



Wind Forcing



Kallberg et al. 1995 (ERA-40 reanalysis)

Ocean Currents

Time-averaged (16-year) ocean circulation







Numerical Simulations:



Haidvogel, McWilliams, Gent, J. Phys. Oceanography 1992

Instantaneous ψ

Haidvogel, McWilliams, Gent, J. Phys. Oceanography 1992

Sea-surface temperature





2007 Mean net surface heat flux (W/m²), Sea surface temperature (contours)



- Surface flow
- Deep flow
- Bottom flow
- Deep Water Formation
- Wind-driven upwelling
- Mixing-driven upwelling
 - Salinity > 36 ‰
 - Salinity < 34 ‰

- Labrador Sea
- Nordic Seas
- W Weddell Sea
 - Ross Sea

L

Ν

R



The surface sources of global ocean waters. Oceanic volume that has originated in each 2° by 2° surface location (11,113 origination sites), scaled by the surface area of each box to make an equivalent thickness, d. The color-scale follows a base ten logarithm of the field.

eWOCE

Potential Temperature θ [°C]









eWOCE

Salinity [pss-78]





Silicate [µmol/kg]

Silicate [µmol/kg]







eWOCE

Depth [m]

6000

40°S

20°S

Nitrate [µmol/kg]





EQ

20°N

40°N

60°N



Phosphate [µmol/kg]

1.5

Silicate [µmol/kg] 150 σ_0 125 Potential Temperature 8 [°C] 20 100 75 10 50 25 36 37 34 35 Salinity [pss-78]

Averaging, across an ocean in longitude

 Meridional Overturning Circulation (MOC) (sometimes called the Thermohaline Circulation, THC)

Atlantic MOC (1 Sv = $10^6 \text{ m}^3/\text{s}$)







Mignotti + Woods, 2015





'Missing' Mixing

Problem: Observations of ocean mixing typically find $\kappa \approx 10^{-5} m^2/s \ll 10^{-4} m^2/s$



Evidence for slow mixing across the pycnocline from an open-ocean tracer-release experiment

James R. Ledwell*, Andrew J. Watson† & Clifford S. Law†

* Applied Ocean Physics and Engineering,
Woods Hole Oceanographic Institution, Woods Hole.
Massachusetts 02543, USA
† Plymouth Marine Laboratory, Prospect Place, West Hoe,
Plymouth PL1 3DH, UK

THE distributions of heat, salt and trace substances in the ocean thermocline depend on mixing along and across surfaces of equal density (isopycnal and diapycnal mixing, respectively). Measurements of the invasion of anthropogenic tracers, such as bomb tritium and ³He (see, for example, refs 1 and 2), have indicated that isopycnal processes dominate diapycnal mixing, and turbulence measurements have suggested that diapycnal mixing is small^{3,4} but it has not been possible to measure accurately the diapycnal diffusivity. Here we report such a measurement, obtained from the vertical dispersal of a patch of the inert compound SF₆ released in the open ocean. The diapycnal diffusivity, averaged over hundreds of kilometres and five months, was 0.11 ± 0.02 cm² s⁻¹, confirming previous estimates¹⁻⁴. Such a low diffusivity can support only a rather small diapycnal flux of nitrate into the euphotic zone; it justifies the neglect of diapycnal mixing in dynamic models of the thermocline²⁵⁻²⁷, and implies that heat, salt and tracers must penetrate the thermocline mostly by transport along, rather than across, density surfaces.



FIG. 1 Evolution of the lateral distribution of the tracer. The injection streaks are shown as short heavy lines near 26° N, 28° W. The contours just to the west show the patch later in May 1992. Heavy lines (further to the west) show tracks for the October survey, where the concentration C at the target surface was >500 fM; light solid lines, C was between 100 and 500 fM; dashed lines, C ~0. Solid triangles indicate bottle stations occupied at the end of the October cruise, with C>300 fM. Station symbols for the November survey are: plus signs, C <30 fM; open circles, C = 30–300 fM; filled circles, C > 300 fM. A fine curve has been drawn to envelop the high C regions for the two surveys. CM marks the location of the central mooring for the Subduction experiment.



FIG. 2 Evolution of the vertical distribution of the tracer. The mean profiles have been scaled so that the widths can be compared.



FIG. 3 Growth of the second moment of the vertical tracer distribution. Squares are for raw M_2 , circles are for the centre of mass shifted to h=0. The line is for $K_z=0.11$ cm² s⁻¹.

Mixing Efficiency

Inoue + Smyth 2009







Boundary Mixing



Armi, 1972



Horizontal convection

Hughes and Griffiths 2008



FIG. 4. Dependence of the time-averaged overturning circulation upon the vertical diffusion coefficient (surface buoyancy flux is fixed, with $Q_0 = 200 \text{ W m}^{-2}$). The maximum streamfunction quoted is that for a two-dimensional flow in a basin of 1-m width, while the density range is $\Delta \rho = \overline{\rho}_{\text{bottom}} - \overline{\rho}_{\text{top}}$. The 20°C isotherm is shown in gray.



Wunsch and Ferrari, 2004



Video: Ryan Abernathy



Matt Colebrook







Figure 1 Changes in surface air temperature caused by a shutdown of North Atlantic Deep Water (NADW) formation in a current ocean-atmosphere circulation model. Note the hemispheric seesaw (Northern Hemisphere cools while the Southern Hemisphere warms) and the maximum cooling over the northern Atlantic. In this particular model (HadCM3)⁷, the surface cooling resulting from switching off NADW formation is up to 6 °C. It is further to the west compared with most models, which tend to put the maximum cooling near Scandinavia. This probably depends on the exact location of deep-water formation (an aspect not well represented in current



coarse-resolution models) and on the sea-ice distribution in the models, as ice-margin shifts act to amplify the cooling. The largest air temperature cooling is thus greater than the largest sea surface temperature (SST) cooling. The latter is typically around 5 °C and roughly corresponds to the observed SST difference between the northern Atlantic and Pacific at a given latitude. In most models, maximum air temperature cooling ranges from 6 °C to 11 °C in annual mean; the effect is generally stronger in winter.

Rahmstorf, Nature, 2002

