## Mathematical Tripos Part II: Michaelmas Term 2023 Numerical Analysis - Lecture 8

We continue to study stability of discretization schemes using the Fourier analytic method from the previous lecture.

Advection equation We look at the advection equation which we already considered in Lecture 6.

$$
\begin{equation*}
u_{t}=u_{x}, \quad t \geq 0, \tag{2.6}
\end{equation*}
$$

where $u=u(x, t)$. It is given with the initial condition $u(x, 0)=\varphi(x)$. The exact solution of (2.6) is simply $u(x, t)=\varphi(x+t)$, a unilateral shift leftwards.

1) Downwind instability: Consider the discretization $\frac{\partial u_{m}(t)}{\partial x} \approx \frac{1}{2 h}\left[u_{m}(t)-u_{m-1}(t)\right]$, so coming to the ODE $u_{m}^{\prime}(t)=\frac{1}{2 h}\left[u_{m}(t)-u_{m-1}(t)\right]$. For the Euler method, the outcome is

$$
u_{m}^{n+1}=u_{m}^{n}+\mu\left(u_{m}^{n}-u_{m-1}^{n}\right), \quad n \in \mathbb{Z}_{+} .
$$

We can analyze the stability of this method using Fourier analysis. The amplification factor is

$$
H(\theta)=1+\mu-\mu e^{-\mathrm{i} \theta} .
$$

We see that for $\theta=\pi / 2,|H(\theta)|^{2}=(1+\mu)^{2}+\mu^{2}>1$, and so the method is unstable for all $\mu>0$.
2) Upwind scheme: If we semidiscretize $\frac{\partial u_{m}(t)}{\partial x} \approx \frac{1}{h}\left[u_{m+1}(t)-u_{m}(t)\right]$, and solve the ODE again by Euler's method, then the result is

$$
\begin{equation*}
u_{m}^{n+1}=u_{m}^{n}+\mu\left(u_{m+1}^{n}-u_{m}^{n}\right), \quad n \in \mathbb{Z}_{+} \tag{2.7}
\end{equation*}
$$

The local error is $\mathcal{O}\left(k^{2}+k h\right)$ which is $\mathcal{O}\left(h^{2}\right)$ for a fixed $\mu$, hence convergence if the method is stable. We can again use Fourier analysis to analyze stability. The amplification factor is

$$
H(\theta)=1-\mu+\mu e^{\mathrm{i} \theta}
$$

and we see that $|H(\theta)|=\left|1-\mu+\mu e^{\mathrm{i} \theta}\right| \leq|1-\mu|+\mu=1$ for $\mu \in[0,1]$. Hence we have stability for $\mu \leq 1$. If $\mu>1$, then note that $|H(\pi)|=|1-2 \mu|>1$, and so we have instability for $\mu>1$.
Matlab demo: Download the Matlab GUI for Solving the Advection Equation, Upwinding and Stability from https://www.damtp.cam.ac.uk/user/hf323/M21-II-NA/demos/index.html and solve the advection equation 2.6 with the different methods provided in the demonstration. Experience what can go wrong when "winding" in the wrong direction!
3) Leap-frog method: We semidicretize 2.6 as $\frac{\partial u_{m}(t)}{\partial x} \approx \frac{1}{2 h}\left[u_{m+1}(t)-u_{m-1}(t)\right]$, but now solve the ODE with the second-order midpoint rule

$$
\boldsymbol{y}_{n+1}=\boldsymbol{y}_{n-1}+2 k \boldsymbol{f}\left(t_{n}, \boldsymbol{y}_{n}\right), \quad n \in \mathbb{Z}_{+} .
$$

The outcome is the two-step leapfrog method

$$
\begin{equation*}
u_{m}^{n+1}=\mu\left(u_{m+1}^{n}-u_{m-1}^{n}\right)+u_{m}^{n-1} . \tag{2.8}
\end{equation*}
$$

The error is now $\mathcal{O}\left(k^{3}+k h^{2}\right)=\mathcal{O}\left(h^{3}\right)$. We analyse stability by the Fourier technique. Thus, proceeding as before,

$$
\begin{equation*}
\widehat{u}^{n+1}(\theta)=\mu\left(\mathrm{e}^{\mathrm{i} \theta}-\mathrm{e}^{-\mathrm{i} \theta}\right) \widehat{u}^{n}(\theta)+\widehat{u}^{n-1}(\theta) \tag{2.9}
\end{equation*}
$$

whence

$$
\widehat{u}^{n+1}(\theta)-2 \mathrm{i} \mu \sin \theta \widehat{u}^{n}(\theta)-\widehat{u}^{n-1}(\theta)=0, \quad n \in \mathbb{Z}_{+},
$$

and our goal is to determine values of $\mu$ such that $\left|\widehat{u}^{n}(\theta)\right|$ is uniformly bounded for all $n, \theta$. This is a difference equation $w_{n+1}+b w_{n}+c w_{n-1}=0$ with the general solution $w_{n}=c_{1} \lambda_{1}^{n}+c_{2} \lambda_{2}^{n}$, where $\lambda_{1}, \lambda_{2}$ are the roots of the characteristic equation $\lambda^{2}+b \lambda+c=0$, and $c_{1}, c_{2}$ are constants, dependent on the initial values $w_{0}$ and $w_{1}$. If $\lambda_{1}=\lambda_{2}$, then solution is $w_{n}=\left(c_{1}+c_{2} n\right) \lambda^{n}$. In our case, we obtain

$$
\lambda_{1,2}(\theta)=\mathrm{i} \mu \sin \theta \pm \sqrt{1-\mu^{2} \sin ^{2} \theta} .
$$

Stability is equivalent to $\left|\lambda_{1,2}(\theta)\right| \leq 1$ for all $\theta$ and this is true if and only if $\mu \leq 1$.

The wave equation Consider the wave equation

$$
\frac{\partial^{2} u}{\partial t^{2}}=\frac{\partial^{2} u}{\partial x^{2}} \quad t \geq 0
$$

given with initial conditions $u(x, 0)$ and $u_{t}(x, 0)=\frac{\partial u}{\partial t}(x, 0)$. The usual approximation looks as follows

$$
u_{m}^{n+1}-2 u_{m}^{n}+u_{m}^{n-1}=\mu\left(u_{m+1}^{n}-2 u_{m}^{n}+u_{m-1}^{n}\right)
$$

with the Courant number being now $\mu=k^{2} / h^{2}$.
The Fourier analysis (for Cauchy problem) provides

$$
\widehat{u}^{n+1}(\theta)-2 \widehat{u}^{n}(\theta)+\widehat{u}^{n-1}(\theta)=-4 \mu \sin ^{2} \frac{\theta}{2} \widehat{u}^{n}(\theta),
$$

with the characteristic equation $\lambda^{2}-2\left(1-2 \mu \sin ^{2} \frac{\theta}{2}\right) \lambda+1=0$. The product of the roots is one, therefore stability (that requires the moduli of both $\lambda$ to be at most one) is equivalent to the roots being complex conjugate, so we require

$$
\left(1-2 \mu \sin ^{2} \frac{\theta}{2}\right)^{2} \leq 1
$$

This condition is achieved if and only if $\mu=k^{2} / h^{2} \leq 1$.

