

## Field Theory in Cosmology: Example Sheet 2

1. Reproduce the constraint equations by varying the action

$$S = \int d^4x \sqrt{h} N \left\{ \frac{M_{\text{Pl}}^2}{2} \left[ {}^{(3)}R + K_{ij}K^{ij} - K^2 \right] + P(X, \phi) \right\}. \quad (1)$$

with respect to  $N$  and  $N^i$ .

2. Derive the linear-order gauge transformations of  $A$ ,  $B$ ,  $\psi$  and  $h_{00}$ , for a generic change of coordinates  $x^\mu \rightarrow x^\mu + \epsilon^\mu$ .
3. Derive the gauge transformation from Newtonian gauge to flat gauge and viceversa. In particular, given some generic perturbations  $\{A^N, h_{00}^N, \varphi^N\}$  in Newtonian gauge, determine the corresponding perturbations  $\{h_{00}^f, \psi^f, \varphi^f\}$  in flat gauge.
4. Solve the  $\delta S / \delta N^i$  constraint to find  $\delta N$ , working in flat gauge to linear order.
5. In the lecture, we prove the conservation of  $\mathcal{R}$  and  $\gamma_{ij}$  on superHubble scales in the presence of a generic energy-momentum tensor  $T_{\mu\nu}$  by working in comoving gauge. Prove again the conservation of  $\mathcal{R}$  by working in Newtonian gauge. In particular, you might want to start with the change of coordinates

$$\epsilon^\mu = \{ \epsilon(t), \lambda x^i \}. \quad (2)$$

and show that the gauge transformations are

$$\begin{aligned} \Phi &= -\dot{\epsilon}, & \Psi &= H\epsilon - \frac{\lambda}{3}, \\ \delta\rho &= -\dot{\rho}\epsilon, & \delta u &= \epsilon, & \pi^S &= 0, \end{aligned} \quad (3)$$

$$\delta p = -\dot{p}\epsilon, \quad \varphi = -\epsilon\dot{\phi}. \quad (4)$$

Then use the scalar part of the  $ij$  components of the Einstein's equation,

$$k_i k_j (\Phi - \Psi) = 0, \quad (5)$$

to impose the physicality condition on  $\epsilon(t)$ . Your final result should be

$$\mathcal{R} = \frac{\lambda}{3}, \quad \varphi = -\dot{\phi} \frac{\mathcal{R}}{a} \int_T^t a(t') dt', \quad \Phi = \Psi = \mathcal{R} \left[ -1 + \frac{H}{a} \int_T^t a(t') dt' \right]. \quad (6)$$

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