Symmetry and Particle Physics, 2

1. Let $T^{\alpha_1...\alpha_{2j}}$ be a symmetric SU(2) tensor for $j=\frac{1}{2},1,\frac{3}{2},\ldots$ Show that the action of the spin operator is given by

$$\mathbf{S}^{\alpha_1 \dots \alpha_{2j}}{}_{\beta_1 \dots \beta_{2j}} T^{\beta_1 \dots \beta_{2j}} = \sum_{i=1}^{2j} \frac{1}{2} (\boldsymbol{\sigma})^{\alpha_i}{}_{\beta} T^{\alpha_1 \dots \alpha_{i-1} \beta \alpha_{i+1} \dots \alpha_{2j}} ,$$

where σ are the Pauli matrices. Define for $m = -j, -j + 1, \dots, j$

$$T^{(jm)} = \left[(j+m)!(j-m)! \right]^{-\frac{1}{2}} T \underbrace{1...1}_{j+m} \underbrace{2...2}_{j-m}.$$

Calculate $S_{\pm}T^{(jm)}$ and $S_3T^{(jm)}$. Show that $\bar{T}_{\alpha_1...\alpha_{2j}}T^{\alpha_1...\alpha_{2j}}=(2j)!\sum_m T^{(jm)*}T^{(jm)}$ for $\bar{T}_{\alpha_1...\alpha_{2j}}$ the conjugate tensor.

2. A field $\phi(x)$ transforms under the action of a Poincaré transformation (Λ, a) such that $U[\Lambda, a]\phi(x)U[\Lambda, a]^{-1} = \phi(\Lambda x + a)$. For an infinitesimal transformation, $\Lambda^{\mu}{}_{\nu} = \delta^{\mu}{}_{\nu} + \omega^{\mu}{}_{\nu}$ and correspondingly $U[\Lambda, a] = 1 - i \frac{1}{2} \omega^{\mu\nu} M_{\mu\nu} - i a^{\mu} P_{\mu}$ show that

$$[M_{\mu\nu}, \phi(x)] = -i(x_{\mu}\partial_{\nu} - x_{\nu}\partial_{\mu})\phi(x), \qquad [P_{\mu}, \phi(x)] = i\partial_{\mu}\phi(x).$$

Verify that $M_{\mu\nu} \to i(x_{\mu}\partial_{\nu} - x_{\nu}\partial_{\mu})$ and $P_{\mu} \to -i\partial_{\mu}$ satisfy the algebra for $[M_{\mu\nu}, M_{\sigma\rho}]$ and $[M_{\mu\nu}, P_{\sigma}]$ expected for the Poincaré group.

3. Show how $B(\theta, \mathbf{n}) \in Sl(2, \mathbb{C})$ where

$$B(\theta,\mathbf{n}) = \cosh \tfrac{1}{2}\theta + \boldsymbol{\sigma} \cdot \mathbf{n} \; \sinh \tfrac{1}{2}\theta \,, \qquad \mathbf{n}^2 = 1 \,,$$

corresponds to a Lorentz boost with velocity $\mathbf{v} = \tanh \theta \mathbf{n}$. Show that

$$(1 + \frac{1}{2}\boldsymbol{\sigma} \cdot \delta \mathbf{v})B(\theta, \mathbf{n}) = B(\theta', \mathbf{n}')R,$$

where, to first order in $\delta \mathbf{v}$,

$$\theta' = \theta + \delta \mathbf{v} \cdot \mathbf{n}$$
, $\mathbf{n}' = \mathbf{n} + \coth \theta (\delta \mathbf{v} - \mathbf{n} \ \mathbf{n} \cdot \delta \mathbf{v})$,

and R is an infinitesimal rotation given by

$$R = 1 + \tanh \frac{1}{2}\theta \, \, \frac{1}{2}i \left(\delta \mathbf{v} \times \mathbf{n} \right) \cdot \boldsymbol{\sigma} = 1 + \frac{\gamma}{\gamma + 1} \, \, \frac{1}{2}i \left(\delta \mathbf{v} \times \mathbf{v} \right) \cdot \boldsymbol{\sigma} \,, \qquad \gamma = (1 - \mathbf{v}^2)^{-\frac{1}{2}} \,.$$

Show that we must have $\mathbf{v}' = \mathbf{v} + \delta \mathbf{v} - \mathbf{v} \ \mathbf{v} \cdot \delta \mathbf{v}$. [Note $\boldsymbol{\sigma} \cdot \mathbf{a} \ \boldsymbol{\sigma} \cdot \mathbf{b} = \mathbf{a} \cdot \mathbf{b} \ 1 + i \boldsymbol{\sigma} \cdot (\mathbf{a} \times \mathbf{b})$.]

