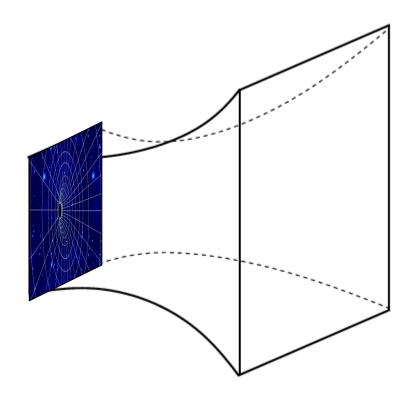
# Optical Conductivity from a Holographic Lattice

**David Tong** 

Based on work with Gary Horowitz and Jorge Santos arXiv:1204.0519



#### Strongly Interacting Stuff



Boundary field theory d=2+1

- Bulk d=3+1 black hole
- Hawking radiation = finite temperature
- Electrically charged (Reissner-Nordstrom) = finite density

**Basics of Conductivity** 

$$\vec{j} = \vec{j}(\omega)e^{-i\omega t}$$

# Obm'de a Wodel:

## Drude Model: AC Conductivity $= \frac{1}{\rho} \frac{1}{1-i\omega\tau}$

Fourier Transform: 
$$\vec{j}(\omega) = \sigma(\omega) \vec{E}(\omega) \quad \vec{j}(\omega) = \vec{j}(\omega) \vec{e}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{e}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) = \vec{j}(\omega) \vec{j}(\omega) = \vec{j}($$

$$\sigma(\omega) = \frac{1}{\rho} \frac{1}{1 - i\omega\tau} \qquad \sigma(\omega) = \frac{1}{\rho} \frac{1}{1 - i\omega\tau}$$

$$ec{E} = N ec{p}(\omega) e^{-i\omega t} + i/\omega \quad \text{as} \quad \omega o \infty$$
 
$$ec{j} = ec{j}(\omega) e^{-i\omega t} \qquad \qquad \sigma(\omega) o i/\omega \quad \text{as} \quad \omega o \infty$$

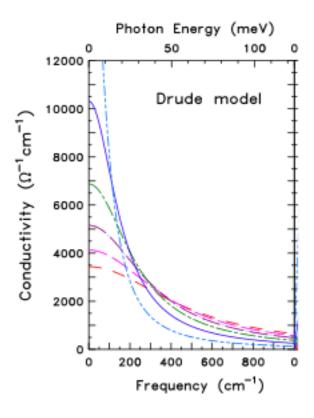
#### **Drude Model**

$$m\frac{d\vec{v}}{dt} + \frac{m}{\tau}\vec{v} = q\vec{E}$$

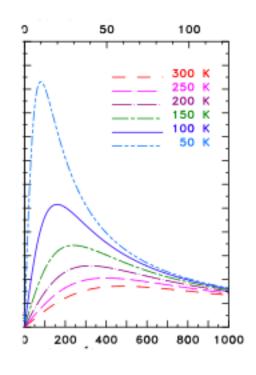
$$\vec{j} = nq\vec{v}$$

$$\vec{j} = nq\vec{v}$$
  $\Longrightarrow$   $\sigma(\omega) = \left(\frac{nq^2\tau}{m}\right)\frac{1}{1 - i\omega\tau}$ 

#### **Drude Model**

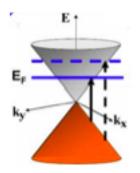


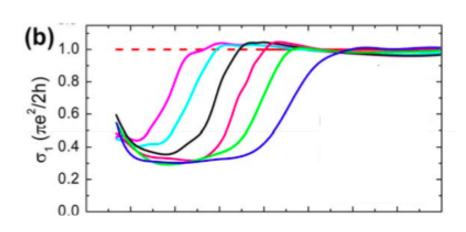
$$Re(\sigma) = \frac{\sigma_0}{1 + \omega^2 \tau^2}$$

