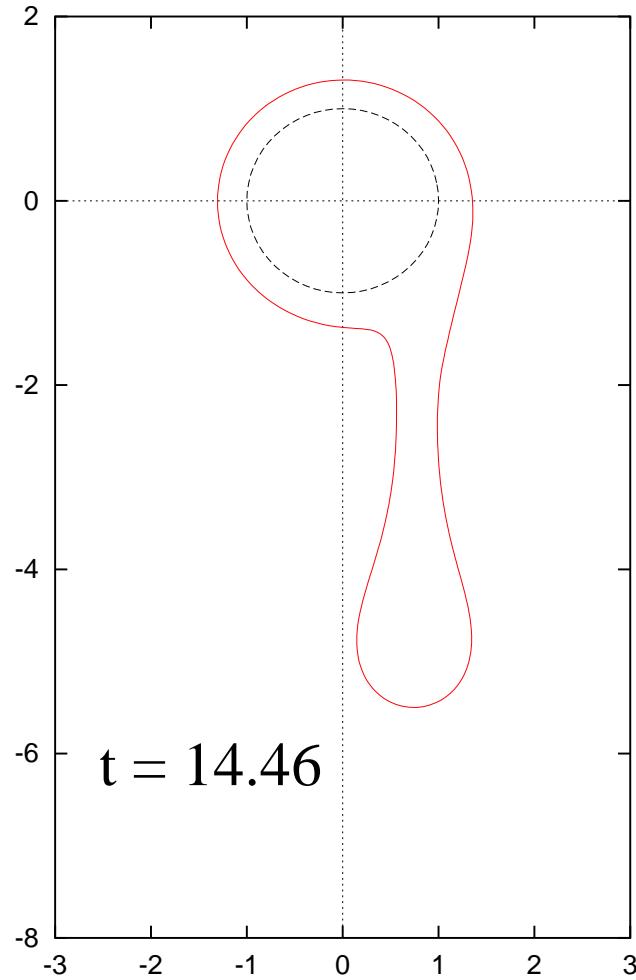


Shocking, the slow changes in HKM's thin coat

by HKM2 = Hinch, Kelmanson, Metcalfe (27 years later)

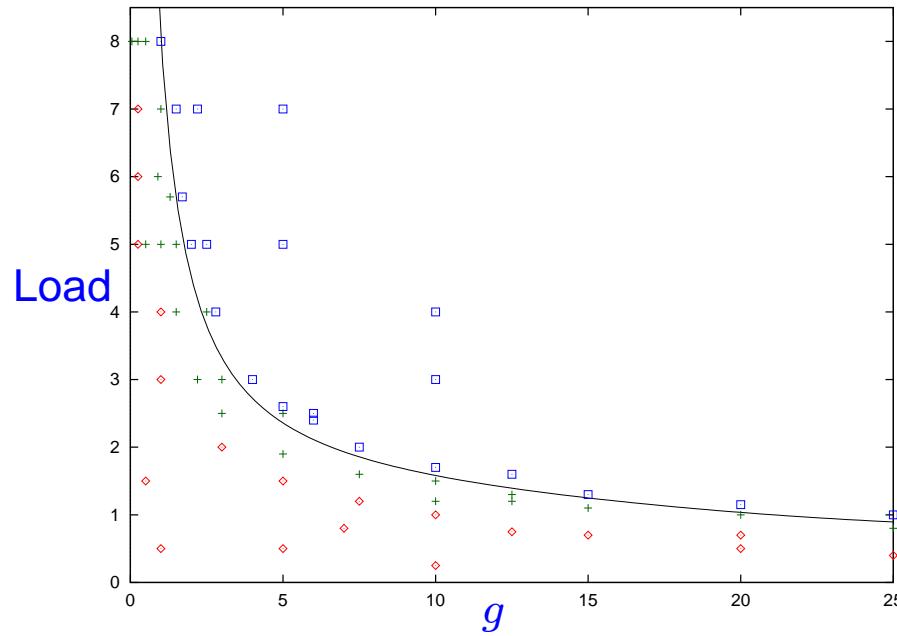
- Coating of a viscous liquid on the outside of a horizontal rotating cylinder
- Steady solution (for maximum load):
 - H.K. Moffatt (1977) *J. Fluid Mech.* 187
 - V.V. Pukhnachev (1977) *Z. Prik. Mekh. Tekh. Fiz.* 3
- Transients?

