

# 3P1a Quantum Field Theory: Example Sheet 1 Michaelmas 2023

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1. Show directly that if  $\phi(x)$  satisfies the Klein-Gordon equation, then  $\phi(\Lambda^{-1}x)$  also satisfies this equation for any Lorentz transformation  $\Lambda$ .
2. The motion of a complex field  $\psi(x)$  is governed by the Lagrangian density

$$\mathcal{L} = \partial_\mu \psi^* \partial^\mu \psi - m^2 \psi^* \psi - \frac{\lambda}{2} (\psi^* \psi)^2 .$$

Write down the Euler-Lagrange field equations for this system. Verify that the Lagrangian is invariant under the infinitesimal transformation

$$\delta\psi = i\alpha\psi , \quad \delta\psi^* = -i\alpha\psi^* .$$

Derive the Noether current associated with this transformation and verify explicitly that it is conserved using the field equations satisfied by  $\psi$ .

3. Verify that the Lagrangian density

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi_a \partial^\mu \phi_a - \frac{1}{2} m^2 \phi_a \phi_a ,$$

for a triplet of real fields  $\phi_a$ , where  $a \in \{1, 2, 3\}$  is invariant under the infinitesimal  $SO(3)$  rotation by  $\theta$

$$\phi_a \rightarrow \phi_a + \theta \epsilon_{abc} \eta_b \phi_c ,$$

where  $\eta_a$  is a unit vector. Compute the Noether current  $j^\mu$ . Deduce that the three quantities

$$Q_a = \int d^3x \epsilon_{abc} \dot{\phi}_b \phi_c$$

are all conserved and verify this directly using the field equations satisfied by  $\phi_a$ .

- 4\*. A Lorentz transformation  $x^\mu \rightarrow x'^\mu = \Lambda^\mu_\nu x^\nu$  is such that it preserves the Minkowski metric  $\eta_{\mu\nu}$ , meaning that  $\eta_{\mu\nu} x^\mu x^\nu = \eta_{\mu\nu} x'^\mu x'^\nu$  for all  $x$ . Show that this implies that

$$\eta_{\mu\nu} = \eta_{\sigma\tau} \Lambda^\sigma_\mu \Lambda^\tau_\nu .$$

Use this result to show that an infinitesimal transformation of the form

$$\Lambda^\mu_\nu = \delta^\mu_\nu + \alpha \omega^\mu_\nu$$

is a Lorentz transformation when  $\omega^{\mu\nu}$  is antisymmetric: i.e.  $\omega^{\mu\nu} = -\omega^{\nu\mu}$  ( $\alpha$  is considered to be infinitesimal).

Write down the matrix form for  $\omega^\mu_\nu$  that corresponds to a rotation through an infinitesimal angle  $\theta$  about the  $x^3$ -axis. Do the same for a boost along the  $x^1$ -axis by an infinitesimal velocity  $v$ .

- 5\* Consider the infinitesimal form of the Lorentz transformation derived in the previous question:  $x^\mu \rightarrow x^\mu + \alpha \omega^\mu_\nu x^\nu$ . Show that a scalar field transforms as

$$\phi(x) \rightarrow \phi'(x) = \phi(x) - \alpha \omega^\mu_\nu x^\nu \partial_\mu \phi(x) ,$$

and hence show that the variation of the Lagrangian density is a total derivative

$$\delta \mathcal{L} = -\alpha \partial_\mu (\omega^\mu_\nu x^\nu \mathcal{L}) .$$

Using Noether's theorem, deduce the existence of the conserved current

$$j^\mu = -\omega^\rho_\nu [T^\mu_\rho x^\nu] .$$

The three conserved charges arising from spatial rotational invariance define the *total angular momentum* of the field. Show that these charges are given by

$$Q_i = \frac{1}{2} \epsilon_{ijk} \int d^3x (x^j T^{0k} - x^k T^{0j}) .$$

Derive the conserved charges arising from invariance under Lorentz boosts. Show that they imply

$$\frac{d}{dt} \int d^3x (x^i T^{00}) = \text{constant} ,$$

and interpret this equation.

6. Maxwell's Lagrangian for the electromagnetic field is

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} ,$$

where  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$  and  $A_\mu$  is the 4-vector potential. Show that  $\mathcal{L}$  is invariant under gauge transformations

$$A_\mu \rightarrow A_\mu + \partial_\mu \xi ,$$

where  $\xi = \xi(x)$  is a scalar field with arbitrary (differentiable) dependence on  $x$ .

Using Noether's theorem, and the spacetime translational invariance of the action, to construct the energy momentum tensor  $T^{\mu\nu}$  for the electromagnetic field. Show that the resulting object is neither symmetric nor gauge invariant. Consider a new tensor given by

$$\Theta^{\mu\nu} = T^{\mu\nu} - F^{\rho\mu} \partial_\rho A^\nu .$$

Show that this object also defines four currents. Moreover, show that it is symmetric, gauge invariant and traceless.

7. The Lagrangian density for a massive vector field  $C_\mu$  is given by

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m^2C_\mu C^\mu ,$$

where  $F_{\mu\nu} = \partial_\mu C_\nu - \partial_\nu C_\mu$ . Derive the equations of motion and show that when  $m \neq 0$  they imply

$$\partial_\mu C^\mu = 0 .$$

Further show that  $C_0$  can be eliminated completely in terms of other fields by

$$\partial_i \partial^i C_0 + m^2 C_0 = \partial^i \dot{C}_i . \quad (1)$$

Construct the canonical momenta  $\Pi_i$  conjugate to  $C_i$  where  $i \in \{1, 2, 3\}$  and show that the canonical momentum conjugate to  $C_0$  is vanishing. Construct the Hamiltonian density  $\mathcal{H}$  in terms of  $C_0, C_i$  and  $\Pi_i$  (NB: don't be concerned that the canonical momentum for  $C_0$  is vanishing.  $C_0$  is non-dynamical; it is determined entirely in terms of the other fields using Eq. (1)).

8. A class of interesting theories is invariant under the simultaneous scaling of all lengths by

$$x^\mu \rightarrow (x')^\mu = \lambda x^\mu \text{ and } \phi(x) \rightarrow \phi'(x) = \lambda^{-D} \phi(\lambda^{-1}x) . \quad (2)$$

Here,  $D$  is called the *scaling dimension* of the field. Consider the action for a real scalar field given by

$$S = \int d^4x \left( \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 - g \phi^p \right) .$$

Find the scaling dimension  $D$  such that the derivative terms remain invariant. For what values of  $m$  and  $p$  is the scaling in Eq. (2) a symmetry of the theory? How do these conclusions change for a scalar field living in an  $(n+1)$ -dimensional spacetime instead of a  $3+1$  dimensional spacetime?

In  $3+1$  dimensions, use Noether's theorem to construct the conserved current  $D^\mu$  associated with scaling invariance.