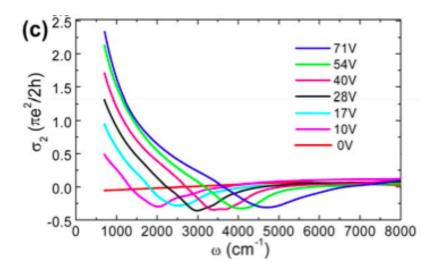


$$Im(\sigma) = \frac{\sigma_0 \omega \tau}{1 + \omega^2 \tau^2}$$

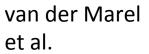
## Graphene

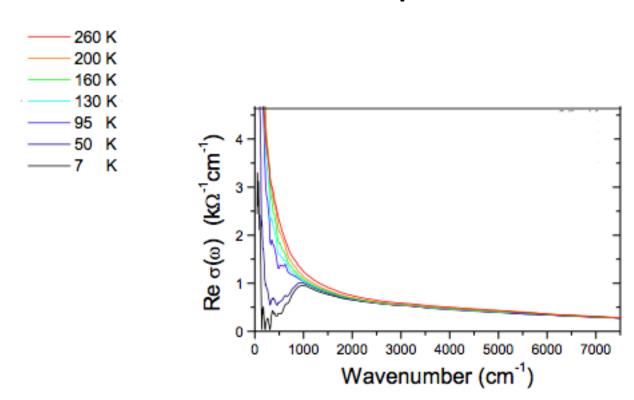






## Cuprates

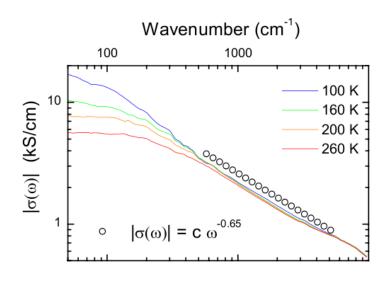


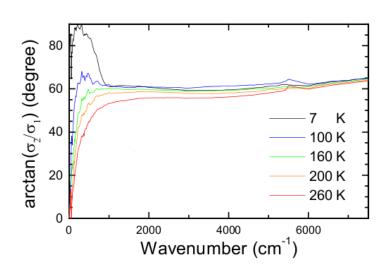


 $Bi_2Sr_2Ca_{0.92}Y_{0.08}Cu_2O_{8+\delta}$ 

#### Cuprates

van der Marel et al.

