## 4.4 Flow past a sphere at low Reynolds number

Uniform flow  $\mathbf{U}$  past a fixed rigid sphere, radius a. There are several methods, all of which have heavy algebra somewhere.

## 4.4.3 Method 3

The linearity of the Stokes equations means that  $\mathbf{u}(\mathbf{x})$  must be linear in U. Further, the problem has spherical symmetry about the centre of the sphere, which is taken as the origin. The velocity and pressure fields must therefore take the forms

$$\begin{aligned} \mathbf{u}(\mathbf{x}) &= \mathbf{U}f(r) + \mathbf{x}(\mathbf{U} \cdot \mathbf{x})g(r), \\ p(\mathbf{x}) &= \mu(\mathbf{U} \cdot \mathbf{x})h(r), \end{aligned}$$

where  $r = |\mathbf{x}|$ , and f, g and h are functions of scalar r to be determined.

Now

$$\frac{\partial u_i}{\partial x_j} = U_i x_j f'/r + \delta_{ij} U_n x_n g + x_i U_j g + x_i x_j U_n x_n g'/r.$$

Contracting i with j, we have the incompressibility condition

$$0 = \nabla \cdot \mathbf{u} = U_n x_n (f'/r + 4g + rg').$$

Differentiating again

$$\mu \nabla^2 u_i = \mu U_i \left( f'' + 2f'/r + 2g \right) + \mu x_i U_n x_n \left( g'' + 6g'/r \right)$$
  
 
$$\nabla_i p(\mathbf{x}) = \mu U_i h + \mu x_i U_n x_n h'/r$$

Hence the governing equations give

$$f'/r + 4g + rg' = 0$$
,  $f'' + 2f'/r + 2g = h$  and  $g'' + 6g'/r = h'/r$ .

Eliminating h and then f yields

$$r^2 g''' + 11rg'' + 24g' = 0.$$

This differential equation is homogeneous in r so there are solutions of the form  $g = r^{\alpha}$ . Substituting, one finds  $\alpha = 0, -3$  and -5, with associated  $f = -(\alpha + 4)r^{\alpha+2}/(\alpha + 2) + f_0$ and  $h = -(\alpha + 5)(\alpha + 2)r^{\alpha}$ . Hence the general solution of the assumed form linear in **U** is

$$\mathbf{u}(\mathbf{x}) = \mathbf{U} \left( -2Ar^2 + B + Cr^{-1} - \frac{1}{3}Dr^{-3} \right) + \mathbf{x}(\mathbf{U} \cdot \mathbf{x}) \left( A + Cr^{-3} + Dr^{-5} \right),$$
  
 
$$p(\mathbf{x}) = \mu(\mathbf{U} \cdot \mathbf{x}) \left( -10A + 2Cr^{-3} \right).$$

We shall need the stress exerted across a spherical surface with unit normal  $\mathbf{n}=\mathbf{x}/r$ 

$$\boldsymbol{\sigma} \cdot \mathbf{n} = \mu \mathbf{U} \left( -3Ar + 2Dr^{-4} \right) + \mu \mathbf{x} (\mathbf{U} \cdot \mathbf{x}) \left( 9Ar^{-1} - 6Cr^{-4} - 6Dr^{-6} \right)$$

Applying the boundary conditions on the rigid sphere and for the far field, we find the coefficients  $D = \frac{3}{2} + \frac{3}{2} +$ 

$$A = 0,$$
  $B = 1,$   $C = -\frac{3}{4}a$  and  $D = \frac{3}{4}a^{3},$ 

 $\mathbf{SO}$ 

$$\mathbf{u} = \mathbf{U}\left(1 - \frac{3a}{4r} - \frac{a^3}{4r^3}\right) + \mathbf{x}(\mathbf{U} \cdot \mathbf{x})\left(-\frac{3a}{4r^3} + \frac{3a^3}{4r^5}\right),$$

$$p = -\frac{3a\mu \mathbf{U} \cdot \mathbf{x}}{2r^3}$$

$$|\boldsymbol{\sigma} \cdot \mathbf{n}|_{r=a} = \frac{3\mu}{2a} \mathbf{U}.$$

Hence the drag on the sphere is

$$\int_{r=a} \boldsymbol{\sigma} \cdot \mathbf{n} \, dS = 4\pi a^2 \frac{3\mu}{2a} \mathbf{U} = 6\pi \mu a \mathbf{U}.$$