Dripping if load too large



FEM solution by Peterson,
Jimack & Kelmanson 2001

Curious oscillations



Load

g

Drip \square , steady +, oscillate \diamond

FEM solutions by Peterson, Jimack & Kelman-
son 2001

Or slowly decaying?

Lubrication theory (2D)

Mass conservation

$$h_t + q_x = 0$$

Flux

$$q = \omega a h - \frac{h^3}{3\mu} (\rho g \cos \theta - p_x)$$

Capillary pressure

$$p = \sigma(h_{xx} + h/a^2)$$

with $x = a\theta$.

Initial condition: $h(\theta, t = 0) = h_0$, uniform thickness.

Naive expansion (page 1 of 8)

[>

[17-term naive expansion solution of Hinch & Kelmanson (2002) eqn (2.4)

$$\begin{aligned}
 Hhuge := & \textcolor{red}{ho} + \textcolor{red}{ho}^3 \left(\frac{1}{3} \cos(\theta) \gamma_0 - \frac{1}{3} \cos(\theta - t) \gamma_0 \right) \\
 & + \textcolor{red}{ho}^5 \left(\frac{1}{6} \cos(2\theta) \gamma_0^2 - \frac{1}{3} \cos(2\theta - t) \gamma_0^2 + \frac{1}{6} \cos(2\theta - 2t) \gamma_0^2 \right) + \textcolor{red}{ho}^7 \left(-\frac{1}{12} \cos(\theta - t) \gamma_0^3 \right. \\
 & \left. + \frac{1}{9} \cos(3\theta) \gamma_0^3 - \frac{1}{9} \cos(\theta - 2t) \gamma_0^3 + \frac{1}{3} \cos(3\theta - 2t) \gamma_0^3 - \frac{1}{3} \cos(3\theta - t) \gamma_0^3 + \frac{2}{9} \cos(\theta) \gamma_0^3 \right. \\
 & \left. - \frac{1}{36} \cos(\theta + t) \gamma_0^3 - \frac{1}{9} \cos(3\theta - 3t) \gamma_0^3 + \frac{5}{18} \gamma_0^3 \sin(\theta - t) t \right) \\
 & + \textcolor{red}{ho}^8 \left(-\frac{1}{3} \sin(2\theta) \alpha \gamma_0^2 + \frac{4}{3} \sin(2\theta - t) \alpha \gamma_0^2 - \sin(2\theta - 2t) \alpha \gamma_0^2 - \frac{2}{3} \alpha \gamma_0^2 \cos(2\theta - 2t) t \right) + \\
 & \textcolor{red}{ho}^9 \left(\frac{55}{108} \cos(4\theta - 2t) \gamma_0^4 + \frac{55}{648} \cos(4\theta) \gamma_0^4 + \frac{95}{324} \cos(2\theta) \gamma_0^4 - \frac{55}{108} \cos(2\theta - t) \gamma_0^4 \right. \\
 & \left. + \frac{5}{18} \gamma_0^4 t \sin(2\theta - t) - \frac{55}{162} \cos(4\theta - t) \gamma_0^4 + \frac{55}{324} \cos(2\theta - 3t) \gamma_0^4 - \frac{5}{162} \cos(2\theta + t) \gamma_0^4 \right. \\
 & \left. + \frac{55}{648} \cos(4\theta - 4t) \gamma_0^4 - \frac{25}{54} \gamma_0^4 \sin(2\theta - 2t) t - \frac{55}{162} \cos(4\theta - 3t) \gamma_0^4 + \frac{25}{324} \cos(2\theta - 2t) \gamma_0^4 \right) + \\
 & \textcolor{red}{ho}^{10} \left(-\frac{5}{12} \gamma_0^3 \sin(\theta - t) \alpha - \frac{11}{9} \gamma_0^3 \sin(3\theta) \alpha - 10 \gamma_0^3 \sin(3\theta - 2t) \alpha + \gamma_0^3 \sin(\theta - 2t) \alpha \right. \\
 & \left. - \gamma_0^3 \alpha t \cos(3\theta - 2t) + \frac{1}{3} \gamma_0^3 \alpha t \cos(\theta - 2t) + \frac{19}{6} \gamma_0^3 \alpha \cos(3\theta - 3t) t + \frac{197}{36} \gamma_0^3 \sin(3\theta - 3t) \alpha \right. \\
 & \left. + \frac{7}{6} \gamma_0^3 \alpha \cos(\theta - t) t + \frac{23}{4} \gamma_0^3 \sin(3\theta - t) \alpha + \frac{1}{12} \gamma_0^3 \sin(\theta + t) \alpha - \frac{2}{3} \gamma_0^3 \sin(\theta) \alpha \right) + \textcolor{red}{ho}^{11} \left(\right. \\
 & \left. - \frac{455}{1296} \cos(5\theta - t) \gamma_0^5 - \frac{91}{432} \cos(3\theta - 4t) \gamma_0^5 + \frac{455}{1296} \cos(5\theta - 4t) \gamma_0^5 + \frac{455}{648} \cos(5\theta - 2t) \gamma_0^5 \right. \\
 & \left. + \frac{5}{162} \cos(3\theta - 3t) \gamma_0^5 + \frac{1}{648} \cos(\theta + 2t) \gamma_0^5 - \frac{115}{1296} \cos(\theta + t) \gamma_0^5 - \frac{7}{216} \cos(3\theta + t) \gamma_0^5 \right. \\
 & \left. + \frac{91}{1296} \cos(5\theta) \gamma_0^5 - \frac{265}{648} \cos(\theta - 2t) \gamma_0^5 + \frac{5}{9} \gamma_0^5 t \sin(\theta - t) + \frac{5}{18} \gamma_0^5 t \sin(\theta - 2t) \right. \\
 & \left. - \frac{5}{6} \gamma_0^5 t \sin(3\theta - 2t) - \frac{5}{216} \gamma_0^5 t \sin(\theta + t) + \frac{5}{18} \gamma_0^5 t \sin(3\theta - t) - \frac{14}{3} \cos(2\theta - 2t) \alpha^2 \gamma_0^2 \right. \\
 & \left. + 4 \gamma_0^2 \alpha^2 \sin(2\theta - 2t) t + \frac{4}{3} \gamma_0^2 \alpha^2 \cos(2\theta - 2t) t^2 + \frac{65}{108} \gamma_0^5 \sin(3\theta - 3t) t + \frac{25}{216} \gamma_0^5 \cos(\theta - t) t^2 \right. \\
 & \left. - \frac{91}{1296} \cos(5\theta - 5t) \gamma_0^5 - \frac{2}{3} \cos(2\theta) \alpha^2 \gamma_0^2 + \frac{16}{3} \cos(2\theta - t) \alpha^2 \gamma_0^2 - \frac{205}{216} \cos(3\theta - t) \gamma_0^5 \right)
 \end{aligned}$$