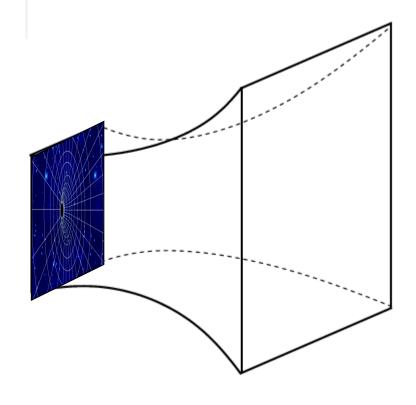




 $Bi_2Sr_2Ca_{0.92}Y_{0.08}Cu_2O_{8+\delta}$ 

**Back to Holography** 

#### Computing the Conductivity



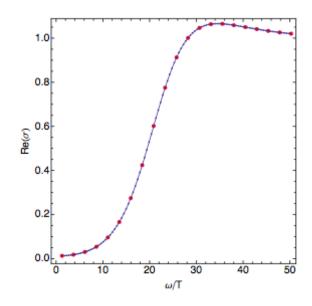
Boundary field theory

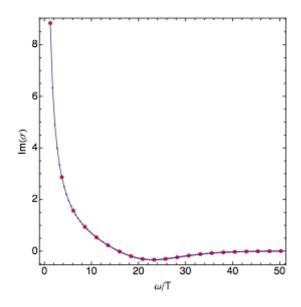
$$z = 0$$

$$A_x = \frac{E}{\omega} + J_x z + \dots$$

### **Optical Conductivity**

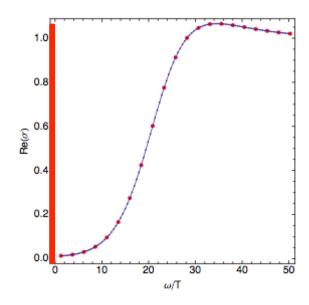
$$j(\omega) = \sigma(\omega)E(\omega)$$



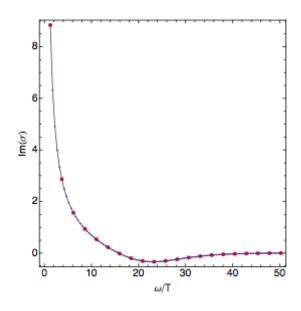


## **Optical Conductivity**

$$j(\omega) = \sigma(\omega)E(\omega)$$



$$\operatorname{Re} \sigma(\omega) \sim K \delta(\omega)$$



$$\operatorname{Im}(\sigma) \to \frac{K}{\pi\omega}$$

#### Resolving the Delta-Function

$$\operatorname{Re} \sigma(\omega) \sim K \delta(\omega)$$

- Due to: Finite density of charge carriers
  - Translational invariance

Two options to resolve

- 1. Dilute charge carriers
- 2. Break translational invariance

#### **Option 1: Diluting Charge Carries**

#### **Probe Branes:**

DC Resistivity:

$$\rho \sim T^{2/z}$$

Optical Conductivity:

$$\sigma(\omega) \to \begin{cases} (i\omega)^{-1} & z < 2\\ (i\omega)^{-2/z} & z > 2 \end{cases}$$

Karch and O'Bannon; Hartnoll, Silverstein, Polchinski, Tong **Breaking Translational Invariance** 

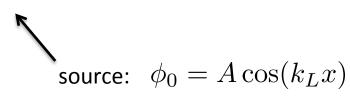
#### How to Build a Lattice

$$\mathcal{L} = \mathcal{L}_{CFT} + \mu \mathcal{Q} + \phi_0(x, y)\mathcal{O}$$

Introduce a neutral, bulk scalar field:  $\Phi \longleftrightarrow \mathcal{O}$ 

Pick  $m_\Phi^2 L^2 = -1 \Longrightarrow$  relevant operator:

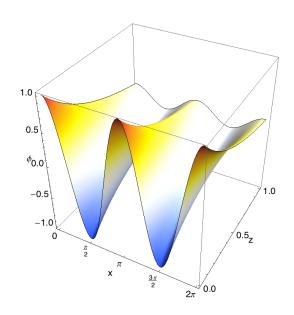
$$\Phi \to z\phi_0 + z^2\phi_1 + \dots$$



#### The Lattice

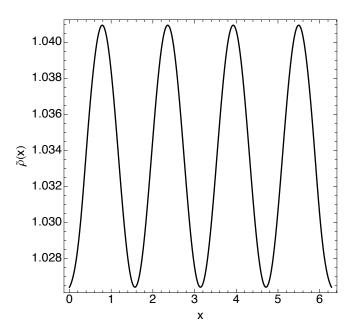
Solve Einstein equations subject to lattice boundary conditions:  $A_0(z,x)$  ,  $\Phi(x,z)$  and

$$ds^{2} = \frac{L^{2}}{z^{2}} \left[ -g_{tt}(z,x)dt^{2} + g_{zz}(z,x)dz^{2} + g_{xx}(z,x)(dx + a(z,x)dz)^{2} + g_{yy}(z,x)dy^{2} \right]$$



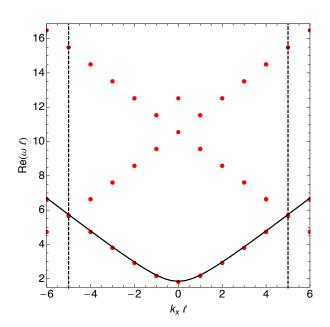
Note:  $\mathcal{O}$  relevant;  $e^{ik_Lx}\mathcal{O}$  irrelevant

