

Residual Dynamic Mode Decomposition

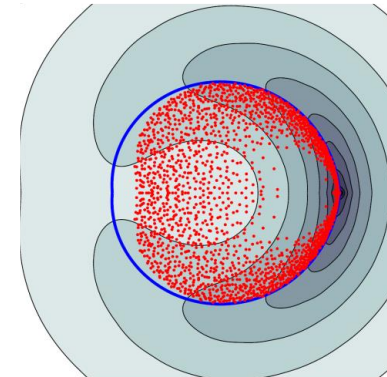
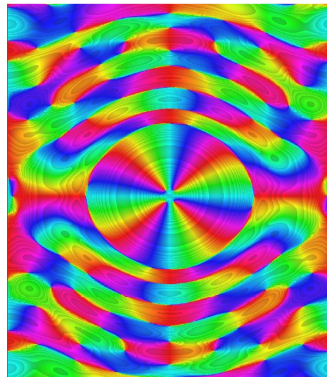
Rigorous data-driven computation of spectral properties
of Koopman operators for dynamical systems

Matthew Colbrook (m.colbrook@damtp.cam.ac.uk)

University of Cambridge

Alex Townsend (townsend@cornell.edu)

Cornell University

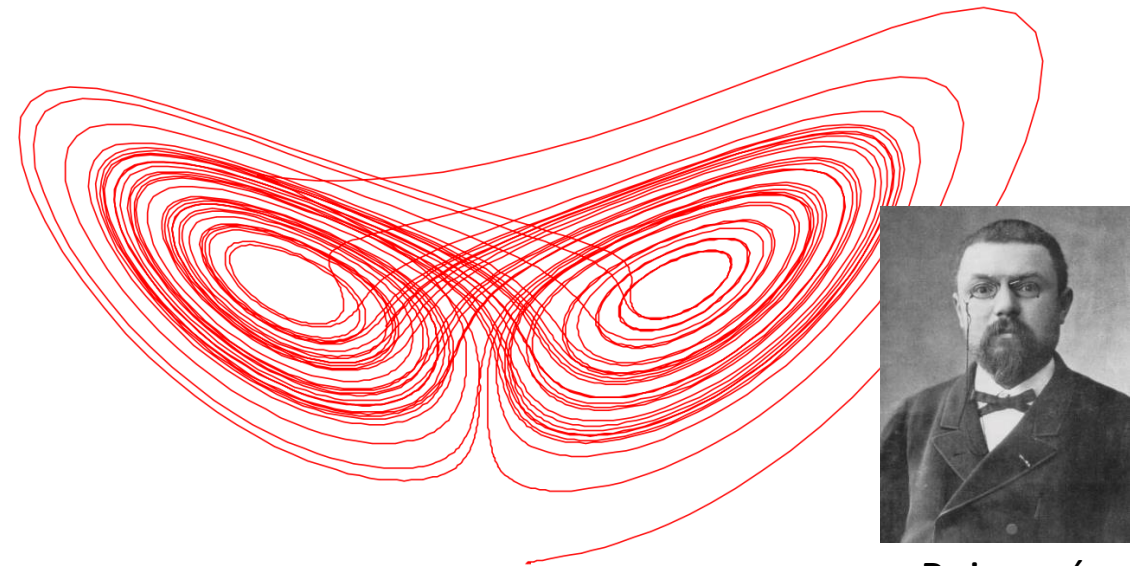


Data-driven dynamical systems

- State $x \in \Omega \subseteq \mathbb{R}^d$, **unknown** function $F: \Omega \rightarrow \Omega$ governs dynamics

$$x_{n+1} = F(x_n)$$

- **Goal:** Learn about system from data $\{x^{(m)}, y^{(m)} = F(x^{(m)})\}_{m=1}^M$
 - **Data:** experimental measurements or numerical simulations
 - E.g., **used for** forecasting, control, design, understanding
- **Applications:** chemistry, climatology, electronics, epidemiology, finance, fluids, molecular dynamics, neuroscience, plasmas, robotics, video processing, etc.



Poincaré

Operator viewpoint

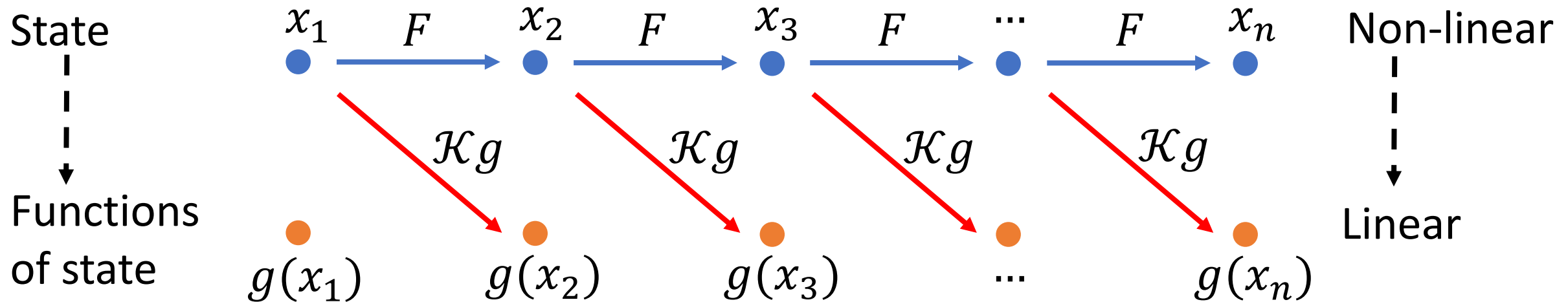
Koopman

von Neumann



- **Koopman operator** \mathcal{K} acts on functions $g: \Omega \rightarrow \mathbb{C}$

$$[\mathcal{K}g](x_n) = g(F(x_n)) = g(x_{n+1})$$
- \mathcal{K} is **linear** but acts on an **infinite-dimensional** space.



- Work in $L^2(\Omega, \omega)$ for positive measure ω , with inner product $\langle \cdot, \cdot \rangle$.

• Koopman, "Hamiltonian systems and transformation in Hilbert space," *Proc. Natl. Acad. Sci. USA*, 1931.

• Koopman, v. Neumann, "Dynamical systems of continuous spectra," *Proc. Natl. Acad. Sci. USA*, 1932.

Why is linear (much) easier?

$$x_{n+1} = F(x_n)$$

- Suppose $F(x) = Ax, A \in \mathbb{R}^{d \times d}, A = V\Lambda V^{-1}$.
- Set $\xi = V^{-1}x$,

$$\xi_n = V^{-1}x_n = V^{-1}A^n x_0 = \Lambda^n V^{-1}x_0 = \Lambda^n \xi_0$$

- Let $w^T A = \lambda w$, set $\varphi(x) = w^T x$,

$$[\mathcal{K}\varphi](x) = w^T Ax = \lambda \varphi(x)$$

Long-time dynamics
become trivial!



Eigenfunction

Much more general (**non-linear** and even **chaotic** F).

Koopman mode decomposition

generalised
eigenfunction of \mathcal{K}

eigenfunction of \mathcal{K}

$$g(x) = \sum_{\text{eigs } \lambda_j} c_{\lambda_j} \varphi_{\lambda_j}(x) + \int_{[-\pi, \pi]_{\text{per}}} \phi_{\theta, g}(x) d\theta$$

$$g(x_n) = [\mathcal{K}^n g](x_0) = \sum_{\text{eigs } \lambda_j} c_{\lambda_j} \lambda_j^n \varphi_{\lambda_j}(x_0) + \int_{[-\pi, \pi]_{\text{per}}} e^{in\theta} \phi_{\theta, g}(x_0) d\theta$$

$$[\mathcal{K}g](x) = g(F(x))$$

Encodes: geometric features, invariant measures, transient behaviour, long-time behaviour, coherent structures, quasiperiodicity, etc.

GOAL: Data-driven approximation of \mathcal{K} and its spectral properties.

- Mezić, “Spectral properties of dynamical systems, model reduction and decompositions,” **Nonlinear Dynam.**, 2005.

Challenges of computing

$$\text{Spec}(\mathcal{K}) = \{\lambda \in \mathbb{C}: \mathcal{K} - \lambda I \text{ is not invertible}\}$$

Truncate: $\mathcal{K} \longrightarrow \mathbb{K} \in \mathbb{C}^{N_K \times N_K}$

- 1) **“Too much”:** Approximate spurious modes $\lambda \notin \text{Spec}(\mathcal{K})$
- 2) **“Too little”:** Miss parts of $\text{Spec}(\mathcal{K})$
- 3) **Continuous spectra.**

Verification: Is it right?

Build the matrix: Dynamic Mode Decomposition (DMD)

Given dictionary $\{\psi_1, \dots, \psi_{N_K}\}$ of functions $\psi_j: \Omega \rightarrow \mathbb{C}$,

$$\{x^{(m)}, y^{(m)} = F(x^{(m)})\}_{m=1}^M$$

$$\langle \psi_k, \psi_j \rangle \approx \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \psi_k(x^{(m)}) = \left[\underbrace{\begin{pmatrix} \psi_1(x^{(1)}) & \dots & \psi_{N_K}(x^{(1)}) \\ \vdots & \ddots & \vdots \\ \psi_1(x^{(M)}) & \dots & \psi_{N_K}(x^{(M)}) \end{pmatrix}}_{\Psi_X} \underbrace{\begin{pmatrix} w_1 & & \\ & \ddots & \\ & & w_M \end{pmatrix}}_W \underbrace{\begin{pmatrix} \psi_1(x^{(1)}) & \dots & \psi_{N_K}(x^{(1)}) \\ \vdots & \ddots & \vdots \\ \psi_1(x^{(M)}) & \dots & \psi_{N_K}(x^{(M)}) \end{pmatrix}}_{\Psi_X} \right]_{jk}$$

$$\langle \mathcal{K}\psi_k, \psi_j \rangle \approx \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \underbrace{\psi_k(y^{(m)})}_{[\mathcal{K}\psi_k](x^{(m)})} = \left[\underbrace{\begin{pmatrix} \psi_1(x^{(1)}) & \dots & \psi_{N_K}(x^{(1)}) \\ \vdots & \ddots & \vdots \\ \psi_1(x^{(M)}) & \dots & \psi_{N_K}(x^{(M)}) \end{pmatrix}}_{\Psi_X} \underbrace{\begin{pmatrix} w_1 & & \\ & \ddots & \\ & & w_M \end{pmatrix}}_W \underbrace{\begin{pmatrix} \psi_1(y^{(1)}) & \dots & \psi_{N_K}(y^{(1)}) \\ \vdots & \ddots & \vdots \\ \psi_1(y^{(M)}) & \dots & \psi_{N_K}(y^{(M)}) \end{pmatrix}}_{\Psi_Y} \right]_{jk}$$

$$\mathcal{K} \longrightarrow \mathbb{K} = (\Psi_X^* W \Psi_X)^{-1} \Psi_X^* W \Psi_Y \in \mathbb{C}^{N_K \times N_K}$$

Recall open problems: too much, too little, continuous spectra, verification

- Schmid, “Dynamic mode decomposition of numerical and experimental data,” **J. Fluid Mech.**, 2010.
- Rowley, Mezić, Bagheri, Schlatter, Henningson, “Spectral analysis of nonlinear flows,” **J. Fluid Mech.**, 2009.
- Kutz, Brunton, Brunton, Proctor, “Dynamic mode decomposition: data-driven modeling of complex systems,” **SIAM**, 2016.
- Williams, Kevrekidis, Rowley “A data-driven approximation of the Koopman operator: Extending dynamic mode decomposition,” **J. Nonlinear Sci.**, 2015.

Residual DMD (ResDMD): Approx. \mathcal{K} and $\mathcal{K}^*\mathcal{K}$

$$\langle \psi_k, \psi_j \rangle \approx \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \psi_k(x^{(m)}) = \left[\underbrace{\Psi_X^* W \Psi_X}_G \right]_{jk}$$

$$\langle \mathcal{K}\psi_k, \psi_j \rangle \approx \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \underbrace{\psi_k(y^{(m)})}_{[\mathcal{K}\psi_k](x^{(m)})} = \left[\underbrace{\Psi_X^* W \Psi_Y}_{K_1} \right]_{jk}$$

$$\langle \mathcal{K}\psi_k, \mathcal{K}\psi_j \rangle \approx \sum_{m=1}^M w_m \overline{\psi_j(y^{(m)})} \psi_k(y^{(m)}) = \left[\underbrace{\Psi_Y^* W \Psi_Y}_{K_2} \right]_{jk}$$

Residuals: $g = \sum_{j=1}^{N_K} \mathbf{g}_j \psi_j$, $\|\mathcal{K}g - \lambda g\|^2 \approx \mathbf{g}^* [K_2 - \lambda K_1^* - \bar{\lambda} K_1 + |\lambda|^2 G] \mathbf{g}$

-
- C., T., “Rigorous data-driven computation of spectral properties of Koopman operators for dynamical systems,” preprint.
 - C., Ayton, Szőke, “Residual Dynamic Mode Decomposition,” **J. Fluid Mech.**, under minor rev.
 - Code: <https://github.com/MColbrook/Residual-Dynamic-Mode-Decomposition>

ResDMD: avoiding “too much”

$$\text{res}(\lambda, \mathbf{g})^2 = \frac{\mathbf{g}^* [K_2 - \lambda K_1^* - \bar{\lambda} K_1 + |\lambda|^2 G] \mathbf{g}}{\mathbf{g}^* G \mathbf{g}}$$

eigenvectors

eigenvalues

Algorithm 1:

1. Compute $G, K_1, K_2 \in \mathbb{C}^{N_K \times N_K}$ and eigendecomposition $K_1 V = G V \Lambda$.
2. For each eigenpair (λ, \mathbf{v}) , compute $\text{res}(\lambda, \mathbf{v})$.
3. **Output:** subset of e-vectors $V_{(\varepsilon)}$ & e-vals $\Lambda_{(\varepsilon)}$ with $\text{res}(\lambda, \mathbf{v}) \leq \varepsilon$ (ε = input tol).

Theorem (no spectral pollution): Suppose quad. rule converges. Then

$$\limsup_{M \rightarrow \infty} \max_{\lambda \in \Lambda^{(\varepsilon)}} \|(\mathcal{K} - \lambda)^{-1}\|^{-1} \leq \varepsilon$$

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$$\limsup_{M \rightarrow \infty} \max_{\lambda \in \Lambda^{(\varepsilon)}} \|(\mathcal{K} - \lambda)^{-1}\|^{-1} \leq \varepsilon$$

BUT: Typically, does not capture all of spectrum! (“too little”)

ResDMD: avoiding “too little”

$$\text{Spec}_\varepsilon(\mathcal{K}) = \bigcup_{\|\mathcal{B}\| \leq \varepsilon} \text{Spec}(\mathcal{K} + \mathcal{B}), \quad \lim_{\varepsilon \downarrow 0} \text{Spec}_\varepsilon(\mathcal{K}) = \text{Spec}(\mathcal{K})$$

Algorithm 2:

First convergent method for general \mathcal{K}

1. Compute $G, K_1, K_2 \in \mathbb{C}^{N_K \times N_K}$.
2. For z_k in comp. grid, compute $\tau_k = \min_{g = \sum_{j=1}^{N_K} \mathbf{g}_j \psi_j} \text{res}(z_k, g)$, corresponding g_k (gen. SVD).
3. **Output:** $\{z_k: \tau_k < \varepsilon\}$ (approx. of $\text{Spec}_\varepsilon(\mathcal{K})$), $\{g_k: \tau_k < \varepsilon\}$ (ε -pseudo-eigenfunctions).

Theorem (full convergence): Suppose the quadrature rule converges.

