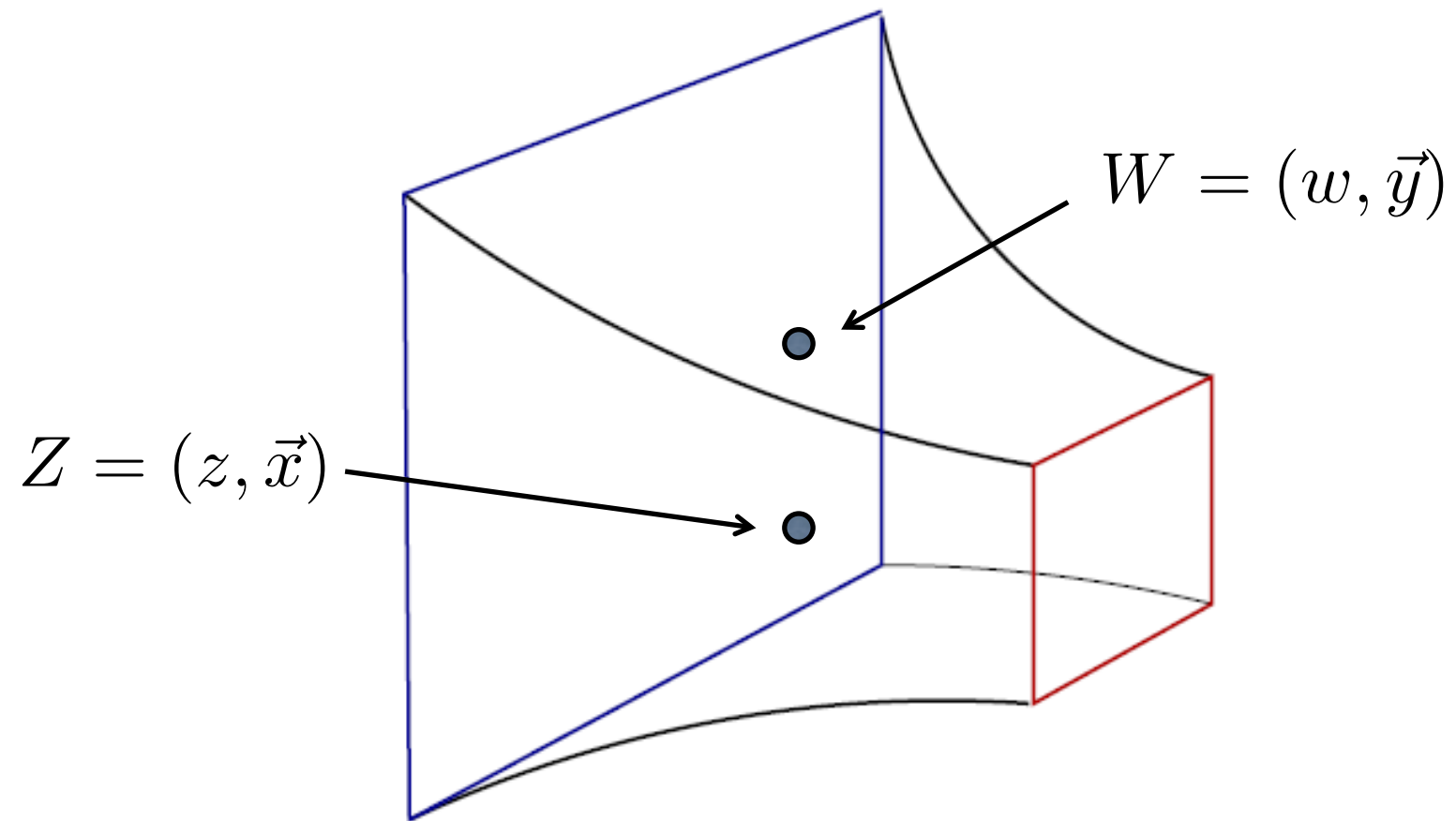


Bulk to bulk propagator for *massive* spinor

$$S(Z, W) = [\text{Stuff}] \times {}_2F_1[\text{Horrible Blah}]$$

Chiral Symmetry Breaking in the Bulk

Allen and Lutken, '86



Bulk to bulk propagator for *massless* spinor

$$S(Z, W) = \frac{1}{2\pi^2} \left(\frac{zw}{L^2} \right)^{3/2} \left[\frac{\not{Z} - \not{W}}{[(z-w)^2 + (\vec{x} - \vec{y})^2]^2} + \frac{\not{Z}\gamma^z + \gamma^z \not{W}}{[(z+w)^2 + (\vec{x} - \vec{y})^2]^2} \right]$$