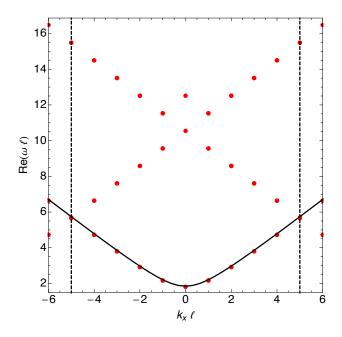
## **Charge Density**



#### **Band Structure**

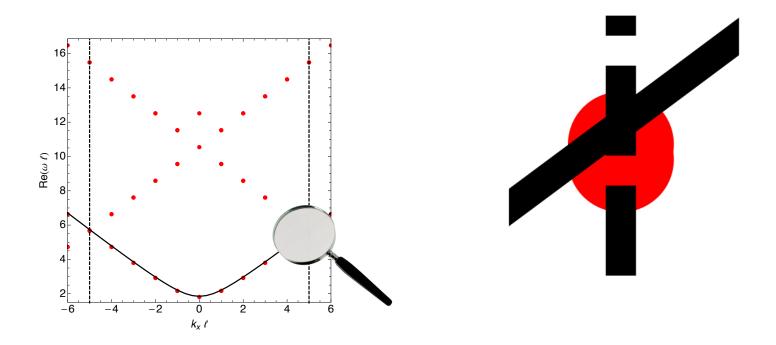
Add a probe scalar field in the lattice background





#### **Band Structure**

Add a probe scalar field in the lattice background



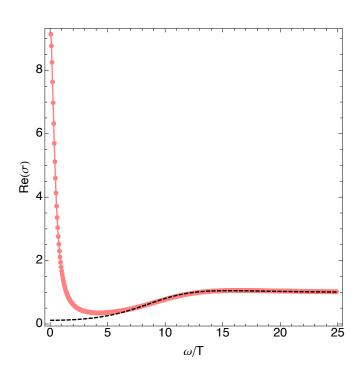
Band gaps negligible

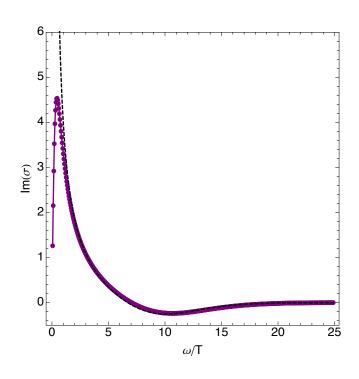
#### Perturbing the Lattice

$$\delta g_{tt}, \ \delta g_{tz}, \ \delta g_{tx}, \ \delta g_{zz}, \ \delta g_{zx}, \ \delta g_{xx}, \ \delta g_{yy}$$
 