Four time scales

- Rotation with cylinder $1/\omega$
- Surface tension makes envelope cylindrical $\mu a^4/\sigma h^3$
- Slow drift of rotation $\omega \mu^2 a^2 / \rho^2 g^2 h^4$
- Slow decay of wobble $\omega^2 \mu^3 a^6 / \rho^2 g^2 \sigma h^7$.

K: ‘Why weeny water waves wobble, wane & wander whilst whirling’

Nondimensionalise

To solve

$$h_t + h_\theta - \gamma(h^3 \cos \theta)_\theta + \alpha(h^3(h_{\theta\theta} + h)_\theta)_\theta = 0$$

with $h(\theta, 0) = 1$.

gravity $\gamma = \frac{\rho g h^2}{3\mu\omega a}$

surface tension $\alpha = \frac{\sigma h^3}{3\mu\omega a^4}$

Multiple scales 1 (easy)

To solve

$$h_t + h_\theta - \gamma(h^3 \cos \theta)_\theta + \alpha(h^3(h_{\theta\theta} + h)_\theta)_\theta = 0$$

with $h(\theta, 0) = 1$, for $\gamma \ll 1$ with $\alpha = O(1)$ small.

Expand

$$h \sim 1 + \gamma h_1 + \gamma^2 h_2 + \gamma^3 h_3$$

with

$$h_n(\theta, \tau = t, T = \gamma^2 t)$$

Multiple scales 1, $O(\gamma)$

$O(\gamma)$:

$$h_{1\tau} + h_{1\theta} + \alpha(h_{1\theta\theta\theta\theta} + h_{1\theta\theta}) = -\sin \theta$$

Solution

$$\begin{aligned} h_1 &= \cos \theta + A(T) \cos(\theta - \tau) + B(T) \sin(\theta - \tau) \\ &+ \sum_{n>1} [A_n \cos n(\theta - \tau) + B_n \sin n(\theta - \tau)] e^{-\alpha n^2(n^2-1)\tau} \end{aligned}$$

with $A(0) = -1$ and $B(0) = 0$.