Chiral Symmetry Breaking in the Bulk

Allen and Lutken, '86

$$S(Z, W) = \frac{1}{2\pi^2} \left(\frac{zw}{L^2} \right)^{3/2} \left[\frac{\not{Z} - \not{W}}{[(z-w)^2 + (\vec{x} - \vec{y})^2]^2} + \frac{\not{Z}\gamma^z - \gamma^z \not{W}}{[(z+w)^2 + (\vec{x} - \vec{y})^2]^2} \right]$$

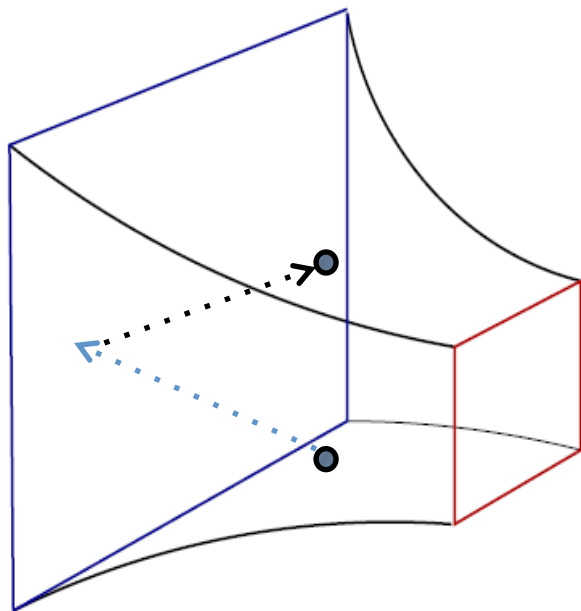
$$\Rightarrow \langle \bar{\psi} \psi \rangle = \lim_{Z \rightarrow W} \text{Tr } S(Z, W) = \frac{1}{4\pi^2 L^3}$$

Changing the Boundary Conditions

Porrati & Giradello, Rattazzi and Redi, '09

Chiral symmetry broken by boundary conditions \Rightarrow circle of boundary conditions

$$S_{\text{bdy}} = \bar{\psi}\psi \longrightarrow \bar{\psi}e^{i\alpha\gamma^5}\psi$$