- **Error control:** $\{z_k: \tau_k < \varepsilon\} \subseteq \text{Spec}_\varepsilon(\mathcal{K})$ (as $M \rightarrow \infty$)
- **Convergence:** Converges locally uniformly to $\text{Spec}_\varepsilon(\mathcal{K})$ (as $N_K \rightarrow \infty$)

The Challenges

~~1) “Too much”: Approximate spurious modes $\lambda \notin \text{Spec}(\mathcal{K})$~~ ✓

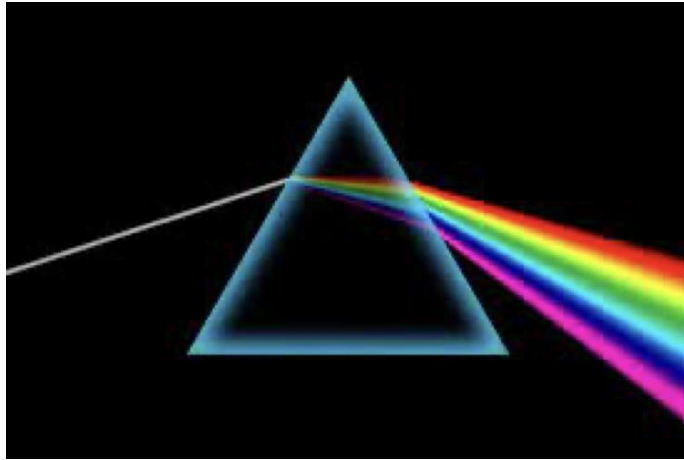
~~2) “Too little”: Miss parts of $\text{Spec}(\mathcal{K})$~~ ✓

3) Continuous spectra.

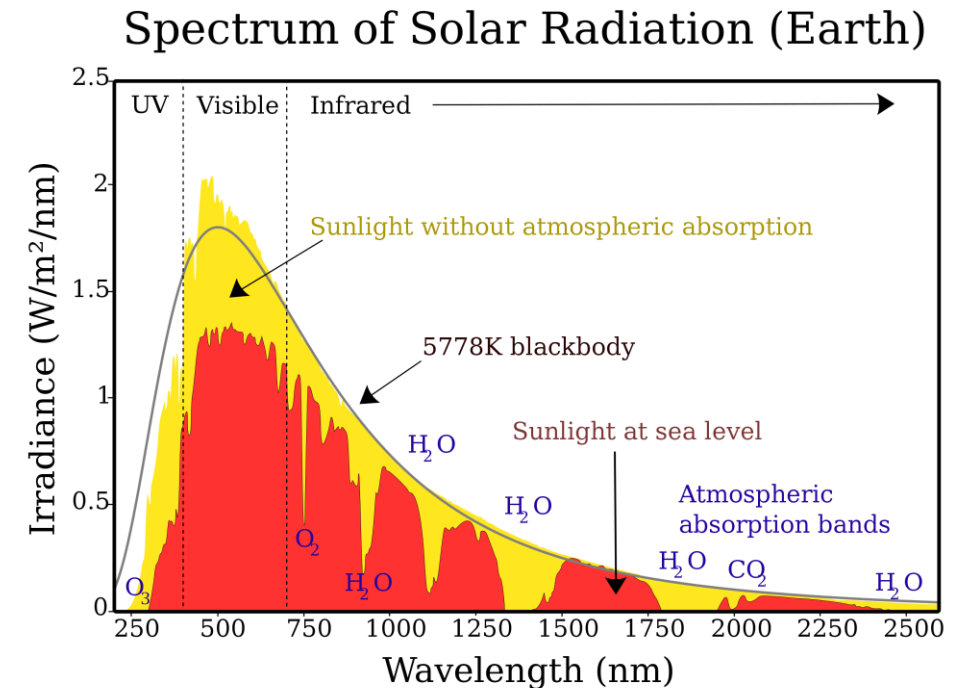
Verification: Is it right?

Continuous spectra

White light contains a continuous spectra



Often interesting to look at the intensity of each wavelength

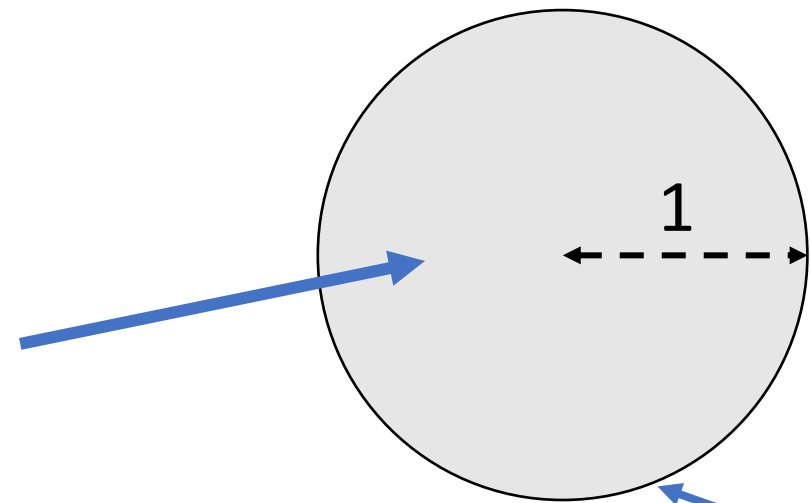


Setup for continuous spectra

Suppose system is measure preserving (e.g., Hamiltonian, ergodic, post-transient etc.)

$$\Leftrightarrow \mathcal{K}^* \mathcal{K} = I \text{ (isometry)}$$

$$\Rightarrow \text{Spec}(\mathcal{K}) \subseteq \{z: |z| \leq 1\}$$



(NB: we consider unitary extensions via Wold decomposition.)

spectral
measure
supp. on
boundary

Spectral decomposition of operators

$A \in \mathbb{C}^{n \times n}$ normal \Rightarrow O.N. basis of eigenvectors v_1, \dots, v_n :

$$v = \left(\sum_{k=1}^n \underset{\substack{\uparrow \\ \text{Projector onto Span}(v_k)}}}{v_k v_k^*} \right) v, \quad Av = \left(\sum_{k=1}^n \underset{\substack{\uparrow \\ \text{eigenvalues}}}{\lambda_k} v_k v_k^* \right) v, \quad v \in \mathbb{C}^n$$

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Energy of “v” in each eigenvector: $\mu_v(\lambda_j) = \langle v_j v_j^* v, v \rangle = |v_j^* v|^2$

This is called the spectral measure with respect to a vector v .

Spectral decomposition of operators

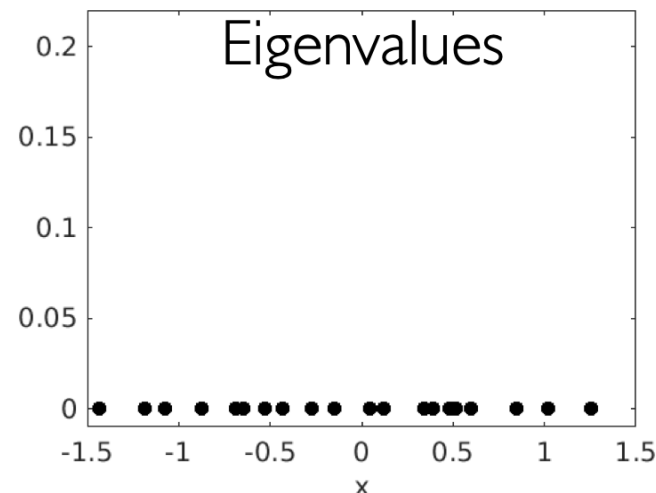
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↑
Projector onto $\text{Span}(v_k)$
↑
eigenvalues

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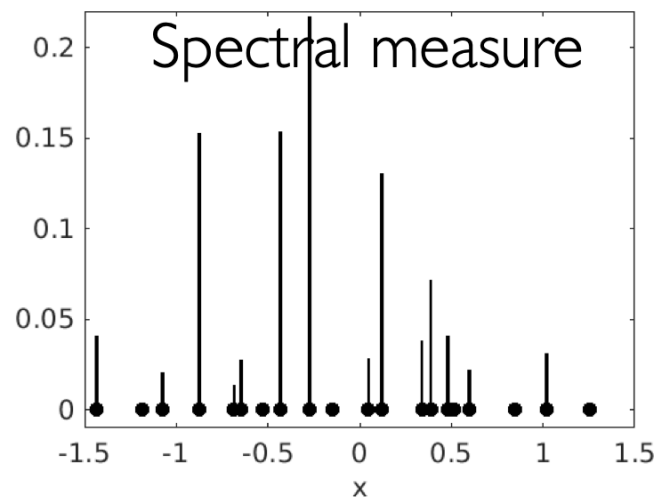
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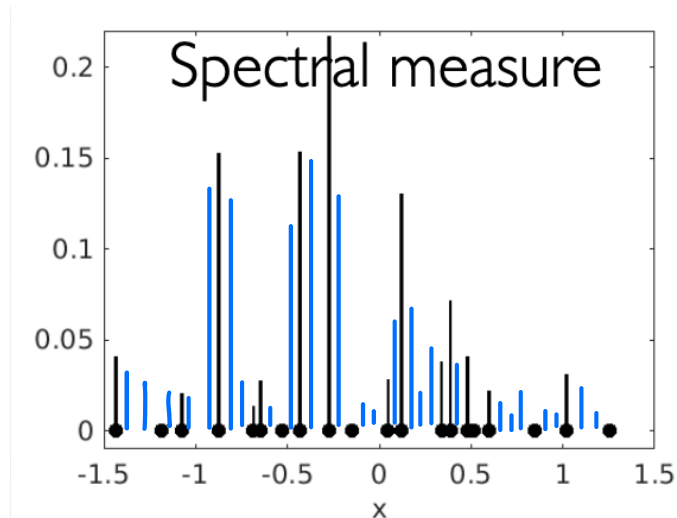
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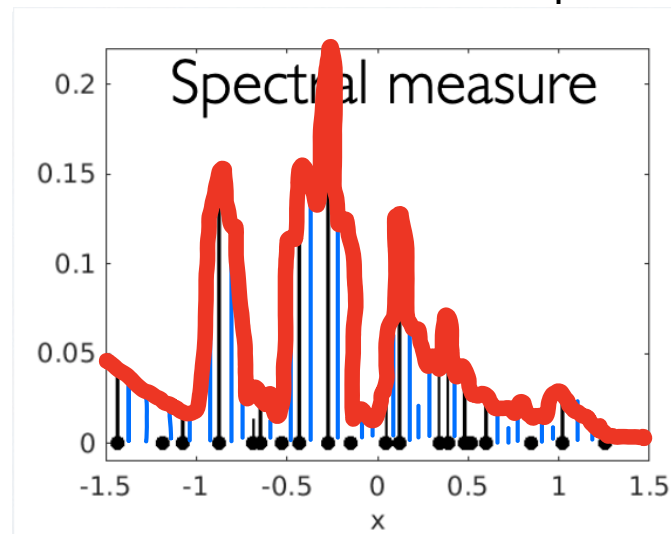
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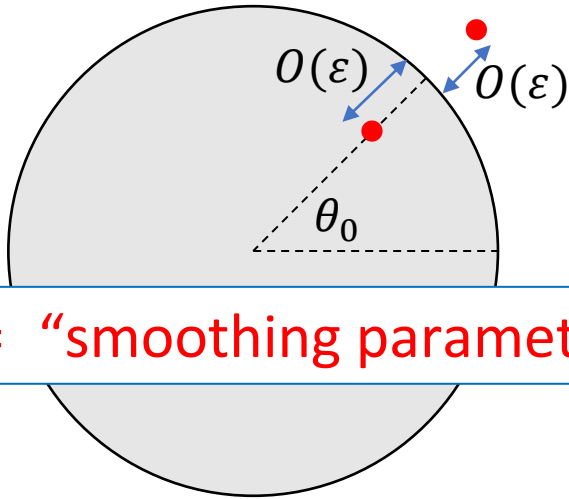
This is called the spectral measure with respect to a vector v .

\mathcal{K} is unitary \Rightarrow projection-valued measure ξ

$$g = \left(\int_{\mathbb{T}} d\xi(y) \right) g, \quad \mathcal{K}g = \left(\int_{\mathbb{T}} y d\xi(y) \right) g$$

Spectral measure $\nu_g(B) = \langle \xi(B)g, g \rangle$

Evaluating spectral measure



$\varepsilon =$ “smoothing parameter”

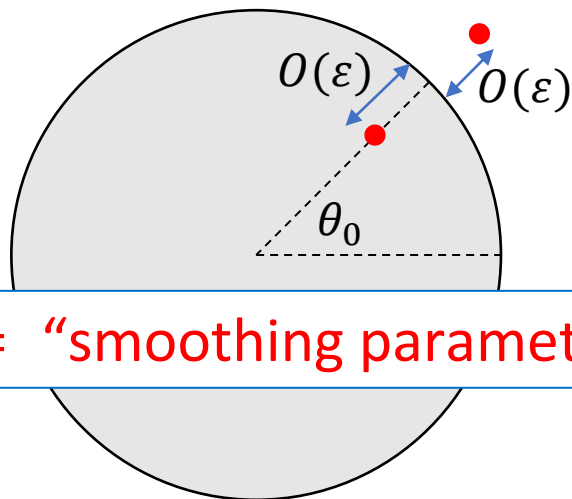
$$[P_\varepsilon * \nu_g](\theta_0) = \int_{[-\pi, \pi]_{\text{per}}} P_\varepsilon(\theta_0 - \theta) d\nu_g(\theta)$$

Smoothing convolution

Poisson kernel for
unit disk

$$P_\varepsilon(\theta_0) = \frac{1}{2\pi} \frac{(1 + \varepsilon)^2 - 1}{1 + (1 + \varepsilon)^2 - 2(1 + \varepsilon)\cos(\theta_0)}$$

Evaluating spectral measur



$\epsilon =$ “smoothing parameter”

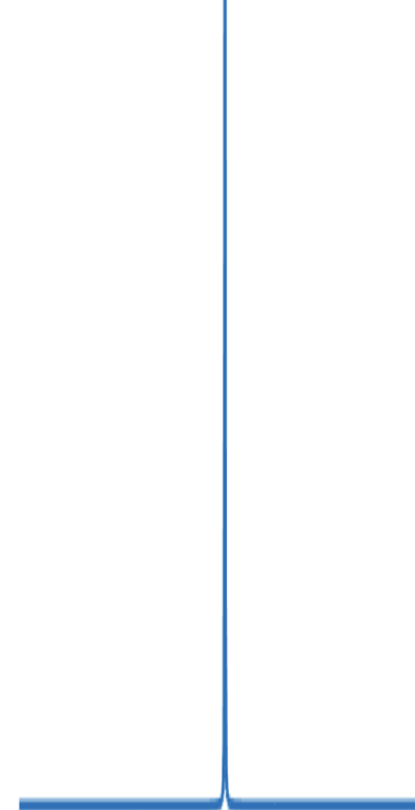
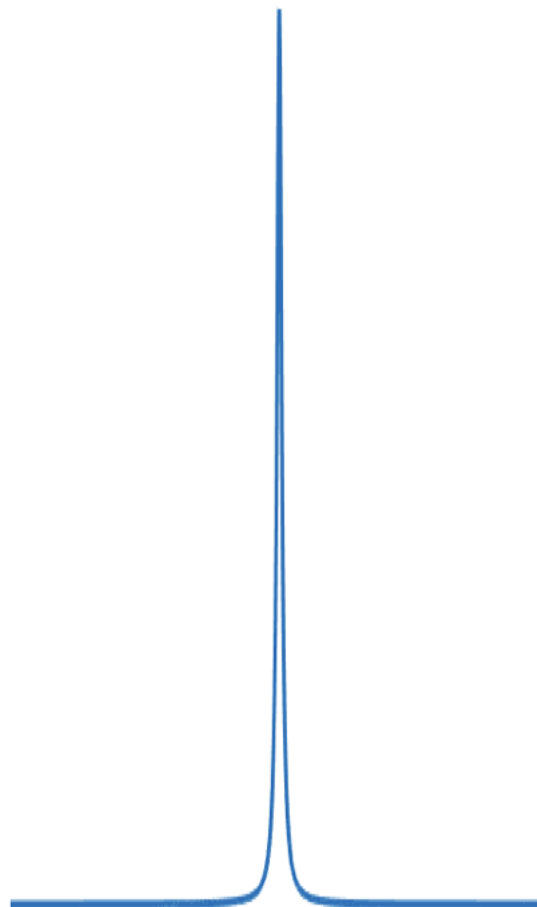
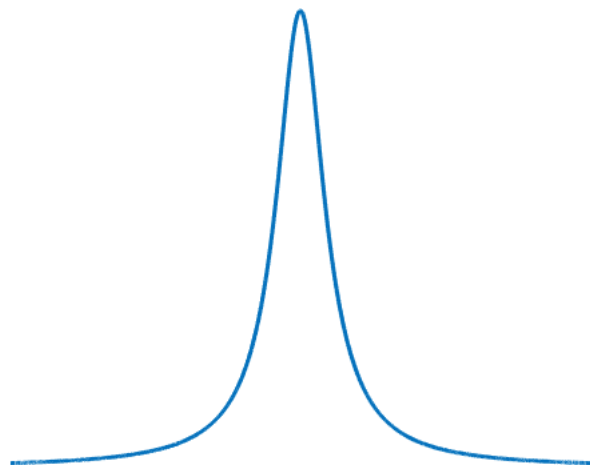
$$[P_\epsilon * \nu_g](\theta_0) = \int_{\mathbb{T}} P_\epsilon(\theta_0 - \theta) \, d\nu_g(\theta)$$