This is steady solution + eccentrically rotating circle + surface tension squeezing to a circle.

Multiple scales 1, $O(\gamma^2)$

$O(\gamma^2)$:

$$h_{2\tau} + h_{2\theta} + \alpha(h_{2\theta\theta\theta\theta} + h_{2\theta\theta}) = -(3h_1 \cos \theta)_\theta$$

Solution – a mess

$$\begin{aligned} h_2 = & \frac{3}{2(1 + 36\alpha^2)} \cos 2\theta - \frac{9\alpha}{1 + 36\alpha^2} \sin 2\theta \\ & + \frac{3(A + 12\alpha B)}{1 + 144\alpha^2} \cos(2\theta - \tau) + \frac{3(B - 12\alpha A)}{1 + 144\alpha^2} \sin(2\theta - \tau) \end{aligned}$$

after others decay by surface tension.

Multiple scales 1, $O(\gamma^3)$

$O(\gamma^3)$:

$$\begin{aligned} & h_{3\tau} + h_{3\theta} + \alpha(h_{3\theta\theta\theta\theta} + h_{3\theta\theta}) \\ &= -h_{1T} - ((3h_2 + 3h_1^2) \cos \theta)_\theta - \alpha(3h_1(h_{2\theta\theta} + h_2)_\theta)_\theta \end{aligned}$$

Secularity condition – resonant $\sin(\theta - \tau)$ and $\cos(\theta - \tau)$

$$\frac{d}{dT} \begin{pmatrix} A \\ B \end{pmatrix} = \begin{pmatrix} -\textcolor{red}{a} & -b \\ b & -\textcolor{red}{a} \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

where

$$a = 81\alpha/(1 + 144\alpha^2) \quad \text{and} \quad b = 3(5 + 72\alpha^2)/2(1 + 144\alpha^2)$$