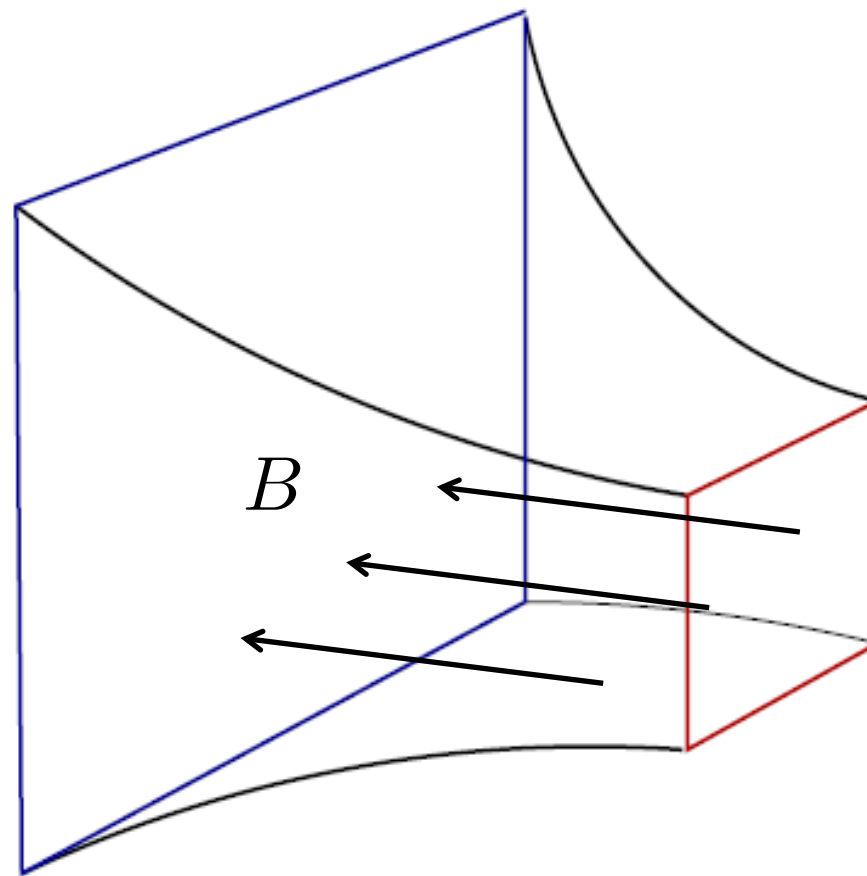
$$\psi_L = \frac{1}{2}(1 - \gamma^5)\psi \longrightarrow e^{-i\alpha}\psi_R$$

$$\psi_R = \frac{1}{2}(1 + \gamma^5)\psi \longrightarrow e^{i\alpha}\psi_L$$

$$\begin{aligned} \Rightarrow \quad 2\langle\bar{\psi}_L\psi_R\rangle &= \langle\bar{\psi}\psi\rangle + i\langle\bar{\psi}\gamma^5\psi\rangle \\ &= \frac{e^{i\alpha}}{4\pi^2 L^3} \end{aligned}$$

A Magnetic Field in AdS

(Should work in Reissner-Nordstrom black hole; ignore backreaction for now)



A Magnetic Field in AdS

Solve the Dirac Equation in background B field

Strategy

- Ignore radial z direction
- Landau levels are given

$$E = \sqrt{2Bn} \quad n \in \mathbf{Z}^+$$

- Zero modes ($n=0$) sector are ξ_α (2 component spinor)
- Promote $\xi_\alpha \rightarrow \xi_\alpha(z, t)$
- Effective 2d action

$$S_{\text{eff}} = \int dt dz \sqrt{-g} i \bar{\xi} \not{D}_{2d} \xi$$

Condensate in Magnetic Field

$$\langle \bar{\xi} \xi \rangle = \frac{e^{i\alpha}}{2\pi} \frac{z^2}{L^3}$$

$$\langle \bar{\psi}_L \psi_R \rangle = \frac{B}{2\pi} \langle \bar{\xi} \xi \rangle$$