Smoothing co

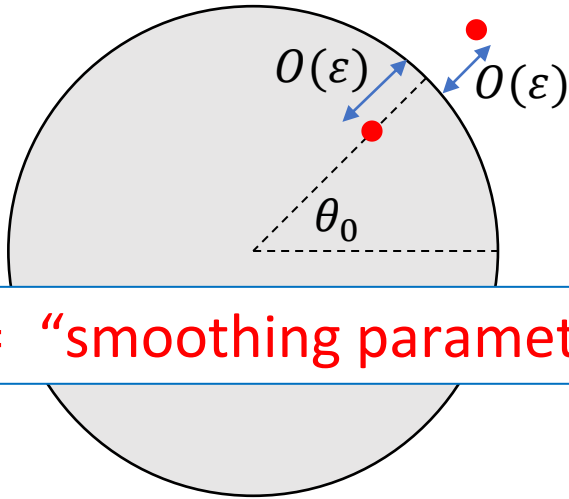
Poisso
unit di

$$\frac{1}{1 + \epsilon}$$

$\overline{0}$



Evaluating spectral measure



$\epsilon = \text{"smoothing parameter"}$

Smoothing convolution

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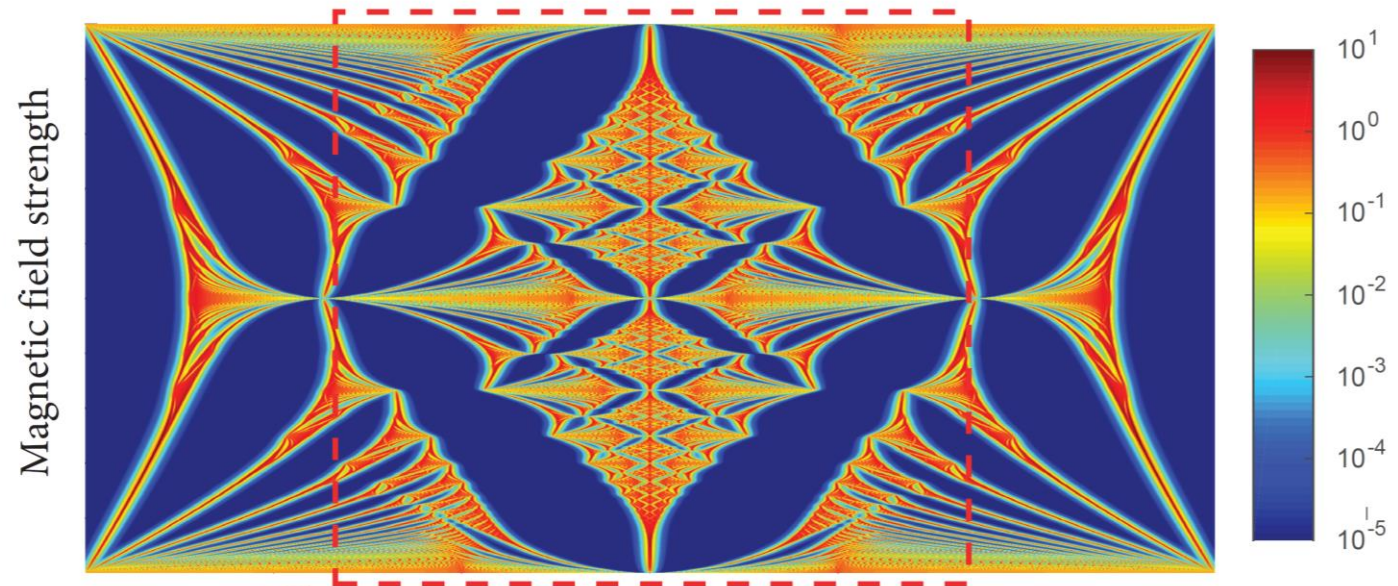
$$[P_\epsilon * \nu_g](\theta_0) = \mathcal{C}_g(e^{i\theta_0}(1 + \epsilon)^{-1}) - \mathcal{C}_g(e^{i\theta_0}(1 + \epsilon))$$

$$\mathcal{C}_g(z) = \int_{[-\pi, \pi]_{\text{per}}} \frac{e^{i\theta} d\nu_g(\theta)}{e^{i\theta} - z} = \begin{cases} \langle (\mathcal{K} - zI)^{-1}g, \mathcal{K}^*g \rangle, & \text{if } |z| > 1 \\ -z^{-1} \langle g, (\mathcal{K} - \bar{z}^{-1}I)^{-1}g \rangle, & \text{if } 0 < |z| < 1 \end{cases}$$

ResDMD computes
with error control

Analogous ideas are common in particle and condensed matter physics for computing spectral measures.

Spectral measures of self-adjoint operators



Horizontal slice = spectral measure at constant magnetic field strength.

Software package

SpecSolve available at <https://github.com/SpecSolve>
Capabilities: ODEs, PDEs, integral operators, discrete operators.

Example

$$\mathcal{K} = \begin{pmatrix} \overline{\alpha_0} & \overline{\alpha_1}\rho_0 & \rho_0\rho_1 & & & \\ \rho_0 & -\overline{\alpha_1}\alpha_0 & -\alpha_0\rho_1 & & & \\ & \overline{\alpha_2}\rho_1 & -\overline{\alpha_2}\alpha_1 & \overline{\alpha_3}\rho_2 & \rho_3\rho_2 & \\ & \rho_2\rho_1 & -\alpha_1\rho_2 & -\overline{\alpha_3}\alpha_2 & -\rho_3\alpha_2 & \ddots \\ & & & \overline{\alpha_4}\rho_3 & -\overline{\alpha_4}\alpha_3 & \ddots \\ & & & \ddots & \ddots & \ddots \end{pmatrix}$$

$$\alpha_j = (-1)^j 0.95^{(j+1)/2}, \quad \rho_j = \sqrt{1 - |\alpha_j|^2}$$

Generalised shift, typical building block of many dynamical systems.

Fix N_K , vary ε : unstable!

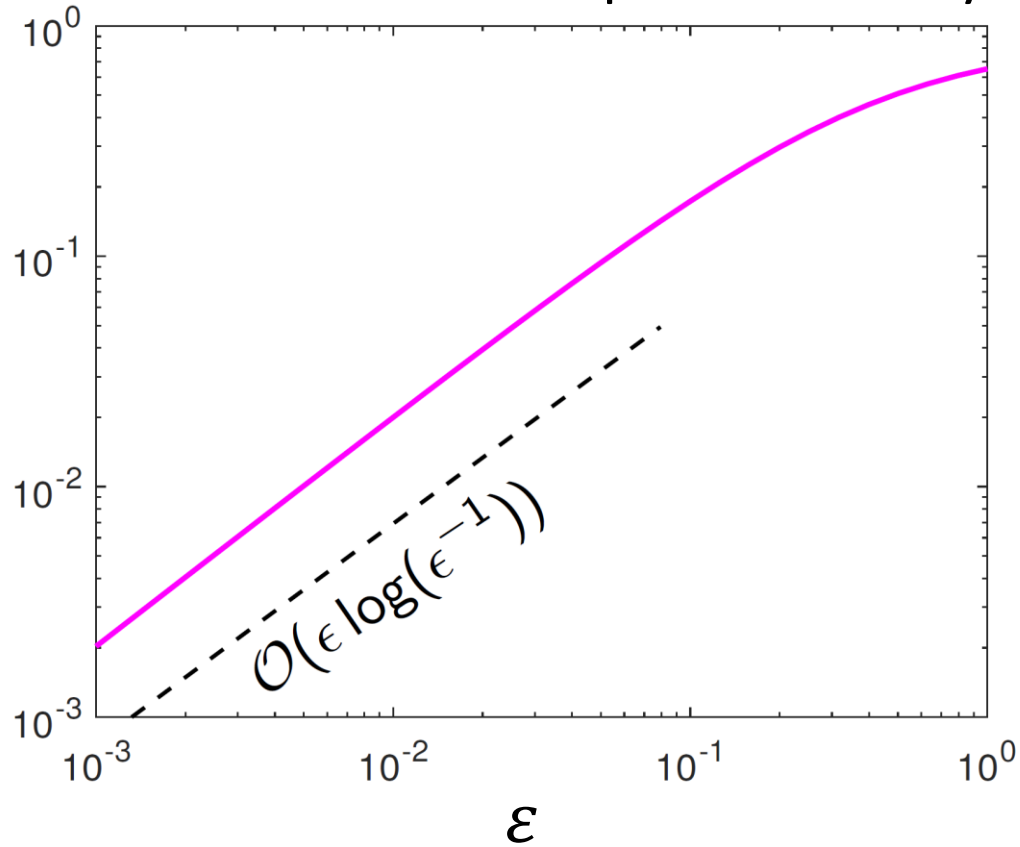
Fix ε , vary N_K : too smooth!

Adaptive: new matrix to compute residuals crucial

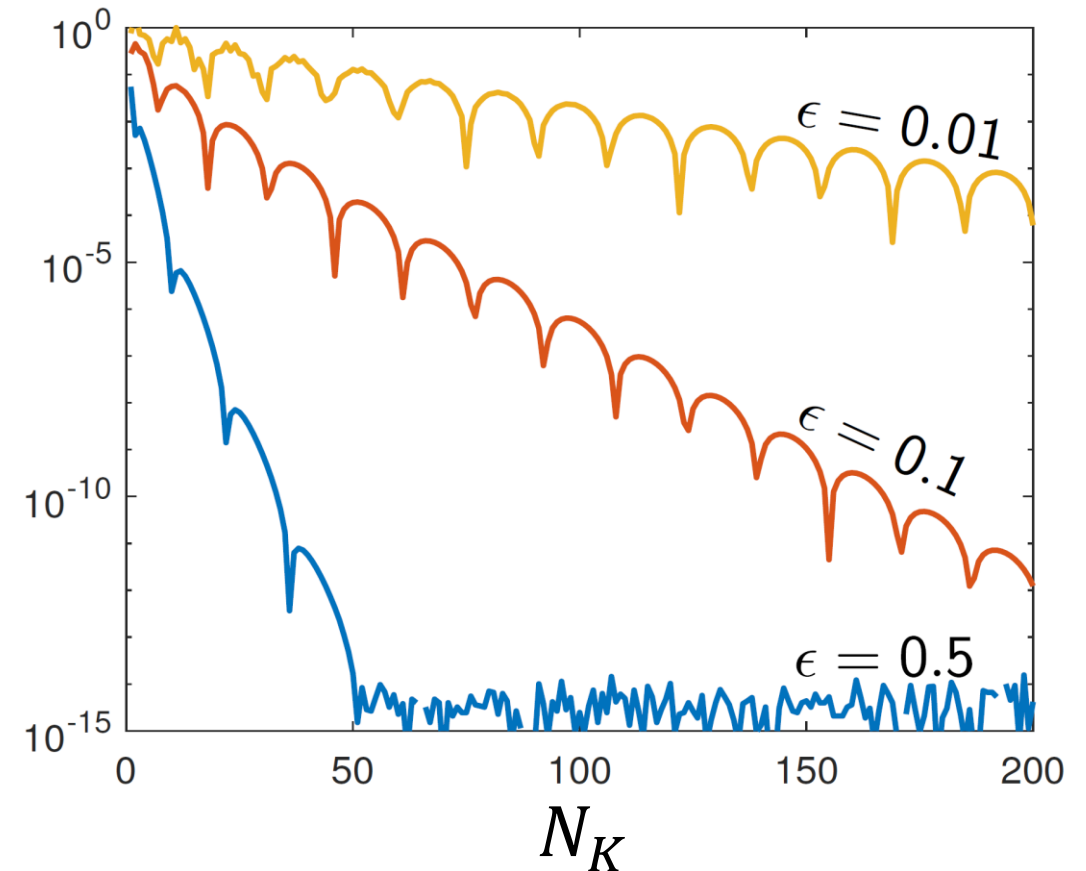
But ... slow convergence

Problem: As $\varepsilon \downarrow 0$, error is $O(\varepsilon \log(1/\varepsilon))$ and $N_K(\varepsilon) \rightarrow \infty$.

Pointwise error for spectral density



Error due to discretization

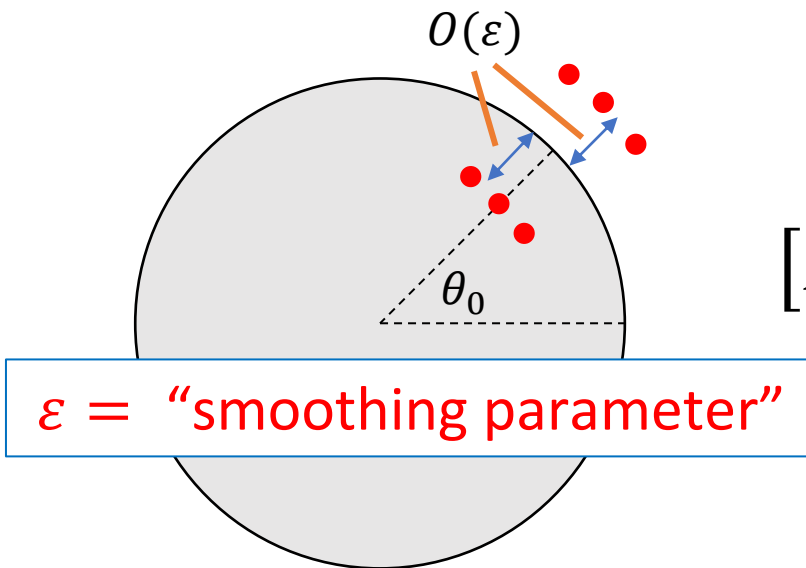


Small N_K critical in data-driven computations. Can we improve convergence rate?