4. The group of four dimensional space-time symmetries may be generalised to conformal transformations $x \to x'$ defined by the requirement

$$\mathrm{d}x'^2 = \Omega(x)^2 \mathrm{d}x^2 \,,$$

where $dx^2 = g_{\mu\nu} dx^{\mu} dx^{\nu}$. For an infinitesimal transformation $x'^{\mu} = x^{\mu} + f^{\mu}(x)$, $\Omega(x)^2 = 1 + 2\sigma(x)$. Obtain in this case

$$\partial_{\mu} f_{\nu} + \partial_{\nu} f_{\mu} = 2\sigma g_{\mu\nu} \quad \Rightarrow \quad 4\sigma = \partial \cdot f \,.$$

Hence obtain

$$4 \,\partial_{\sigma} \partial_{\mu} f_{\nu} = g_{\mu\nu} \,\partial_{\sigma} \partial \cdot f + g_{\sigma\nu} \,\partial_{\mu} \partial \cdot f - g_{\sigma\mu} \,\partial_{\nu} \partial \cdot f \,.$$

From this obtain $2\partial_{\sigma}\partial_{\mu}\partial \cdot f = -g_{\sigma\mu}\partial^{2}\partial \cdot f$ and hence show that we must have $\partial_{\sigma}\partial_{\mu}\partial \cdot f = 0$. Why does it then follow that $f_{\mu}(x)$ can only be quadratic in x? Show that $f^{\mu}(x)$ must then have the general form

$$f^{\mu}(x) = a^{\mu} + \omega^{\mu}_{\ \nu} x^{\nu} + \lambda x^{\mu} + b^{\mu} x^2 - 2b \cdot x \, x^{\mu} \,, \quad \omega_{\mu\nu} = -\omega_{\nu\mu} \,.$$

Show also that an inversion $x'^{\mu} = x^{\mu}/x^2$ is a conformal transformation. Calculate the finite conformal transformation obtained by an inversion followed by a translation by b^{μ} followed by another inversion and show that it is compatible with the result for $f^{\mu}(x)$.

5. A four-dimensional space is defined in terms of 6-vectors $\eta^A = (\eta^\mu, \eta^+, \eta^-), \mu = 0, 1, 2, 3,$ subject to the relations

$$\eta \cdot \eta = g_{AB} \eta^A \eta^B = g_{\mu\nu} \eta^\mu \eta^\nu - \eta^+ \eta^- = 0, \qquad \eta^A \sim C \eta^A.$$

Using $\eta^{\pm} = \eta^4 \pm \eta^5$ show that this space is invariant under transformations $\eta^A \to G^A{}_B \eta^B$ where $[G^A{}_B]$ are matrices belonging to the group SO(4,2). For an infinitesimal transformation, $G^A{}_B = \delta^A{}_B + \omega^A{}_B$, show that $\omega^A{}_B$ may be decomposed in the form

$$\begin{bmatrix} \omega^A{}_B \end{bmatrix} = \begin{pmatrix} \omega^\mu{}_\nu & a^\mu & b^\mu \\ 2b_\nu & -\lambda & 0 \\ 2a_\nu & 0 & \lambda \end{pmatrix}, \qquad \omega_{\mu\nu} = -\omega_{\nu\mu}.$$

Suppose, for $\eta^+ \neq 0$, $\eta^A = \eta^+(x^\mu, 1, x^2)$. Using $\delta \eta^A = \omega^A{}_B \eta^B$ determine the corresponding δx^μ . What transformation corresponds to $\eta^+ \leftrightarrow \eta^-$? For four points x_i , i = 1, 2, 3, 4, calculate $\eta_1 \cdot \eta_2 \, \eta_3 \cdot \eta_4 / (\eta_1 \cdot \eta_3 \, \eta_2 \cdot \eta_4)$. Why is this a conformal invariant?

6. Consider the subgroup of the Galilean group corresponding to translations and boosts where

$$t' = t + b$$
, $\mathbf{x}' = \mathbf{x} + \mathbf{a} + \mathbf{v}t$.

Denoting the corresponding group element by $(b, \mathbf{a}, \mathbf{v})$ work out the group multiplication law and show that $(b_2, \mathbf{a}_2, \mathbf{v}_2)^{-1}(b_1, \mathbf{a}_1, \mathbf{v}_1)^{-1}(b_2, \mathbf{a}_2, \mathbf{v}_2)(b_1, \mathbf{a}_1, \mathbf{v}_1) = (0, 0, b_1\mathbf{v}_2 - b_2\mathbf{v}_1)$. Suppose, for infinitesimal $b, \mathbf{a}, \mathbf{v}$, the associated unitary operator has the form

$$U[b, \mathbf{a}, \mathbf{v}] = 1 + ib H - i\mathbf{a} \cdot \mathbf{P} + i\mathbf{v} \cdot \mathbf{K},$$

work out the corresponding commutators. For general \mathbf{a}, \mathbf{v} require $U[0, \mathbf{a}, \mathbf{v}] = T[\mathbf{a}]U_B[\mathbf{v}]$ and assume now $T[\mathbf{a}]U_B[\mathbf{v}] = e^{im \mathbf{a} \cdot \mathbf{v}} U_B[\mathbf{v}] T[\mathbf{a}]$ for some positive constant m. Show that this leads to the modified commutation relation $[K_i, P_i] = im \delta_{ij}$.

Suppose $\mathbf{P}|\mathbf{0}\rangle = \mathbf{0}$ and define $|m\mathbf{v}\rangle = U_B[\mathbf{v}]|\mathbf{0}\rangle$. Show that is an eigenvector of \mathbf{P} and that $H|\mathbf{p}\rangle = (E_0 + \frac{\mathbf{p}^2}{2m})|\mathbf{p}\rangle$.