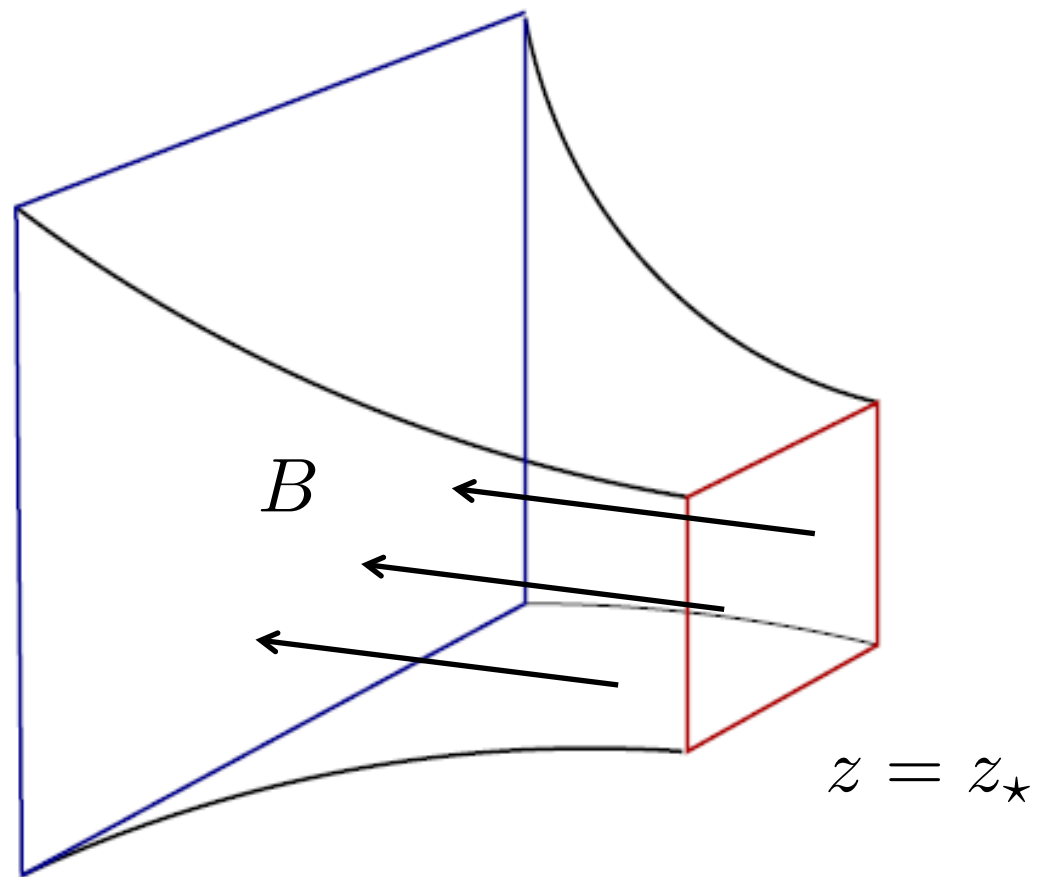


Density of states in lowest Landau level

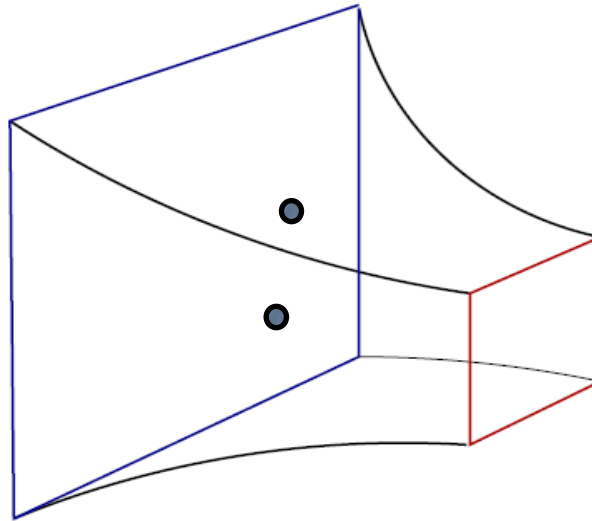
Caveat

- Gap to higher levels is warped in AdS to $\sim \sqrt{B} L z$
- Higher Landau levels important near boundary

Condensate in Magnetic Field with Hard Wall



Condensate in Magnetic Field with Hard Wall



First going left:

1 bounce



3 bounces



5 bounces



$$\langle \bar{\xi}_L \xi_R \rangle = \frac{1}{\pi} \left(\frac{z}{L} \right)^3 \left[\frac{e^{i\alpha}}{2z} + \frac{e^{2i\alpha}}{2z + 2z_*} + \frac{e^{3i\alpha}}{2(z + 2z_*)} + \dots \right]$$

$$- \frac{1}{\pi} \left(\frac{z}{L} \right)^3 \left[\frac{1}{2(z_* - z)} + \frac{e^{-i\alpha}}{2(2z_* - z)} + \frac{e^{-2i\alpha}}{2(3z_* - z)} + \dots \right]$$



1 bounce



3 bounces



5 bounces

First going right:

Condensate in Magnetic Field with Hard Wall

$$\langle \bar{\psi}_L \psi_R \rangle = e^{i\alpha} \frac{B}{4\pi^2} \left(\frac{z}{L} \right)^3 \sum_{n=-\infty}^{+\infty} \frac{e^{in\alpha}}{z + nz_\star}$$

Series is conditionally convergent, but not absolutely convergent.

$$\lim_{\alpha \rightarrow 0^\pm} \langle \bar{\psi}_L \psi_R \rangle = \text{real} \pm \frac{i}{2\pi} \frac{Bz^3}{L^3 z_\star}$$

$$\text{c.f.} \quad \int_{-\infty}^{+\infty} dn \frac{e^{i\alpha n}}{z + nz_\star} = \text{real} \pm \frac{i\pi}{z_\star} \text{sign}(\alpha)$$

Punchline of Story I

$$\lim_{\alpha \rightarrow 0^\pm} \langle \bar{\psi} \gamma^5 \psi \rangle = \pm \frac{1}{2\pi} \frac{B z^3}{L^3 z_\star}$$

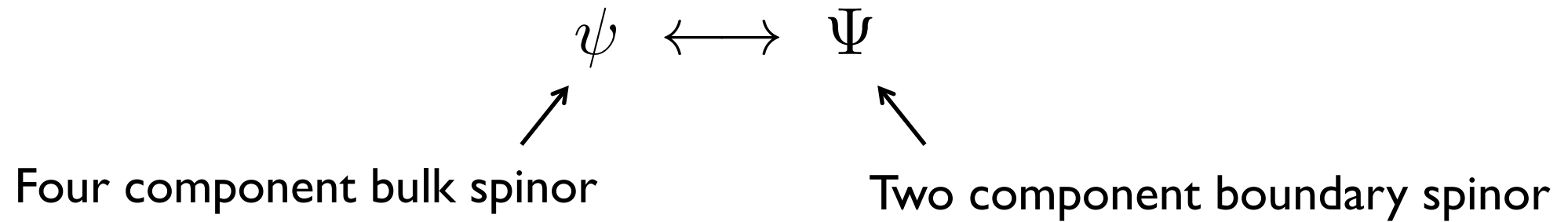
“Chiral” Symmetry Breaking in $d=3+1$

	B	$\bar{\psi}\psi$	$\bar{\psi}\gamma^5\psi$
P	-	+	-
C	-	+	+
CP	+	+	-

Spontaneous breaking of CP

Story 2: Magnetic Catalysis in the boundary

Usual AdS/CFT Dictionary



$$\Delta[\Psi] = \frac{3}{2} + mL = \frac{3}{2}$$

But...

$$\bar{\psi}\psi \longleftrightarrow ?$$

$$\bar{\psi}\gamma^5\psi \longleftrightarrow ?$$

Our Claim

$$\bar{\psi}\gamma^5\psi \iff \bar{\Psi}\Psi$$

Symmetries

	B	$\bar{\psi}\psi$	$\bar{\psi}\gamma^5\psi$	$\bar{\Psi}\Psi$
P	-	+	-	-
C	-	+	+	+
CP	+	+	-	-

Implication

$$\bar{\psi}\gamma^5\psi \longrightarrow \frac{\sin \alpha}{4\pi^2 L^3} + \frac{1}{2\pi} \frac{B}{L^3 z_\star} z^3 + \dots$$

Source for $\bar{\Psi}\Psi$
(Arises due to boundary terms)

Expectation value $\langle \bar{\Psi}\Psi \rangle$

Punchline of Story 2

$$\langle \bar{\Psi} \Psi \rangle = \pm \frac{3}{2} B M$$

with $M = \frac{1}{z_{\star}}$

Summary and Open Questions

- Hard Wall + Magnetic Field = Magnetic Catalysis
 - Spontaneous CP breaking
- $\bar{\psi}\psi \longleftrightarrow ?$
- Other effects due to bulk condensates?

Comments and Remarks

Caveat

- Gap to higher levels is warped in AdS to $\sim \sqrt{BL}z$
- Higher Landau levels important near boundary

Note: RN black hole

- Non-extremal: condensate diverges at the horizon
- Extremal: condensate constant at the horizon