High-order rational kernels

m th order rational kernels:

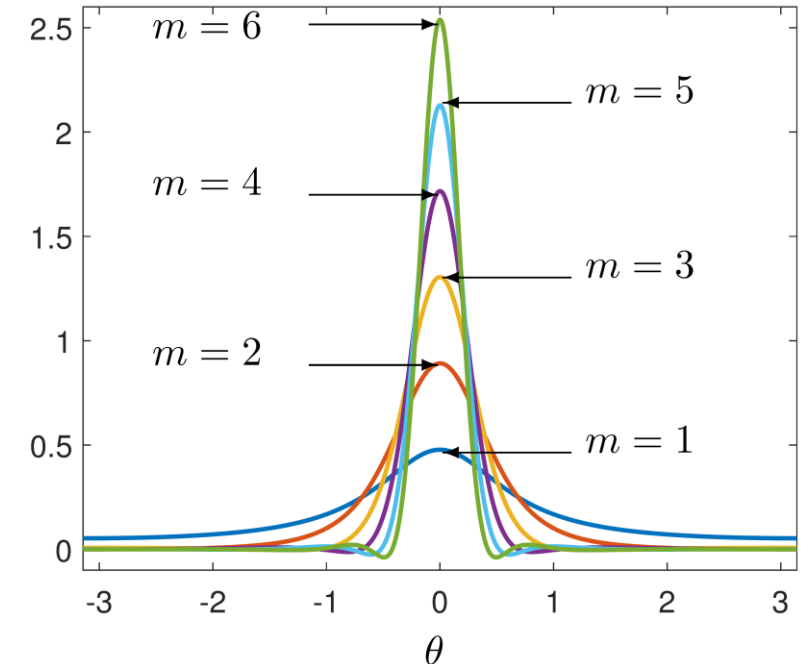
$$K_\varepsilon(\theta) = \frac{e^{-i\theta}}{2\pi} \sum_{j=1}^m \left[\frac{c_j}{e^{-i\theta} - (1 + \varepsilon \bar{z}_j)^{-1}} - \frac{d_j}{e^{-i\theta} - (1 + \varepsilon z_j)} \right]$$



ResDMD computes
with error control

$$[K_\varepsilon * v_g](\theta_0) = \sum_{j=1}^m \left[c_j \mathcal{C}_g(e^{i\theta_0}(1 + \varepsilon \bar{z}_j)^{-1}) - d_j \mathcal{C}_g(e^{i\theta_0}(1 + \varepsilon z_j)) \right]$$

Kernels



Smaller N_K (larger ε)

Convergence

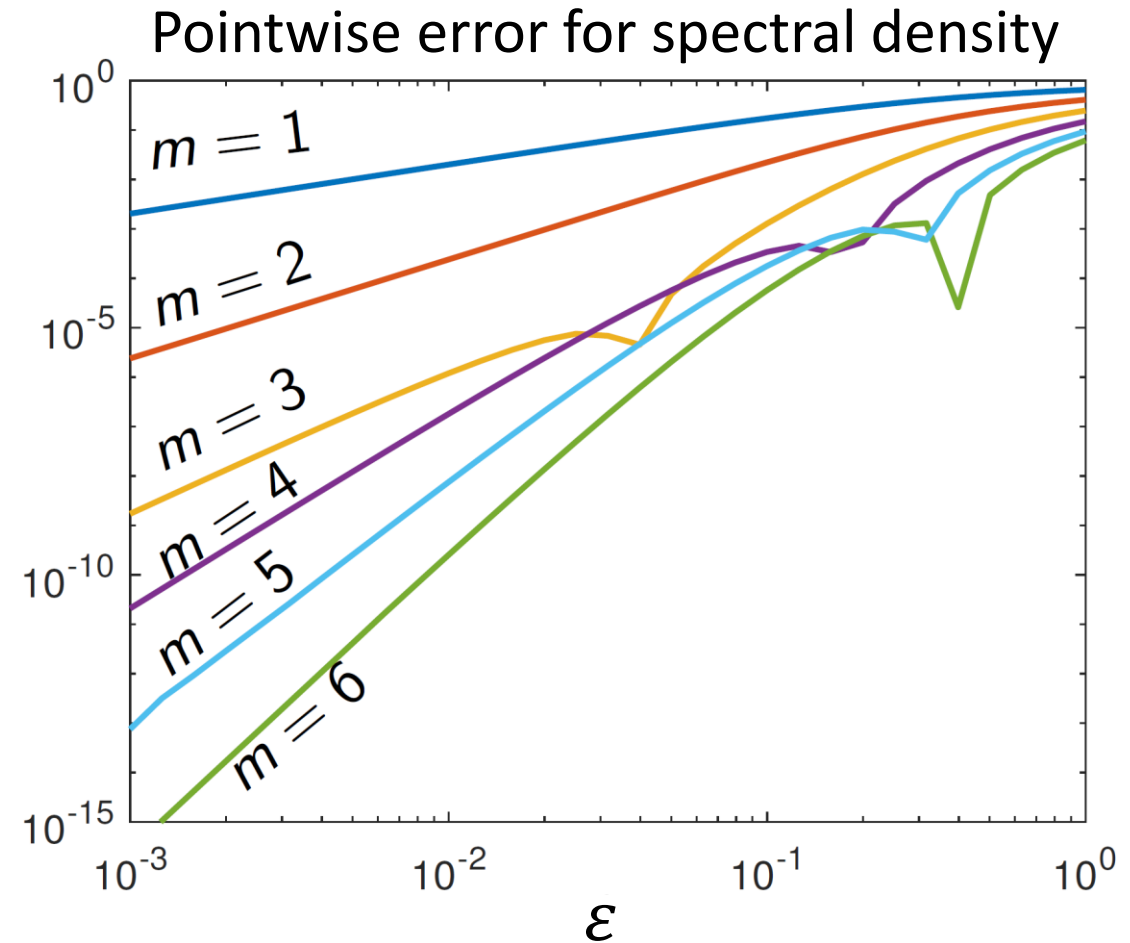
Theorem: Automatic selection of $N_K(\varepsilon)$ with $O(\varepsilon^m \log(1/\varepsilon))$ convergence:

- Density of continuous spectrum ρ_g .
(pointwise and L^p)
- Integration against test functions.
(weak convergence)

$$\int_{[-\pi, \pi]_{\text{per}}} h(\theta) [K_\varepsilon * \nu_g](\theta) \, d\theta$$

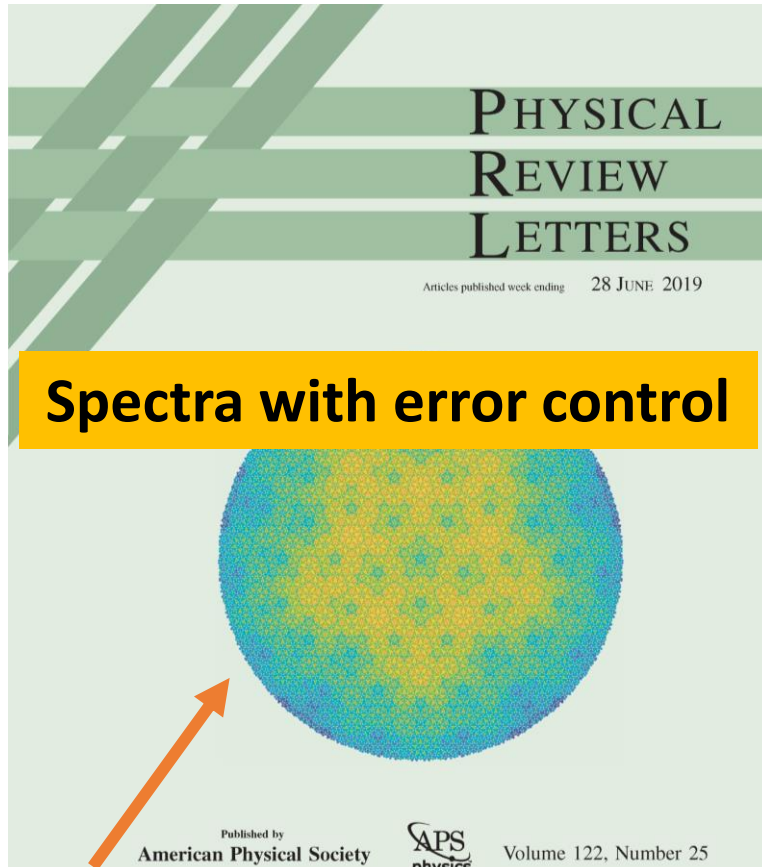
$$= \int_{[-\pi, \pi]_{\text{per}}} h(\theta) \, d\nu_g(\theta) + O(\varepsilon^m \log(1/\varepsilon))$$

Also recover discrete spectrum.



Is it right?

The importance of verification

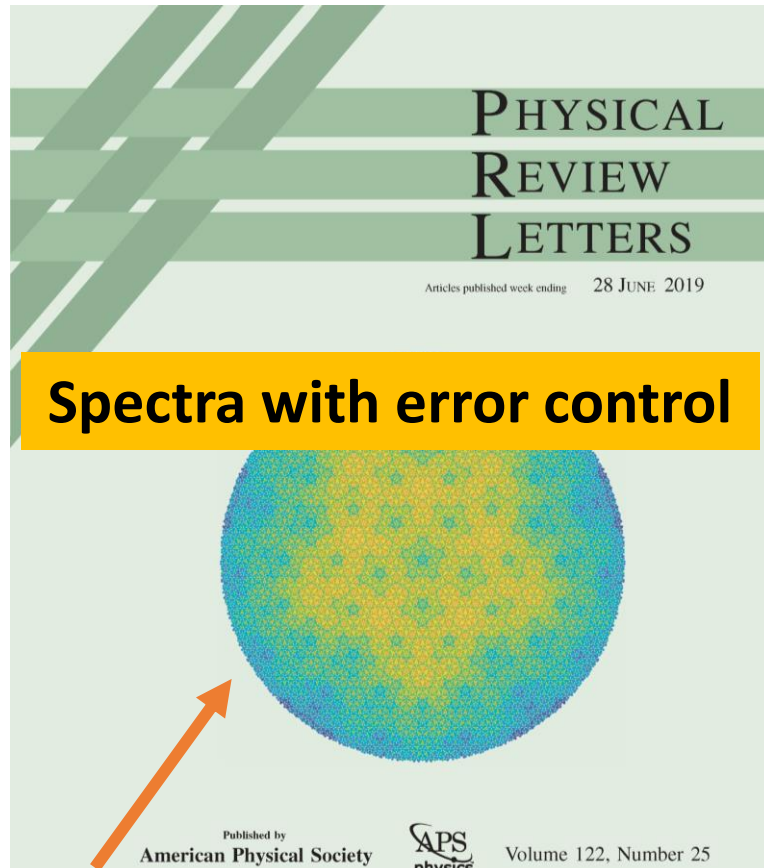


E.g., ground state of quasicrystal

- C., Roman, Hansen, “How to compute spectra with error control,” *Phys. Rev. Lett.*, 2019.

Is it right?

The importance of verification



E.g., ground state of quasicrystal



**Certainty in computed
spectral properties**

▼ PHYSICAL REVIEW B
covering condensed matter and materials physics

Highlights

Editors' Suggestion

Bulk localized transport states in infinite and finite quasicrystals via magnetic aperiodicity
Phys. Rev. B

B **BLT**

E.g., new physical phenomena:
bulk localised transport states

- C., Roman, Hansen, "How to compute spectra with error control," **Phys. Rev. Lett.**, 2019.
- Johnstone, C., Nielsen, Öhberg, Duncan, "Bulk Localised Transport States in Infinite and Finite Quasicrystals via Magnetic Aperiodicity," **Phys. Rev. B**, 2022.

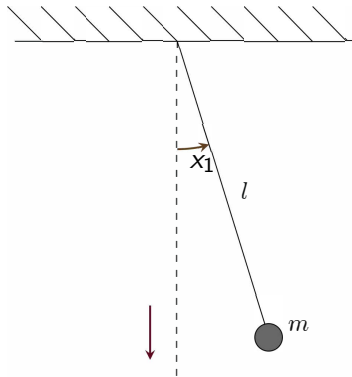
The Challenges

- ~~1) “Too much”: Approximate spurious modes $\lambda \notin \text{Spec}(\mathcal{K})$~~ ✓
- ~~2) “Too little”: Miss parts of $\text{Spec}(\mathcal{K})$~~ ✓
- ~~3) Continuous spectra.~~ ✓

Verification: Is it right?

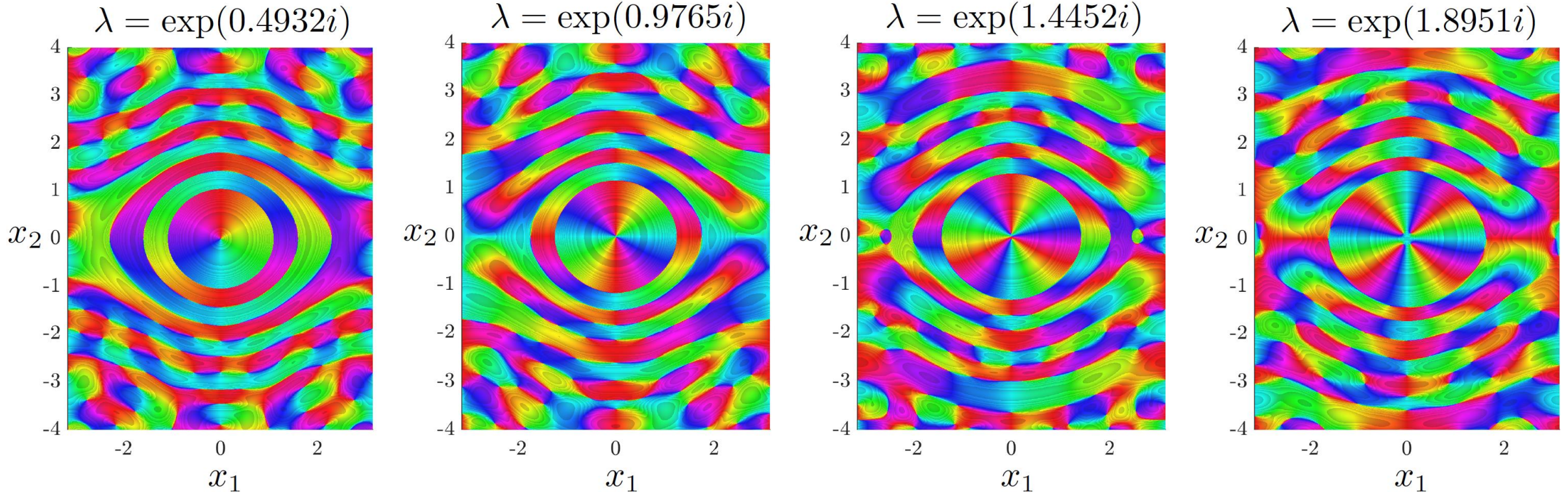
Example: non-linear pendulum

$$\dot{x}_1 = x_2, \quad \dot{x}_2 = -\sin(x_1), \quad \Omega = [-\pi, \pi]_{\text{per}} \times \mathbb{R}$$



Computed pseudospectra ($\varepsilon = 0.25$). Eigenvalues of \mathbb{K} shown as dots (spectral pollution).

Approximate eigenfunctions



Colour represents complex argument, constant modulus shown as shadowed steps.
All residuals smaller than $\varepsilon = 0.05$ (made smaller by increasing N_K).

Quadrature with trajectory data

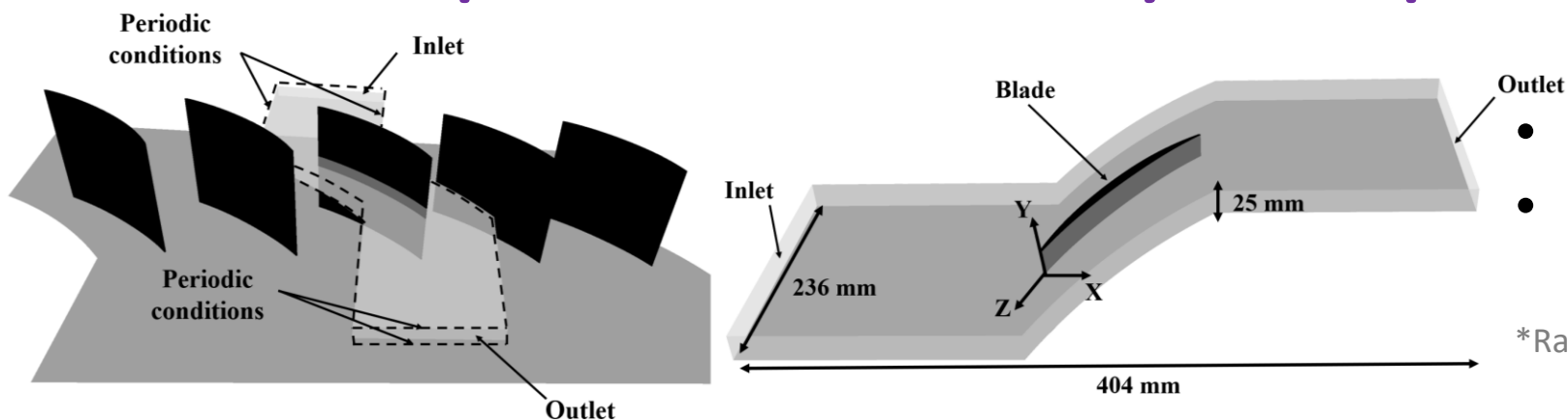
$$\text{E.g., } \langle \mathcal{K}\psi_k, \psi_j \rangle = \lim_{M \rightarrow \infty} \sum_{m=1}^M w_m \overline{\psi_j(x^{(m)})} \underbrace{\psi_k(y^{(m)})}_{[\mathcal{K}\psi_k](x^{(m)})}$$

Three examples:

- **High-order quadrature:** $\{x^{(m)}, w_m\}_{m=1}^M$ M -point quadrature rule.
 Rapid convergence. Requires free choice of $\{x^{(m)}\}_{m=1}^M$ and small d .
- **Random sampling:** $\{x^{(m)}\}_{m=1}^M$ selected at random.
 Large d . Slow Monte Carlo $O(M^{-1/2})$ rate of convergence.
- **Ergodic sampling:** $x^{(m+1)} = F(x^{(m)})$.
 Single trajectory, large d . Requires ergodicity, convergence can be slow.