$$\delta A_t, \ \delta A_z, \ \delta A_x, \ \delta \Phi$$

### **Optical Conductivity**

$$j(\omega) = \sigma(\omega)E(\omega)$$

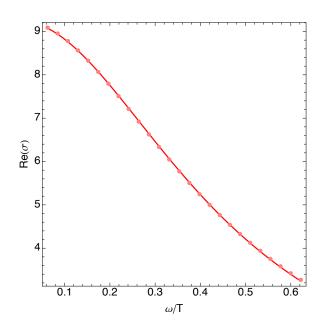


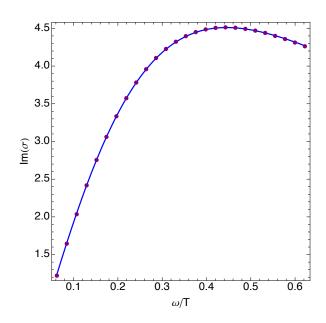


$$\mu = 1.4$$
  $T = 0.115\mu$   $k_L = 2$   $A = 1.5$ 

#### Low Frequency Behaviour



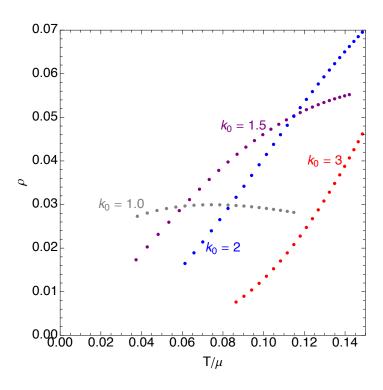




$$\sigma(\omega) = \frac{K\tau}{1 - i\omega\tau}$$

Note: No quasiparticles

#### **DC** Resistivity



$$\rho = \frac{1}{K\tau}$$

#### DC Resistivity

Low energy excitations around black hole governed by *locally critical* theory

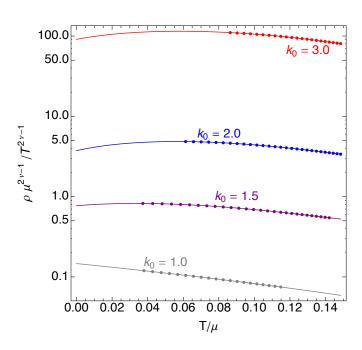
- Field Theory:  $z \to \infty$
- Geometry:  $AdS_2 \times \mathbf{R}^2$

Hartnoll and Hofman:

$$\rho \sim T^{2\nu - 1}$$

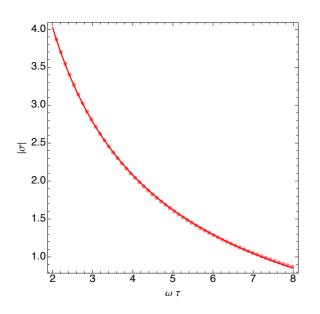
$$\nu = \frac{1}{2}\sqrt{5 + 2(k/\mu)^2 - 4\sqrt{1 + (k/\mu)^2}}$$

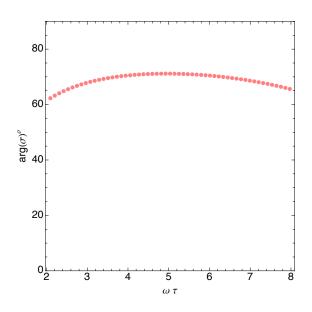
## **DC** Resistivity



## Mid-Frequency Behaviour

$$\omega \gtrsim T$$
$$2 < \omega \tau < 8$$

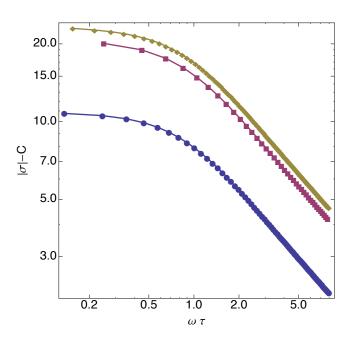




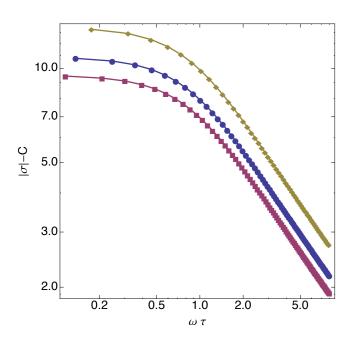
$$|\sigma(\omega)| = \frac{B}{\omega^{2/3}} + C$$

#### **Robust Power-Law**

#### Log-log plots



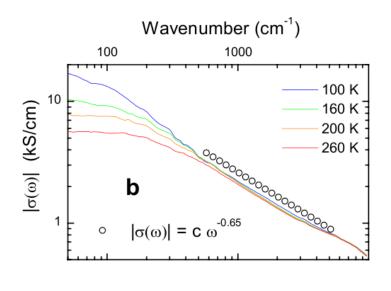
$$k_L = 3, k_L = 1, k_L = 2$$

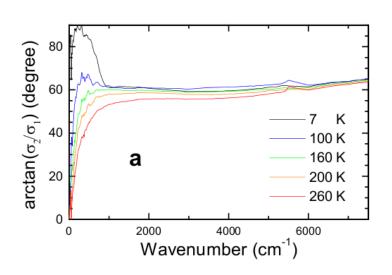


$$T = 0.98\mu, T = 0.115\mu, T = 0.13\mu$$

#### Comparison to Cuprates

van der Marel et al.





 $Bi_2Sr_2Ca_{0.92}Y_{0.08}Cu_2O_{8+\delta}$ 

Note: No offset, C

## Summary

Why?