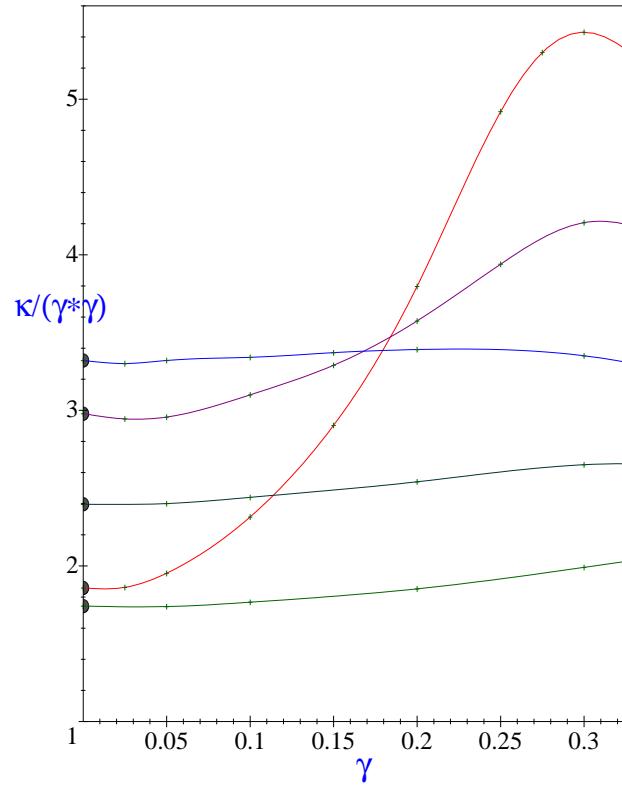
Multiple scales 1, decay rate

Hence decay rate

$$\kappa = 81\gamma^2\alpha/(1 + 144\alpha^2)$$

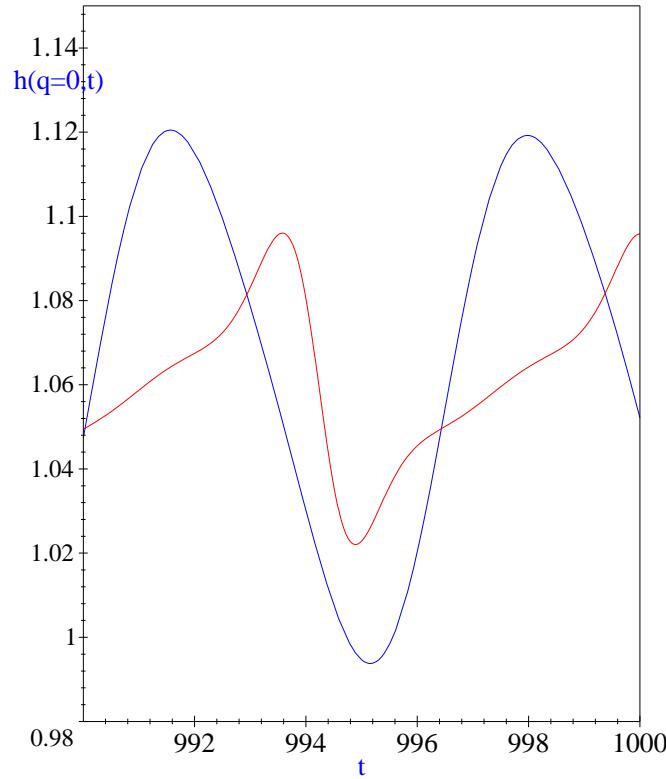
$$\alpha = 0.025, 0.05, 0.1, 0.2, 0.3$$

$$[\gamma^2\alpha = \omega^2\mu^3a^6/\rho^2g^2\sigma h^7]$$



- Maximum load becomes steady in 3 revolutions, half-maximum in 300 and quarter-maximum in 3000, because $\kappa \propto h^{-7}$.

M's problem

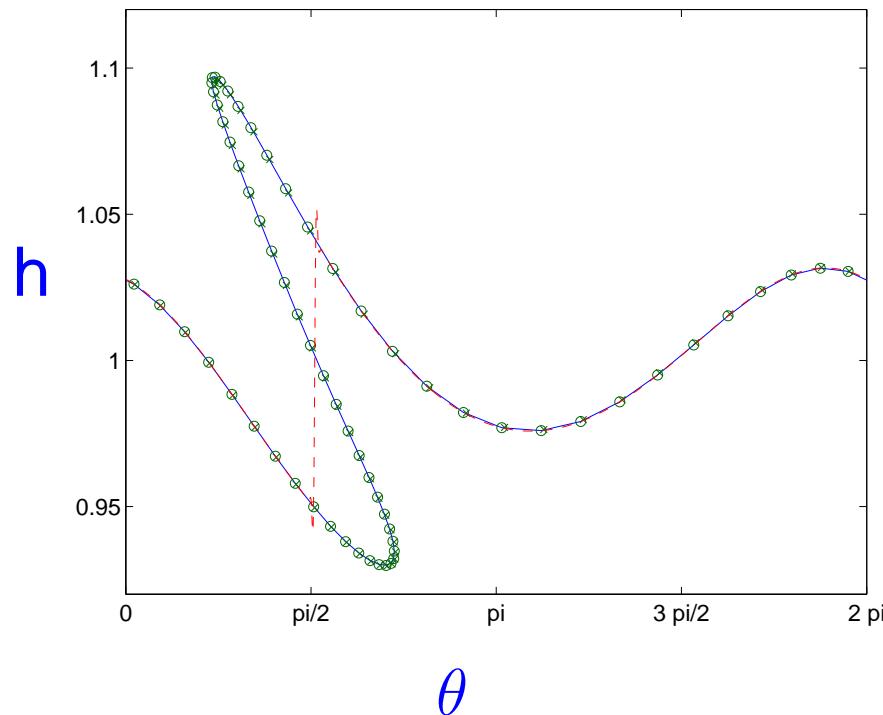


$h(\theta = 0, t)$ for
 $\gamma = 5.32 \cdot 10^{-2}$, $\alpha = 4.8 \cdot 10^{-5}$

Profile not
theoretical $\cos + \cos 2$
but shocking!

First study valid for $\gamma^2 \ll \alpha, \gamma \ll 1$ What if $\alpha \ll \gamma^2$?

No surface tension



$t = 500$

$\alpha = 0$ ——— numerical, \circ , \times action-angle integration
 $\alpha = 10^{-9}$ - - - - numerical (spectral)

Multiple scales 2

To solve

$$h_t + h_\theta - \gamma(h^3 \cos \theta)_\theta + \alpha(h^3(h_{\theta\theta} + h)_\theta)_\theta$$

for $\gamma \ll 1$ and $\beta = \alpha/30\gamma^3 = O(1)$.