Most common

Example: Trustworthy computation for large d

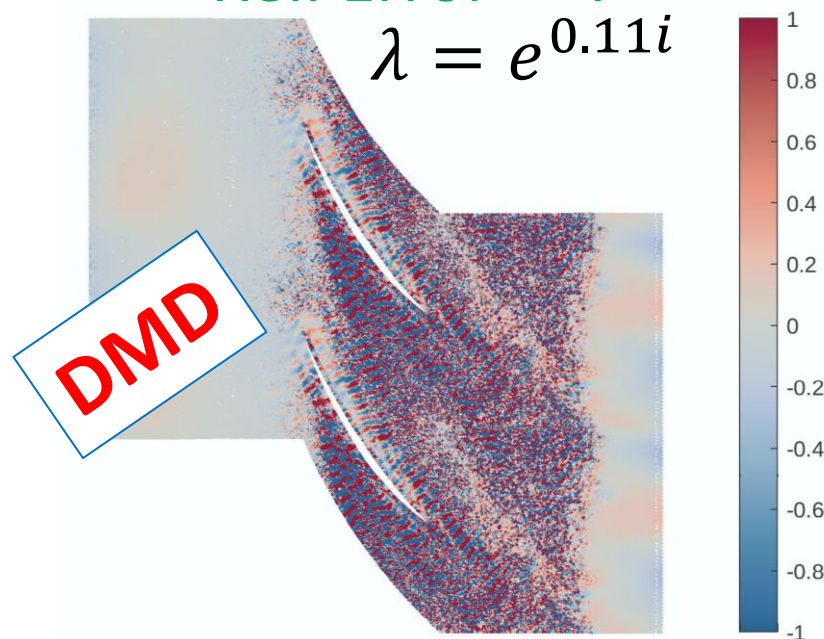


- Reynolds number $\approx 3.9 \times 10^5$
- Ambient dimension (d) $\approx 300,000$ (number of measurement points)

*Raw measurements provided by Stephane Moreau (Sherbrooke)

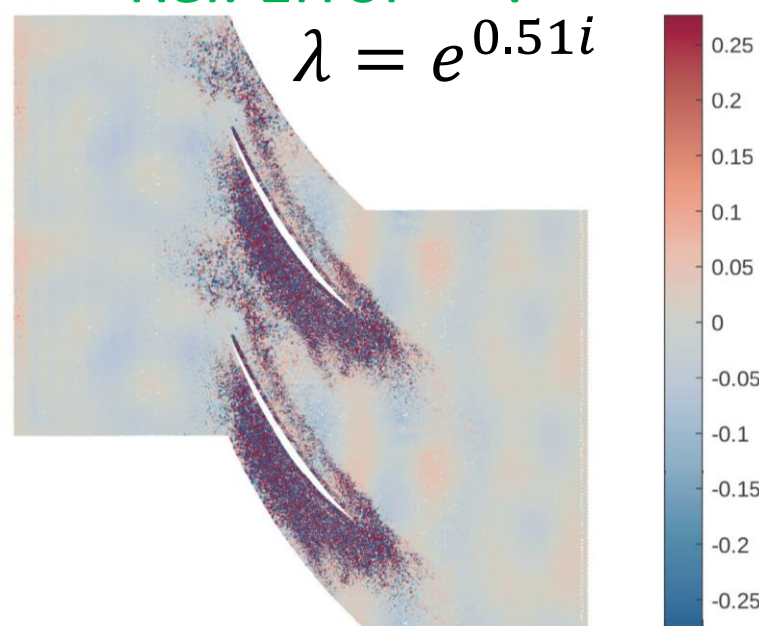
Rel. Error = ?

$$\lambda = e^{0.11i}$$



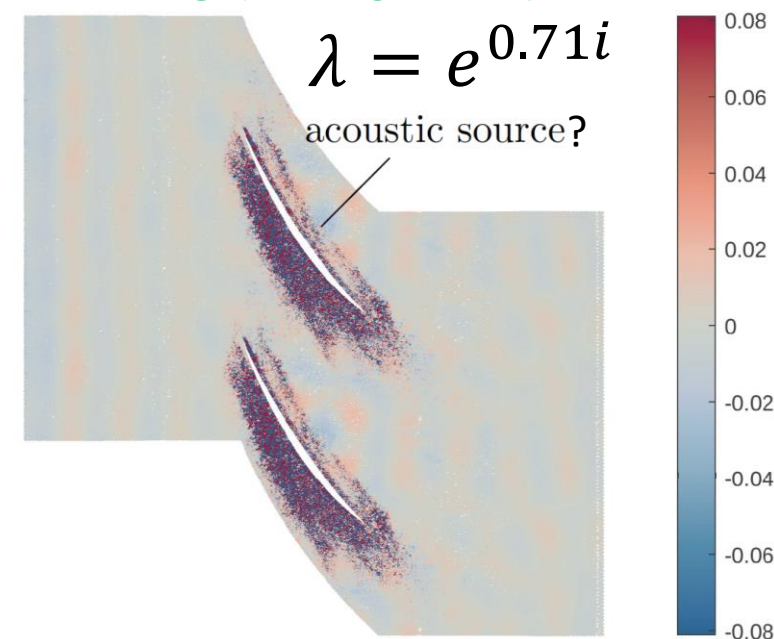
Rel. Error = ?

$$\lambda = e^{0.51i}$$

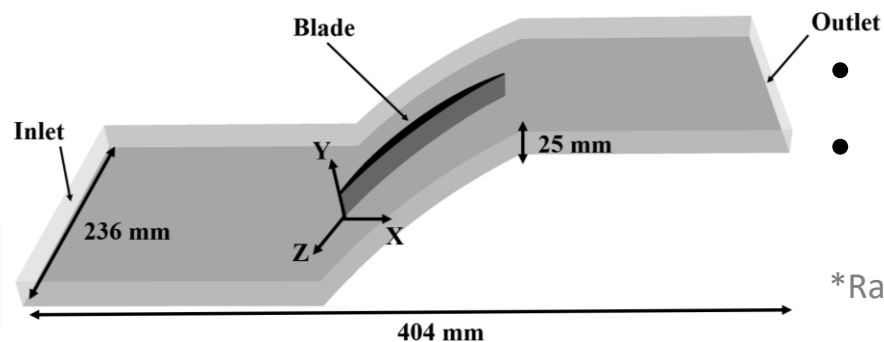
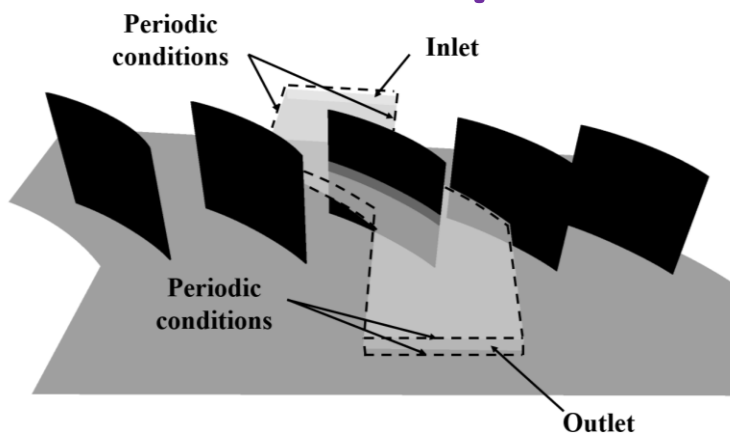


Rel. Error = ?

$$\lambda = e^{0.71i}$$



Example: Trustworthy computation for large d



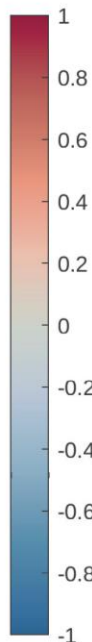
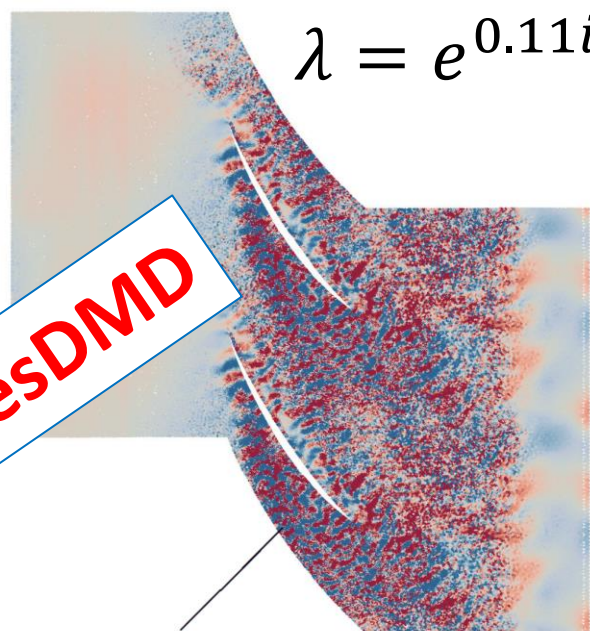
- Reynolds number $\approx 3.9 \times 10^5$
- Ambient dimension (d) $\approx 300,000$ (number of measurement points)

*Raw measurements provided by Stephane Moreau (Sherbrooke)

Rel. Error ≤ 0.0054

$$\lambda = e^{0.11i}$$

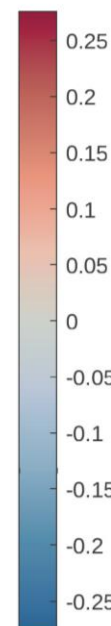
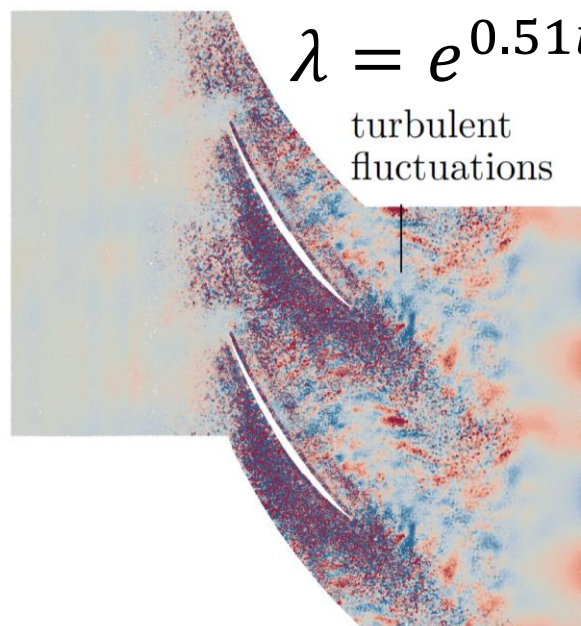
ResDMD



Rel. Error ≤ 0.0128

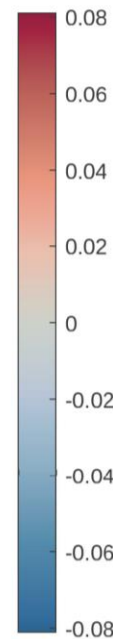
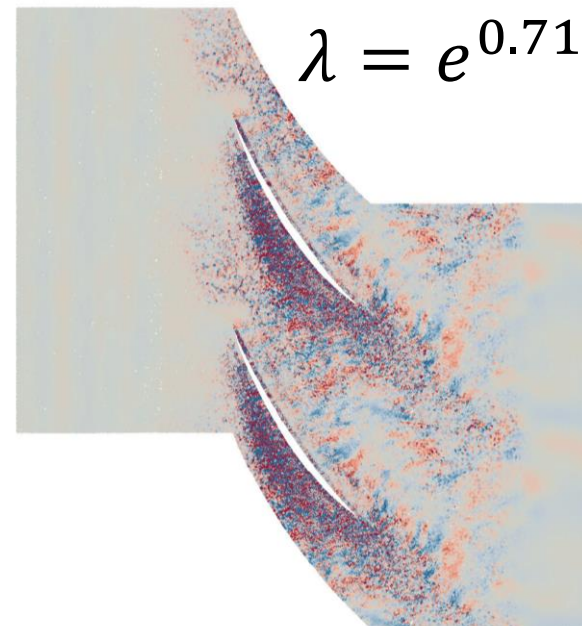
$$\lambda = e^{0.51i}$$

turbulent
fluctuations



Rel. Error ≤ 0.0196

$$\lambda = e^{0.71i}$$



acoustic vibrations

- C., T., "Rigorous data-driven computation of spectral properties of Koopman operators for dynamical systems," preprint.

Large d ($\Omega \subseteq \mathbb{R}^d$): robust and scalable

Popular to learn dictionary $\{\psi_1, \dots, \psi_{N_K}\}$

E.g., DMD with truncated SVD (linear dictionary, most popular), kernel methods (this talk), neural networks, etc.

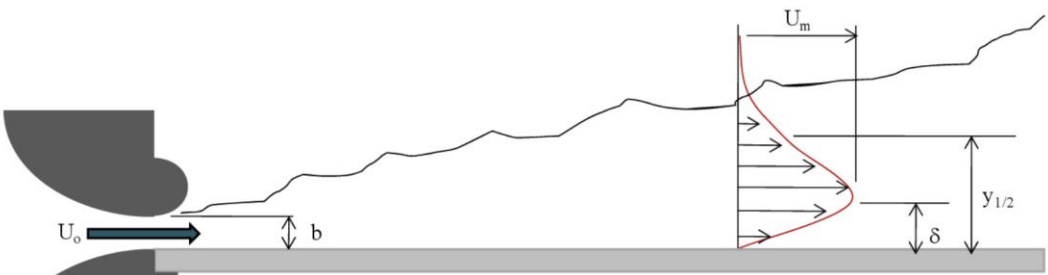
Q: Is discretisation $\text{span}\{\psi_1, \dots, \psi_{N_K}\}$ large/rich enough?

Above algorithms:

- Pseudospectra: $\{z_k: \tau_k < \varepsilon\} \subseteq \text{Spec}_\varepsilon(\mathcal{K})$ **error control**
- Spectral measures: $\mathcal{C}_g(z)$ and smoothed measures **adaptive check**

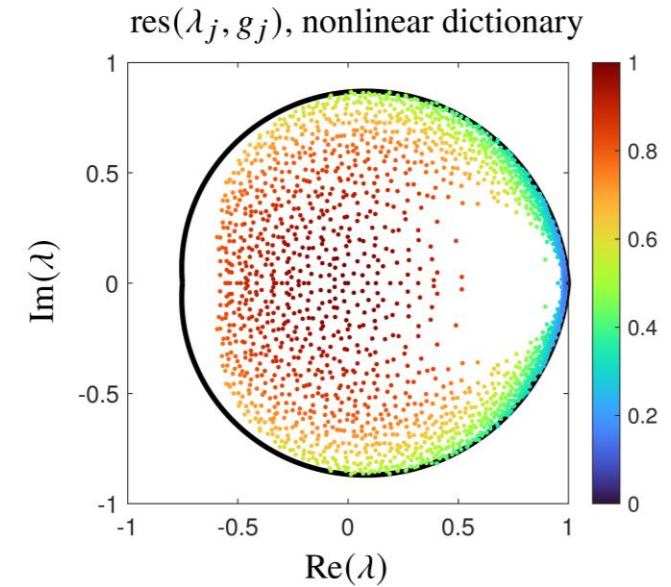
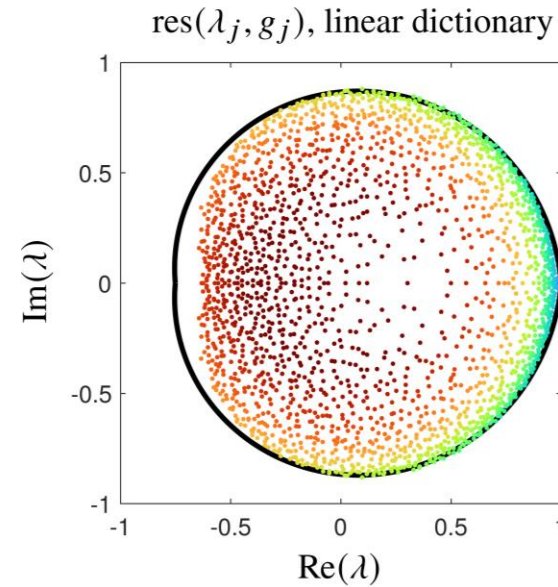
\Rightarrow Rigorously **verify** learnt dictionary $\{\psi_1, \dots, \psi_{N_K}\}$

