

Mathematical Tripos Part III
String Theory. Lent 2018: Example Sheet 1

1 Verify that the relativistic point particle action

$$I = \int_{t_a}^{t_B} dt \left\{ \dot{X}^m P_m - \frac{1}{2} e (P^2 + m^2) \right\}$$

is invariant (ignoring boundary terms) under the Diff_1 gauge transformations

$$\delta_\xi X = \xi \dot{X}, \quad \delta_\xi P = \xi \dot{P}, \quad \delta_\xi e = \partial_t(e\xi),$$

where $\xi(t)$ is an arbitrary infinitesimal function. Now verify that it is also invariant (ignoring boundary terms) under the “canonical” gauge transformation

$$\delta_\alpha X = \alpha P, \quad \delta_\alpha P = 0, \quad \delta e = \dot{\alpha},$$

where $\alpha(t)$ is an arbitrary infinitesimal function.

Show that any action functional $I[\psi, \phi]$ is invariant under the gauge transformation

$$\delta_f \psi = f \frac{\delta I}{\delta \phi}, \quad \delta_f \phi = -f \frac{\delta I}{\delta \psi},$$

for *arbitrary* infinitesimal function f . Such “trivial” gauge invariances have no physical effect. Use this observation to explain why the “ Diff_1 ” and “canonical” gauge invariances of the point particle action are physically equivalent.

2 A free non-relativistic point particle of mass m is described by the action

$$I[\mathbf{x}] = \int dt \left\{ \frac{1}{2} m \left| \frac{d\mathbf{x}}{dt} \right|^2 \right\}. \quad (\dagger)$$

Show, by elimination of variables, that an equivalent action is

$$I[t, \mathbf{x}, E, \mathbf{p}; e] = \int d\lambda \left\{ \dot{\mathbf{x}} \cdot \mathbf{p} - \dot{t} E + e \left(E - \frac{|\mathbf{p}|^2}{2m} \right) \right\} \quad \left(\dot{\mathbf{x}} = \frac{d\mathbf{x}}{d\lambda} \right) \quad (\star)$$

where λ is an arbitrary worldline parameter. Find the gauge transformations of the canonical variables (\mathbf{x}, \mathbf{p}) and (t, E) generated by the constraint function, and verify that the action (\star) is invariant for an appropriate gauge transformation of e . Now show that this gauge invariance can be fixed by imposing the temporal gauge condition $t = \lambda$, and thereby recover the action (\dagger) .

Let \mathbf{P} be the Noether charge corresponding to translation invariance $\mathbf{x} \rightarrow \mathbf{x} + \mathbf{a}$. Show that $\mathbf{P} = \mathbf{p}$. Verify that the action (\star) is invariant (neglecting a time-boundary term) under the Galilean boost transformation

$$\delta_{\mathbf{v}} \mathbf{x} = \mathbf{v} t, \quad \delta_{\mathbf{v}} \mathbf{p} = m \mathbf{v}, \quad \delta_{\mathbf{v}} E = \mathbf{v} \cdot \mathbf{p}.$$

Deduce that the corresponding Noether charge is $\mathbf{B} = t\mathbf{p} - m\mathbf{x}$, and verify that $\dot{\mathbf{B}} = 0$ as a consequence of the equations of motion. Show that $\{P, B\}_{PB} = m$.

Apply Dirac’s method of quantization with constraints to (\star) to show that the wavefunction $\psi(t, \mathbf{x})$ of the quantum non-relativistic free particle satisfies the time-dependent Schrodinger equation

$$-\frac{1}{2m} \nabla^2 \psi = \frac{i}{\hbar} \partial_t \psi.$$

3 In the temporal gauge $x^0(t) = t$ the relativistic point particle has the action

$$I = \int dt \left\{ \dot{\vec{x}} \cdot \vec{p} - p^0 \right\}, \quad p^0 = \sqrt{|\vec{p}|^2 + m^2}.$$

Verify that this is invariant under the infinitesimal transformation

$$\delta \vec{x} = \vec{\Lambda} t - (\vec{\Lambda} \cdot \vec{x}) \vec{p}/p_0, \quad \delta \vec{p} = \vec{\Lambda} p_0,$$

for constant vector parameter $\vec{\Lambda}$. Find the corresponding Noether charge. What is the significance of this symmetry?

4 The Hamiltonian form of the action for a relativistic p -brane of tension T_p is

$$I = \int dt \int d^p \sigma \left\{ \dot{X}^m P_m - \frac{1}{2} e (P^2 + T_p^2 \det k) - u^i \partial_i X^m P_m \right\}$$

where k is the induced p -metric on the brane in coordinates σ^i ($i = 1, \dots, p$). By elimination of P_m and the Lagrange multipliers (e, u^i) , verify that this action is equivalent to the Dirac action $S = -T_p \int dt \int d^p \sigma \sqrt{-\det g}$, where g is the induced worldvolume $(p+1)$ -metric. [Hint: use the identity $\det g \equiv (g_{tt} - k^{ij} g_{ti} g_{tj}) \det k$.]

5 Explain why the following configuration describes a circular closed string of time-dependent radius $R(t) = R_0 |\cos(t/R)|$.

$$X^0 = t, \quad X^1 = R_0 \cos\left(\frac{t}{R_0}\right) \cos\left(\frac{\sigma}{R_0}\right), \quad X^2 = R_0 \cos\left(\frac{t}{R_0}\right) \sin\left(\frac{\sigma}{R_0}\right),$$

with R_0 a positive constant and $X^3 = X^4 = \dots = 0$. Compute the induced worldsheet metric and hence show (i) that the above configuration solves the NG equations of motion and (ii) that $2\pi R(t)$ is the proper length of the string. Write down an expression for the total energy E and show that $E = (2\pi R_0)T$. Deduce that the rest-energy density of the string is T .

6 Write down the Hamiltonian form of the NG string action. In the case of an open string with free ends, what are the boundary conditions on X ? Why do the ends of the string move at the speed of light?

Explain why the following configuration describes a rigidly rotating straight string in which the ends of the string are at $\sigma = 0$ and $\sigma = \pi L/2$.

$$X^0 = t \quad X^1 = \frac{L}{2} \sin\left(\frac{2t}{L}\right) \cos\left(\frac{2\sigma}{L}\right), \quad X^2 = \frac{L}{2} \cos\left(\frac{2t}{L}\right) \cos\left(\frac{2\sigma}{L}\right),$$

with $X^3 = X^4 = \dots = 0$. Compute the induced worldsheet metric, and show that it is conformally flat. Use your result to show that L is the proper length of the string, and that the above configuration solves the Nambu-Goto equations of motion.

Write down expressions for the conserved total energy E and angular-momentum J of an open string of tension T with free ends, and use it to show that a rigidly rotating open string of proper length L has

$$E = \frac{L}{4\alpha'}, \quad J = \frac{L^2}{16\alpha'} \quad \left(\alpha' = \frac{1}{2\pi T} \right).$$

Hence show that $J = \alpha' E^2$. [In a “Regge plot” of J against E^2 the slope is α' , so α' is called the “Regge slope parameter”.]