Expand

$$h \sim 1 + \gamma h_1 + \gamma^2 h_2 + \gamma^3 h_3 + \gamma^4 h_4$$

with $h_n(\theta, \tau, T)$

$$\tau = t \left(1 - \frac{15}{2}\gamma^2 - \frac{1215}{8}\gamma^4 \right) \quad \text{and} \quad T = 30\gamma^3 t$$

Multiple scales 2, $O(\gamma)$

$O(\gamma)$:

$$h_{1\tau} + h_{1\theta} = -\sin \theta$$

Solution

$$h_1 = \cos \theta + \textcolor{red}{F}(\eta = \theta - \tau, T)$$

with

$$\textcolor{red}{F}(\eta, 0) = -\cos \eta$$

Disturbance rotating with cylinder. To find evolution.

Multiple scales 2, $O(\gamma^2)$

$O(\gamma^2)$:

$$h_{2\tau} + h_{2\theta} = (3h_1 \cos \theta)_\theta$$

Solution

$$\begin{aligned} h_2 = & 3F \cos \theta + 3F' \sin \theta \\ & + \text{steady state} + G(\eta, T) \end{aligned}$$

Multiple scales 2, $O(\gamma^3)$

$O(\gamma^3)$:

$$h_{3\tau} + h_{3\theta} = \left(3(h_2 + h_1^2) \cos \theta\right)_\theta$$

Solution

$$\begin{aligned} h_3 = & \frac{15}{2} F \cos 2\theta + 3F^2 \cos \theta + \frac{33}{4} F' \sin \theta - \frac{9}{4} F'' \cos 2\theta \\ & + \text{steady state} + \text{terms G} + H(\eta, T) \end{aligned}$$

Multiple scales 2, $O(\gamma^4)$

$O(\gamma^3)$:

$$\begin{aligned} h_{4\tau} + h_{4\theta} = & -h_{1T} - 30\beta(h_{1\theta\theta} + h_1)_{\theta\theta} \\ & + ((3h_3 + 6h_2h_1 + h_1^3)\cos\theta)_\theta \end{aligned}$$

Secularity : Kuramoto-Sivashinsky equation

$$F_T - FF_\eta + \beta(F_{\eta\eta\eta\eta} + F_{\eta\eta}) = 0$$

with

$$F(\eta, 0) = -\cos\eta$$

Governing equation now has constant coefficients

KS with $\beta \ll 1$

$$F_T - FF_\eta + \beta(F_{\eta\eta\eta\eta} + F_{\eta\eta}) = 0 \quad \text{with} \quad F(\eta, 0) = -\cos \eta$$

Set $\beta = 0$.