Example: Verify the dictionary

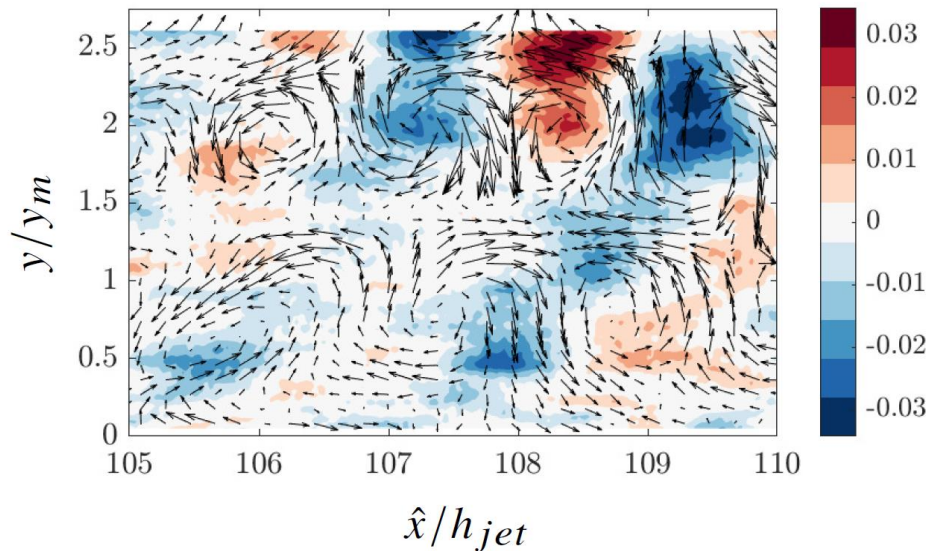


- Reynolds number $\approx 6.4 \times 10^4$
- Ambient dimension (d) $\approx 100,000$ (velocity at measurement points)

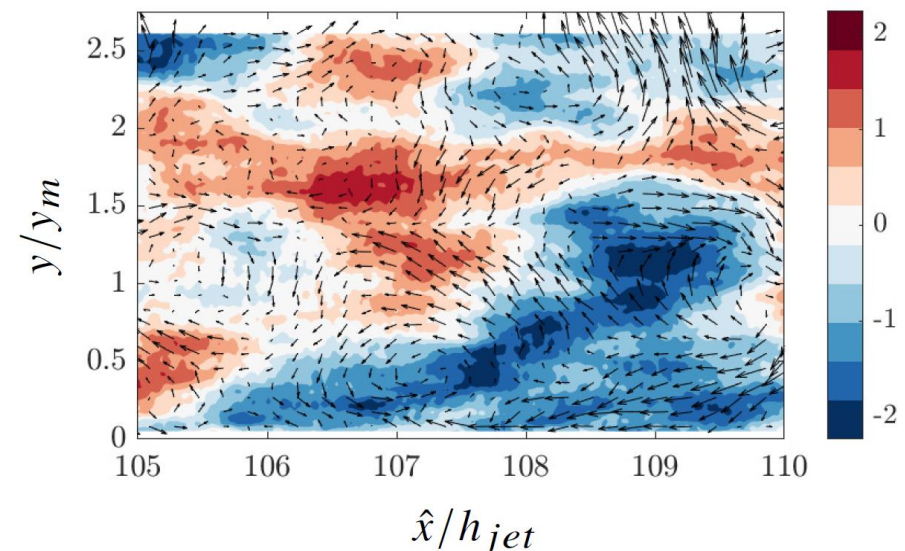
*Raw measurements provided by Máté Szőke (Virginia Tech)



$$\lambda = 0.9439 + 0.2458i, \text{ error} \leq 0.0765$$

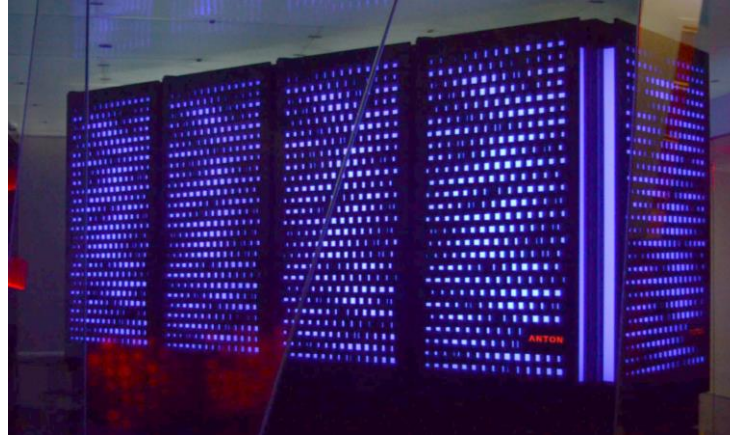
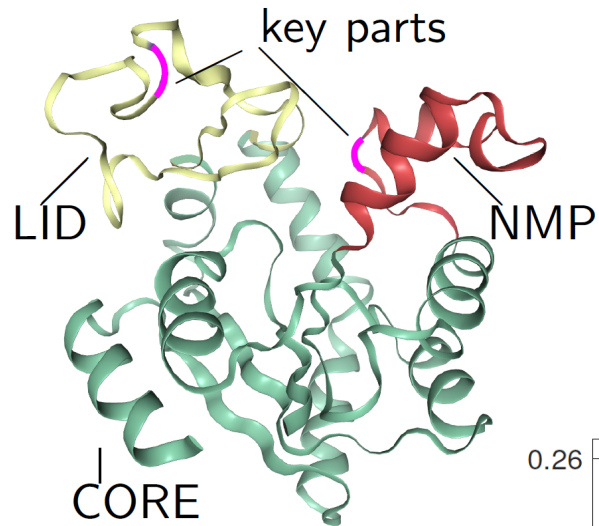


$$\lambda = 0.8948 + 0.1065i, \text{ error} \leq 0.1105$$



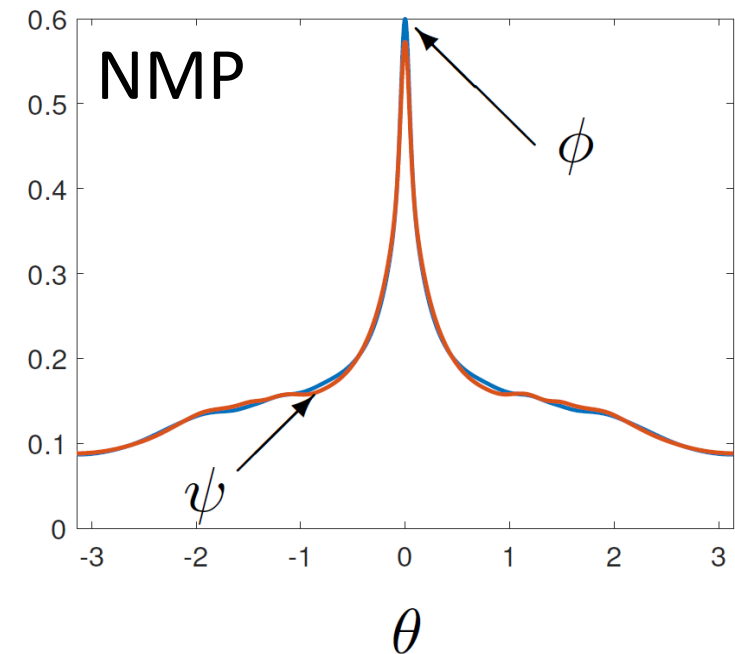
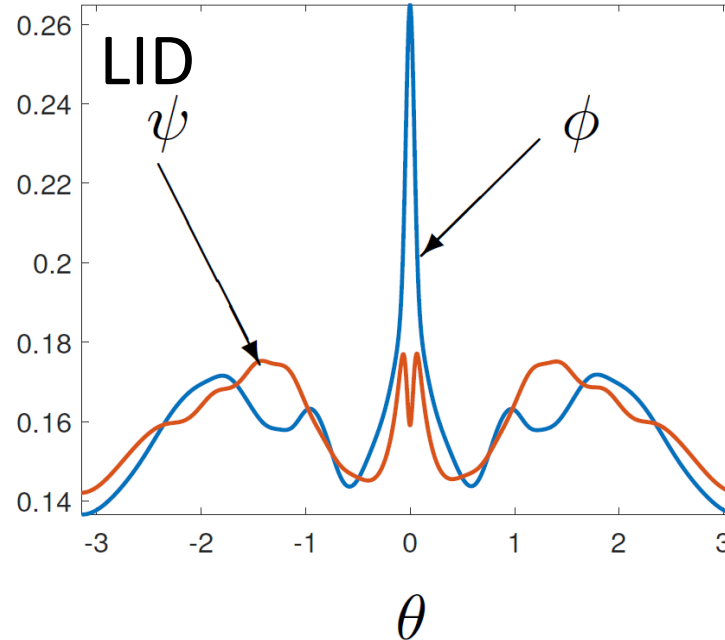
Example: molecular dynamics (Adenylate Kinase)

Adenylate Kinase

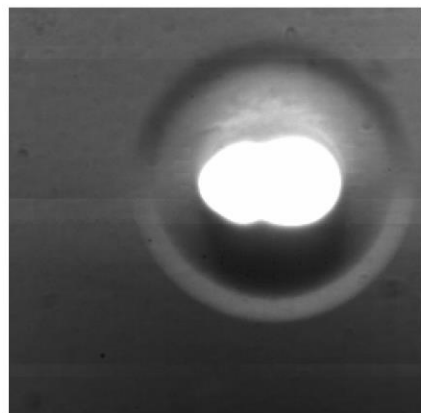
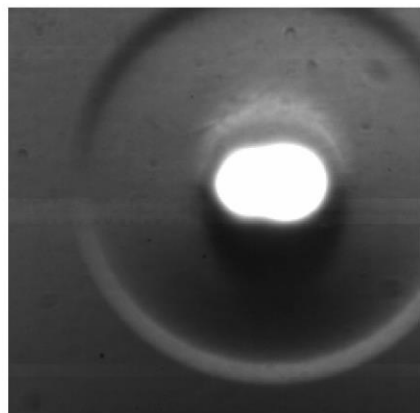
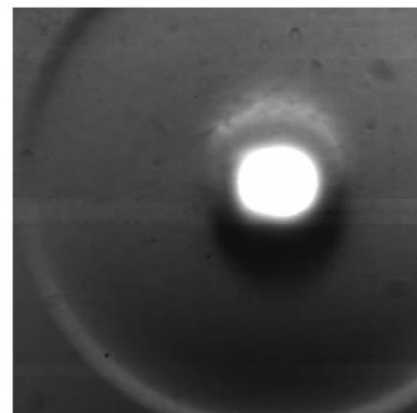
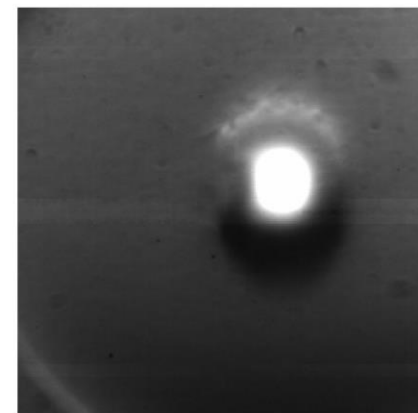
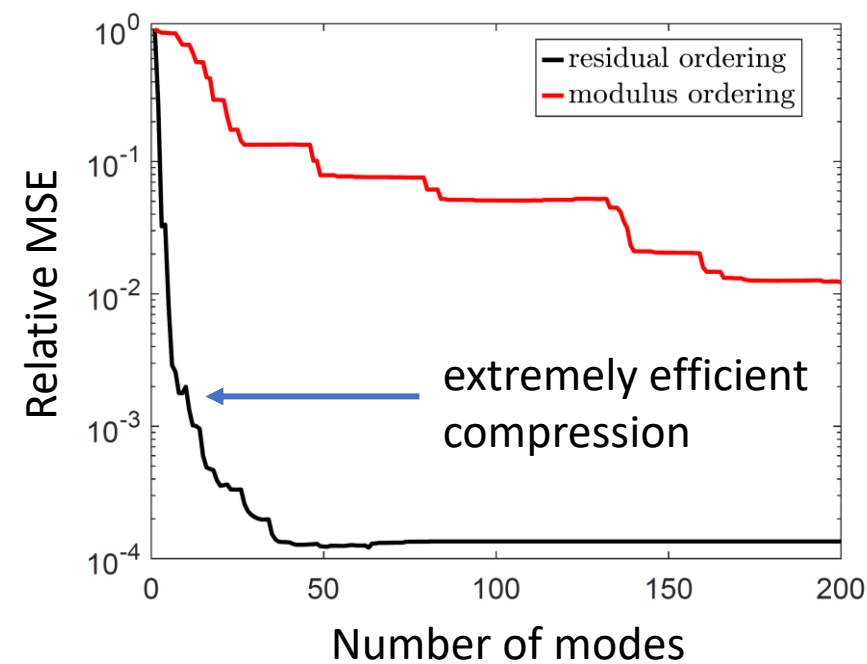
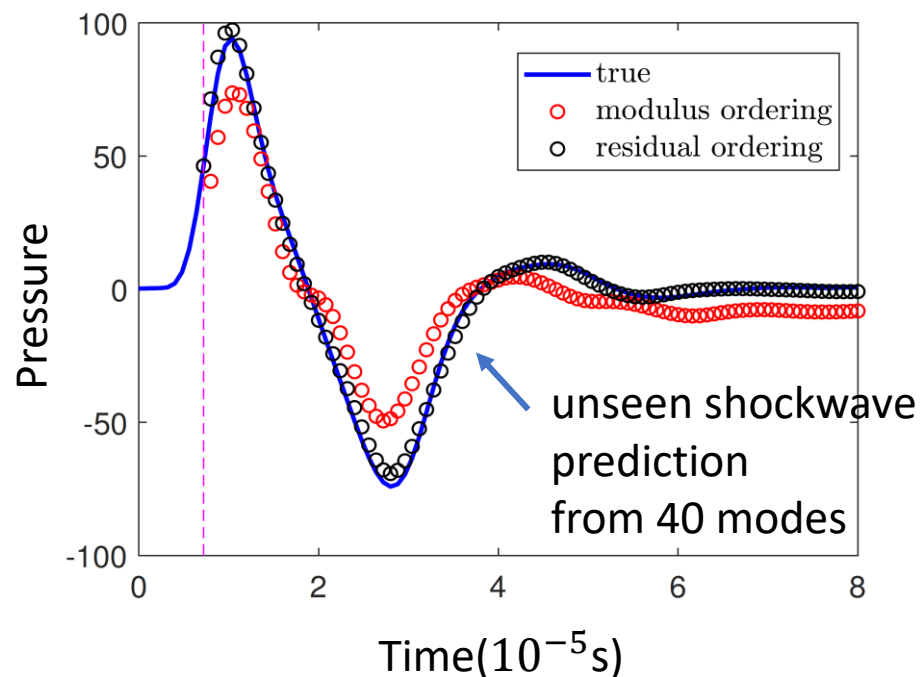


- Ambient dimension (d) $\approx 20,000$ (positions and momenta of atoms)
- 6th order kernel (spec res 10^{-6})

*Dataset: www.mdanalysis.org/MDAnalysisData/adk_equilibrium.html



Example: Trustworthy Koopman mode decomposition

a) $t = 5 \mu\text{s}$ b) $t = 10 \mu\text{s}$ c) $t = 15 \mu\text{s}$ d) $t = 20 \mu\text{s}$ 

- C., Ayton, Szőke, “Residual Dynamic Mode Decomposition,” **J. Fluid Mech.**, under minor rev.

Wider programme

- Inf.-dim. computational analysis \Rightarrow **Compute spectral properties rigorously.**
- Continuous linear algebra \Rightarrow **Avoid the woes of discretisation**
- Solvability Complexity Index hierarchy \Rightarrow **Classify diff. of comp. problems, prove algs are optimal.**
- **Extends to:** Foundations of AI, optimization, computer-assisted proofs, and PDE learning.

-
- C., “On the computation of geometric features of spectra of linear operators on Hilbert spaces,” **Found. Comput. Math.**, to appear.
 - Boullé, T., “Learning elliptic partial differential equations with randomized linear algebra”, **Found. Comput. Math.**, 2022.
 - Boullé, Kim, Shi, T., “Learning Green's functions associated with parabolic partial differential equations”, **JMLR**, to appear.
 - C., Horning, T. “Computing spectral measures of self-adjoint operators,” **SIAM Rev.**, 2021.
 - C., Hansen, “The foundations of spectral computations via the solvability complexity index hierarchy,” **J. Eur. Math. Soc.**, 2022.
 - C., Antun, Hansen, “The difficulty of computing stable and accurate neural networks: On the barriers of deep learning and Smale’s 18th problem,” **Proc. Natl. Acad. Sci. USA**, 2022.
 - C., “Computing spectral measures and spectral types,” **Comm. Math. Phys.**, 2021.
 - C., Roman, Hansen, “How to compute spectra with error control,” **Phys. Rev. Lett.**, 2019.
 - C., “Computing semigroups with error control,” **SIAM J. Numer. Anal.**, 2022.
 - Gilles, T., “Continuous analogues of Krylov methods for differential operators,” **SIAM J. Numer. Anal.**, 2019.
 - Horning, T., “FEAST for Differential Eigenvalue Problems,” **SIAM J. Numer. Anal.**, 2020.

Summary: rigorous data-driven Koopmanism!

- “Too much” or “Too little”

Idea: New matrix for residual \Rightarrow **ResDMD** for computing spectra.

- Continuous spectra and spectral measures:

Idea: Convolution with rational kernels via resolvent and **ResDMD**.

- Is it right?

Idea: Use **ResDMD** to verify computations. E.g., learned dictionaries.

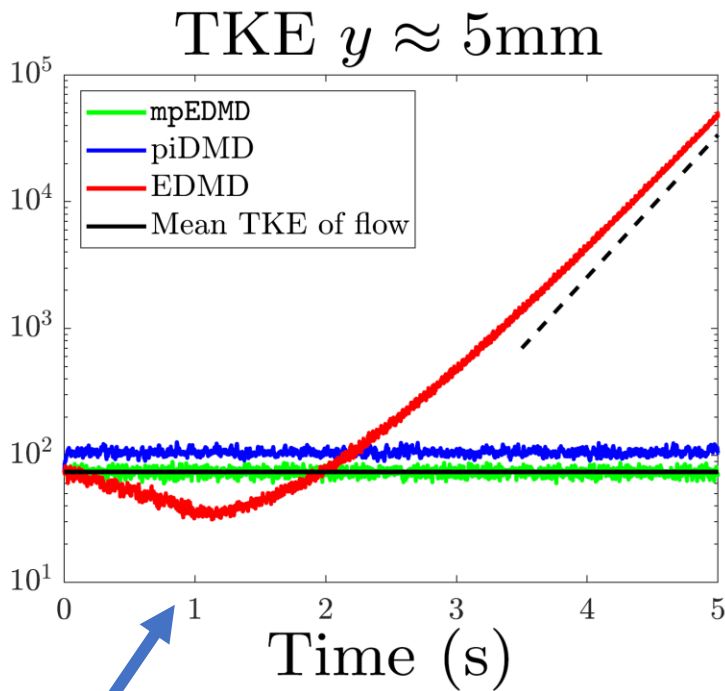
Code:

<https://github.com/MColbrook/Residual-Dynamic-Mode-Decomposition>

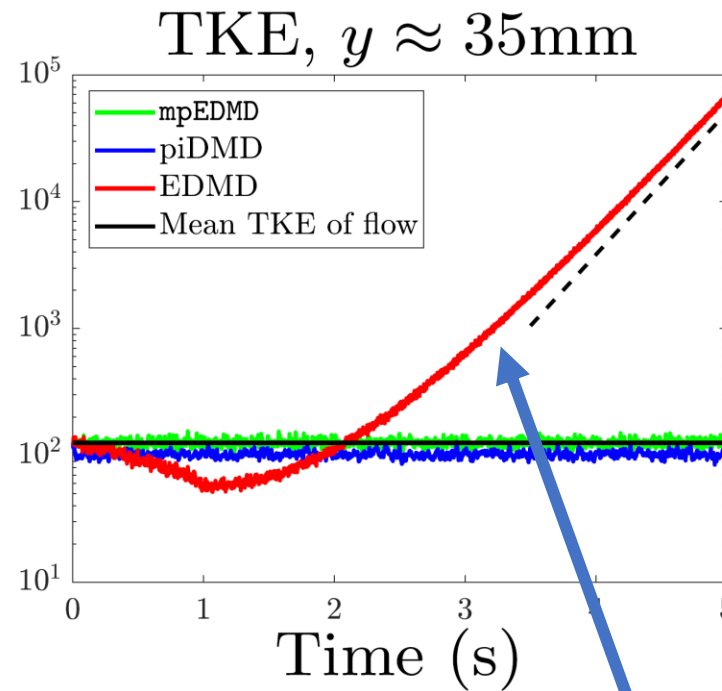
Additional slides...

measure-preserving EDMD...

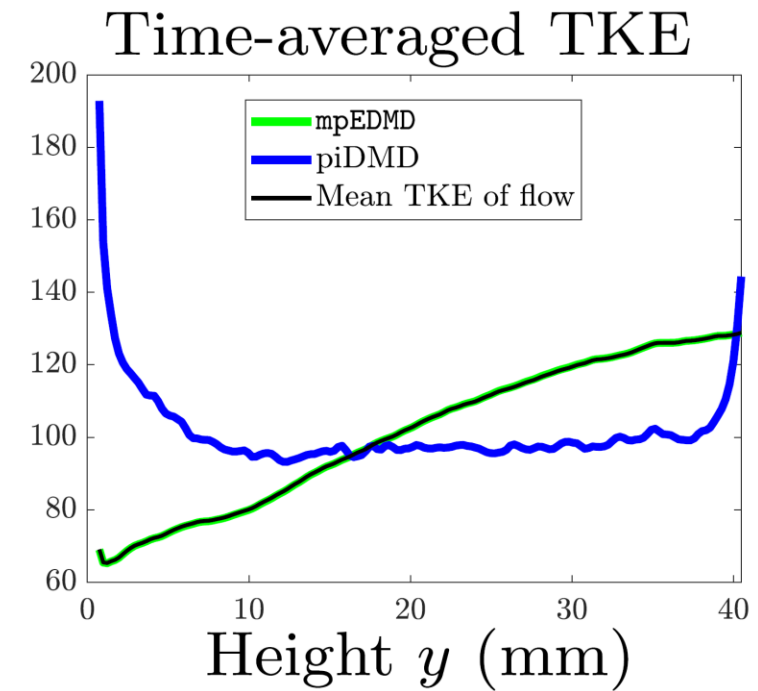
- Polar decomposition of \mathcal{K} . Easy to combine with any DMD-type method!
- Converges for spectral measures, spectra, Koopman mode decomposition.
- Measure-preserving discretization for arbitrary measure-preserving systems.




Snapshots collected over 1s



EDMD unstable!



Solvability Complexity Index Hierarchy

Class $\Omega \ni A$, want to compute $\Xi: \Omega \rightarrow (\mathcal{M}, d)$  metric space

- Δ_0 : Problems solved in finite time (v. rare for cts problems).

- Δ_1 : Problems solved in “one limit” with full error control:

$$d(\Gamma_n(A), \Xi(A)) \leq 2^{-n}$$

- Δ_2 : Problems solved in “one limit”:

$$\lim_{n \rightarrow \infty} \Gamma_n(A) = \Xi(A)$$

- Δ_3 : Problems solved in “two successive limits”:

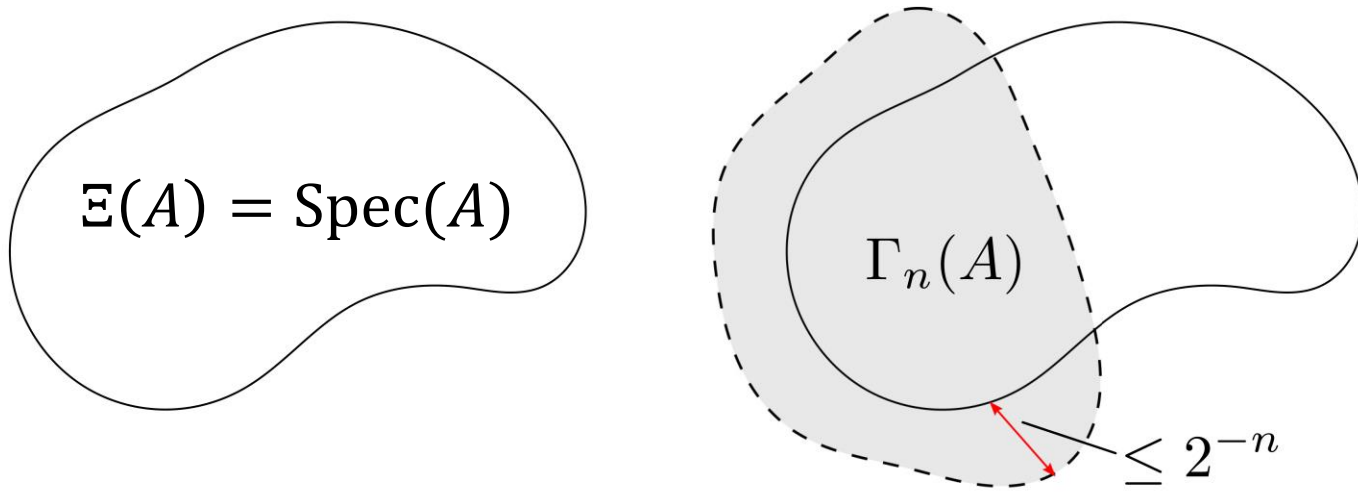
$$\lim_{n \rightarrow \infty} \lim_{m \rightarrow \infty} \Gamma_{n,m}(A) = \Xi(A)$$

⋮

-
- Ben-Artzi, C., Hansen, Nevanlinna, Seidel, “*On the solvability complexity index hierarchy and towers of algorithms*,” preprint.
 - Hansen, “*On the solvability complexity index, the n -pseudospectrum and approximations of spectra of operators*,” **J. Amer. Math. Soc.**, 2011.
 - McMullen, “*Families of rational maps and iterative root-finding algorithms*,” **Ann. of Math.**, 1987.
 - Doyle, McMullen, “*Solving the quintic by iteration*,” **Acta Math.**, 1989.
 - Smale, “*The fundamental theorem of algebra and complexity theory*,” **Bull. Amer. Math. Soc.**, 1981.

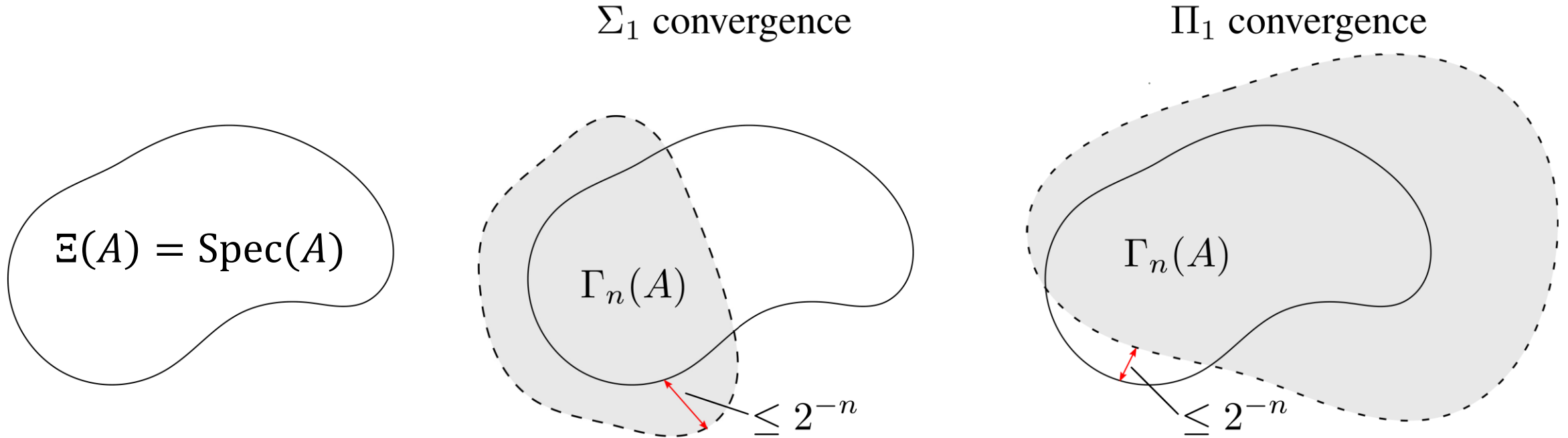
Error control for spectral problems

Σ_1 convergence



- $\Sigma_1: \exists \text{ alg. } \{\Gamma_n\} \text{ s.t. } \lim_{n \rightarrow \infty} \Gamma_n(A) = \Xi(A), \max_{z \in \Gamma_n(A)} \text{dist}(z, \Xi(A)) \leq 2^{-n}$

Error control for spectral problems



- $\Sigma_1: \exists$ alg. $\{\Gamma_n\}$ s.t. $\lim_{n \rightarrow \infty} \Gamma_n(A) = \Xi(A)$, $\max_{z \in \Gamma_n(A)} \text{dist}(z, \Xi(A)) \leq 2^{-n}$
- $\Pi_1: \exists$ alg. $\{\Gamma_n\}$ s.t. $\lim_{n \rightarrow \infty} \Gamma_n(A) = \Xi(A)$, $\max_{z \in \Xi(A)} \text{dist}(z, \Gamma_n(A)) \leq 2^{-n}$

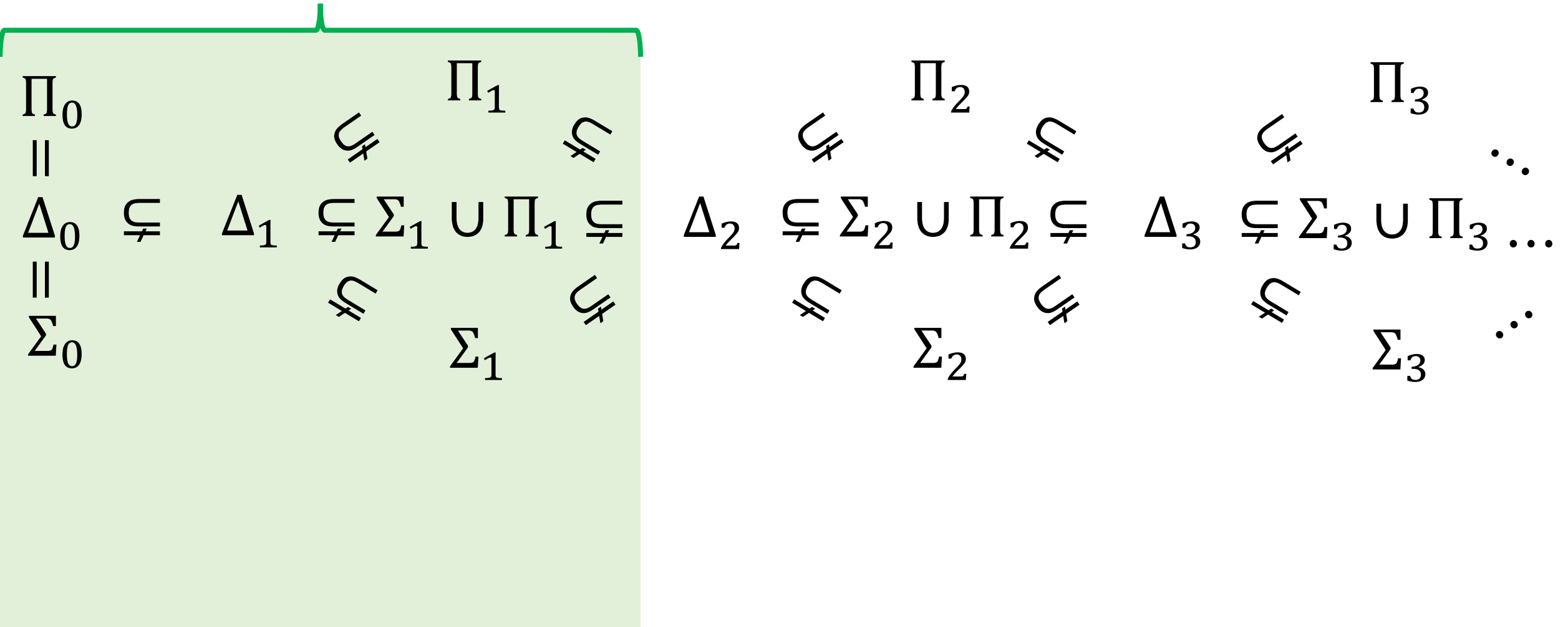
Such problems can be used in a proof!