Shock forms at

$$T = 1 \quad \text{where} \quad \eta = \frac{\pi}{2}$$

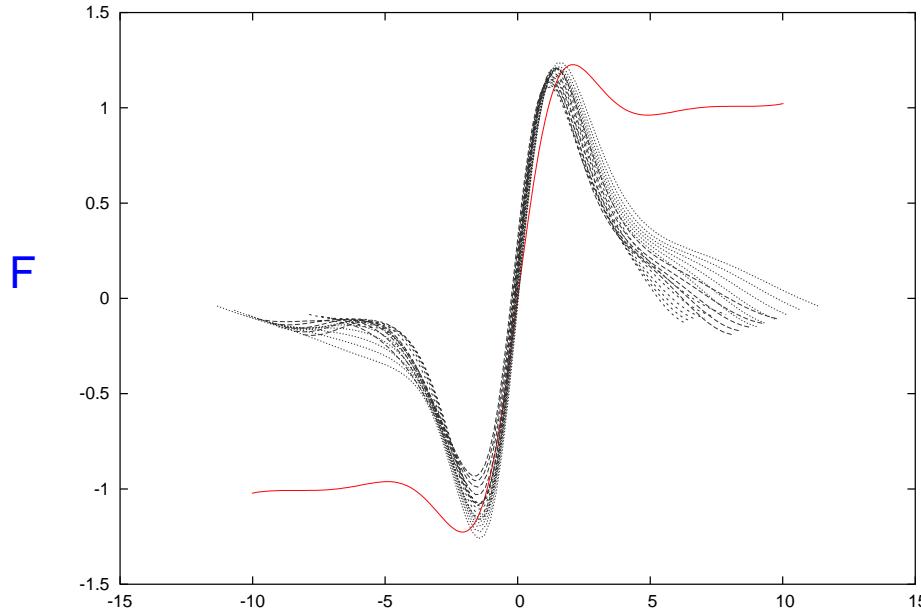
Afterwards, decay of a N-shock wave

$$F \sim \frac{\frac{\pi}{2} + \eta}{T}$$

Structure of the shock

Quasi-steady, with $F \rightarrow \pm A(T)$ outside, $A = \pi/T$

$$F^2 - A^2 = 2\beta F'''$$



at different times.
for $\gamma = 0.1, 0.08$
and $\alpha = 10^{-4}, 3 \cdot 10^{-4}$

$$(\eta - \frac{\pi}{2}) / (\mathcal{A}/2\beta)^{1/3}$$

But width $(2\beta/A)^{1/3}$ no longer thin when $T = 1/\beta$.

KS with $\beta \gg 1$

Multiple scales 3. Slow-slow time $\textcolor{red}{T} = \beta^{-1} T$

$$F \sim A(\textcolor{red}{T}) \cos \eta - \frac{1}{24} \beta^{-1} A^2 \sin 2\eta$$

$O(\beta^{-2})$: secularity

$$A_{\text{T}} = \frac{1}{48} A^3 \quad \text{with} \quad A(0) = -1$$

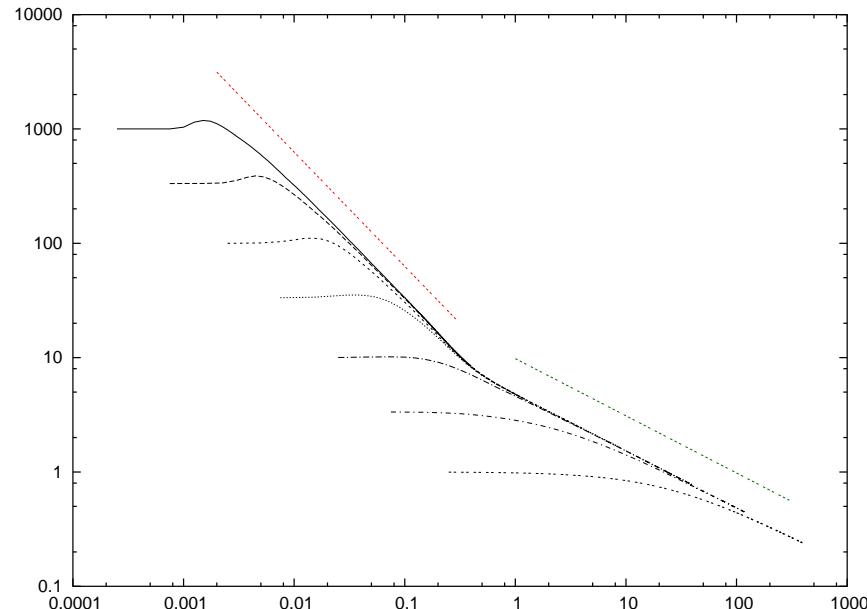
Solution

$$A = - \left(1 + \frac{\textcolor{red}{T}}{24} \right)^{-1/2}$$

Can also use for $\beta \ll 1$ once $\beta T \gg 1$.

Decay of KS solutions

$\frac{F_{\max}}{\beta}$



for 7 β : 0.003 to 1

Decay: T^{-1} , then $T^{-1/2}$, and then ...

...and then

After algebraic decay of KS equation, eventually
exponential decay of lubrication equation

$$F_{\max} \sim \frac{\alpha}{\left(\alpha^2 + \frac{25}{108}\gamma^4\right)^{1/2}} e^{-81\alpha\gamma^2 t}$$

NB: first study needed $\alpha \gg \gamma^2$.

Conclusions

- Transients decay
 - $\gamma^2 \ll \alpha, \gamma \ll 1$: exponentially
 - $\gamma^3 \ll \alpha \ll \gamma^2$: $T^{-1/2} \rightarrow$ exponentially
 - $\gamma^4 \ll \alpha \ll \gamma^3$: $T^{-1} \rightarrow T^{-1/2} \rightarrow$ exponentially
- HKM's steady state much simpler!