Small sample of classification theorems

Increasing difficulty



Error control

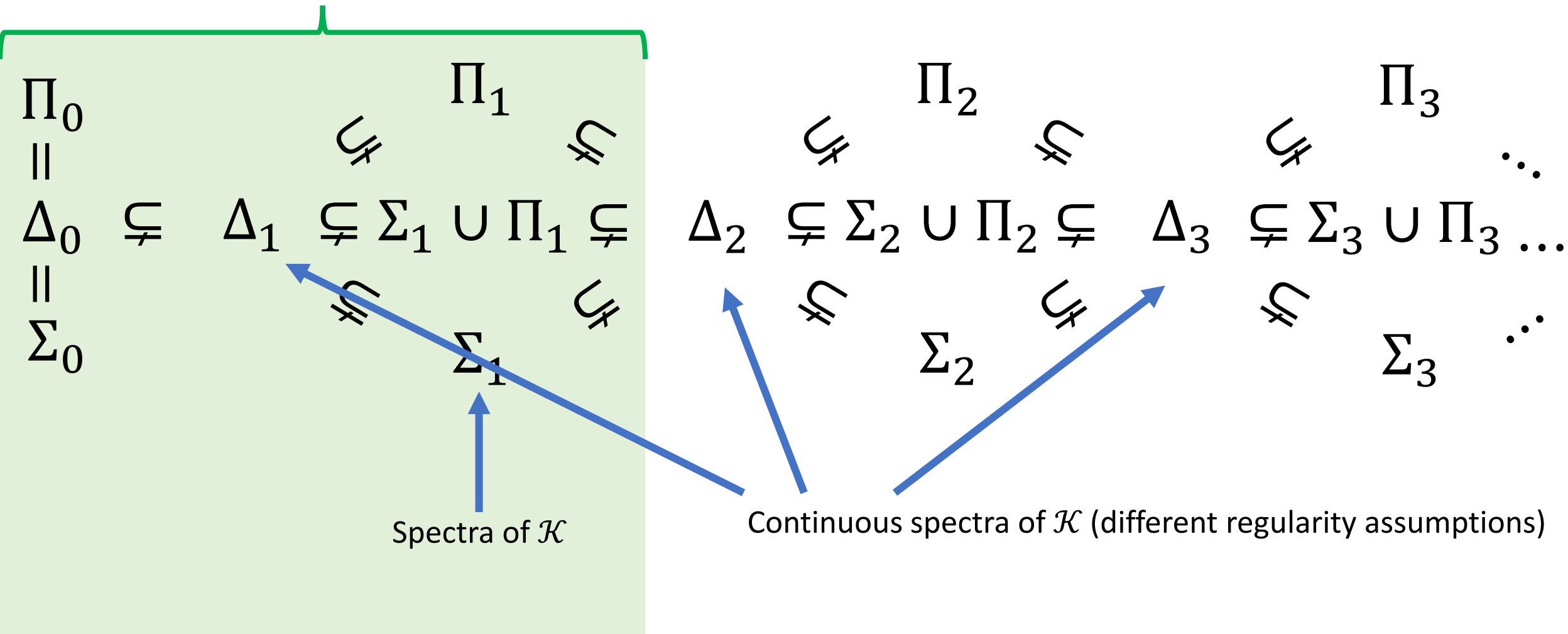


Small sample of classification theorems

Increasing difficulty



Error control



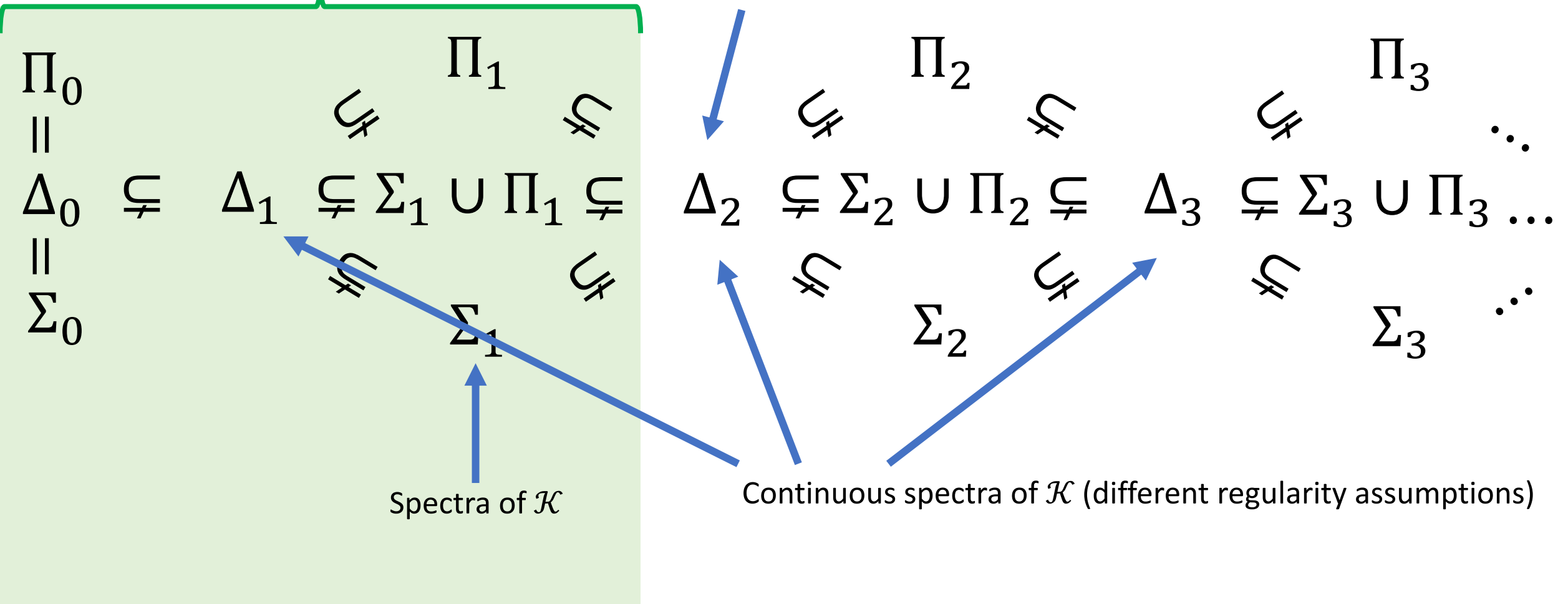
Small sample of classification theorems

Increasing difficulty



Error control

Spectra of compact operators

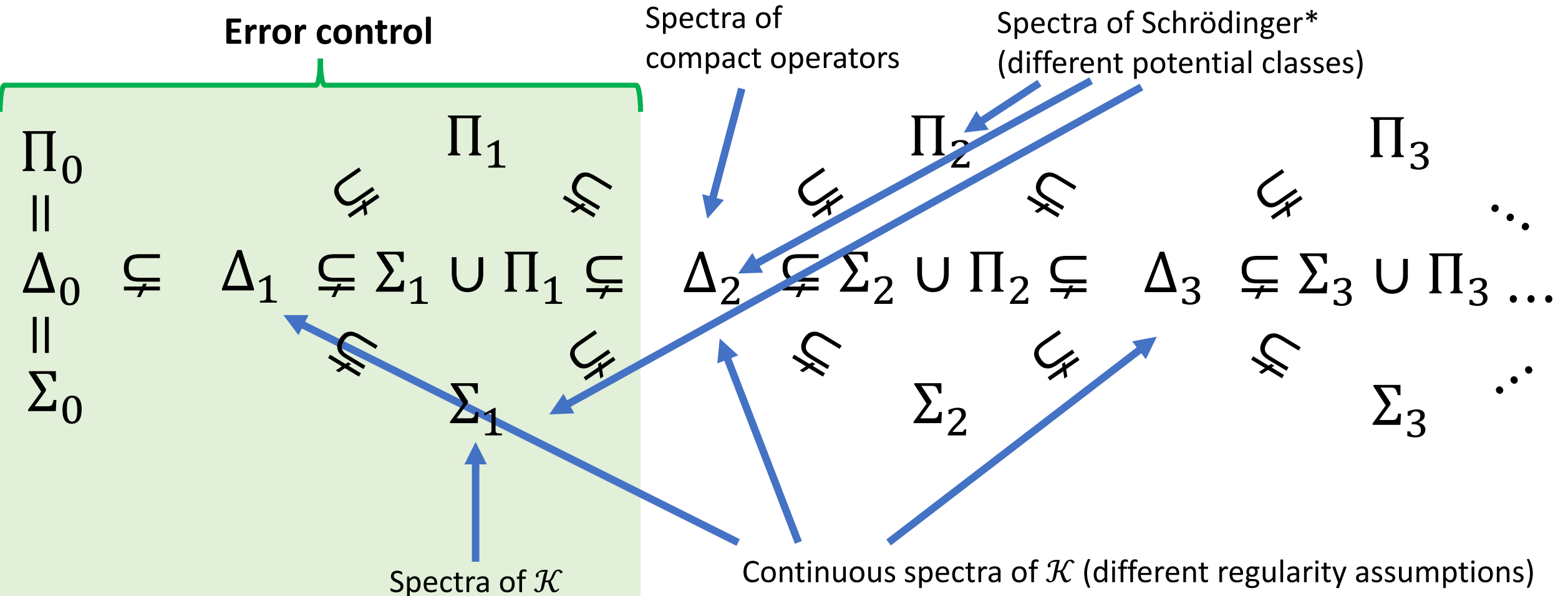


Small sample of classification theorems

Increasing difficulty



Error control



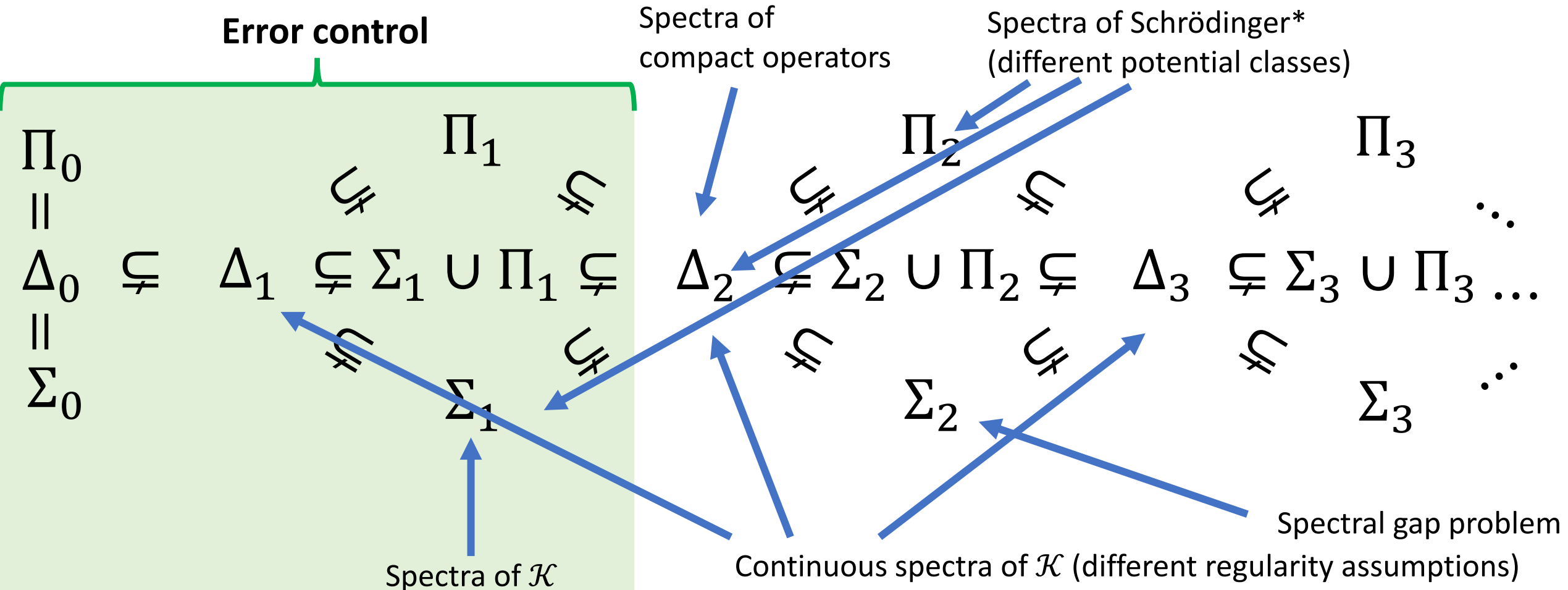
*Open problem of Schwinger: "The special canonical group," "Unitary operator bases," PNAS, 1960.

Small sample of classification theorems

Increasing difficulty



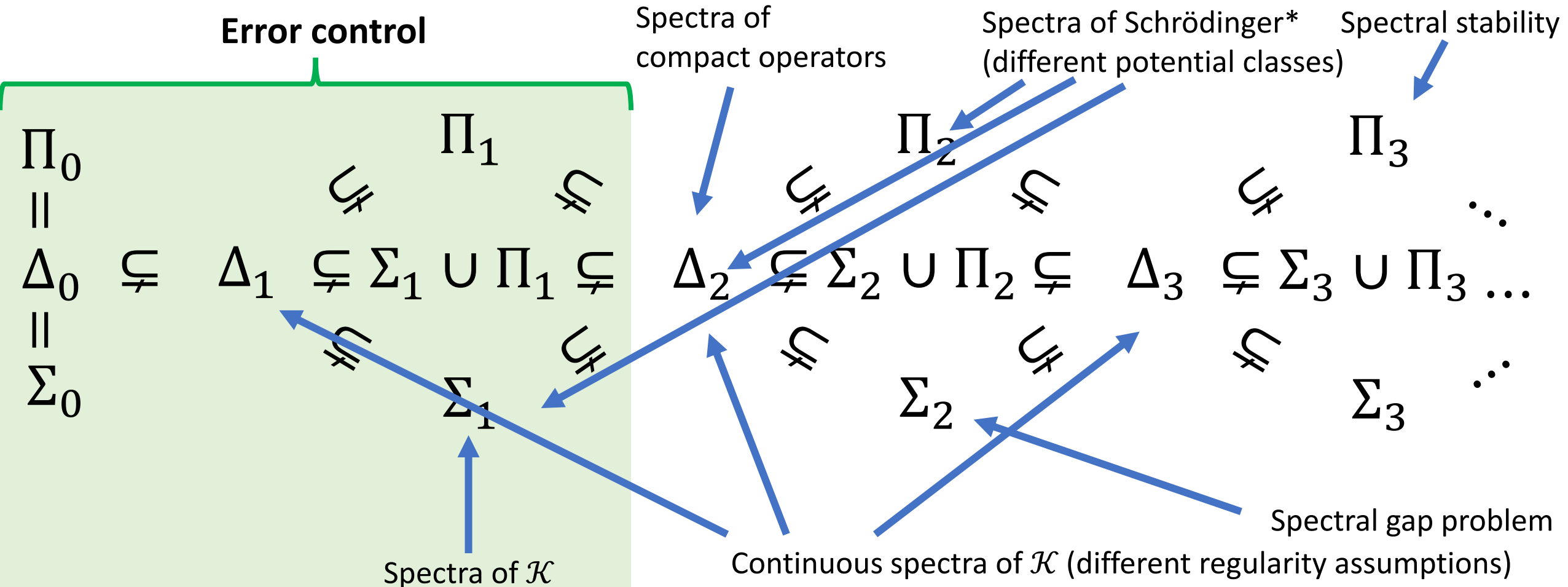
Error control



*Open problem of Schwinger: "The special canonical group," "Unitary operator bases," PNAS, 1960.

Small sample of classification theorems

Increasing difficulty



**Open problem of Schwinger*: “The special canonical group,” “Unitary operator bases,” PNAS, 1960.