

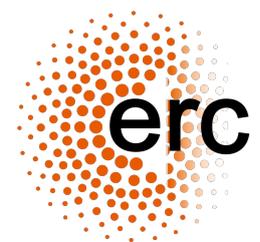
On Particle-Vortex Duality

David Tong

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Based on work with Andreas Karch

Also related work by Seiberg, Senthil, Witten and Wang



Particle-Vortex Duality for Bosons

Theory A:

$$S = \int d^3x |(\partial_\mu - iA_\mu)\phi|^2 - V(\phi)$$

Theory B:

$$S = \int d^3x |(\partial_\mu - ia_\mu)\Phi|^2 - \tilde{V}(\Phi) + \frac{1}{2\pi} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu a_\rho$$

dynamical gauge field

background gauge field

Boson \longleftrightarrow Monopole operator $\int f = 2\pi$

Particle-Vortex Duality for Fermions

Theory A:

$$S = \int d^3x \, i\bar{\psi}\gamma^\mu(\partial_\mu - iA_\mu)\psi$$

Theory B:

$$S = \int d^3x \, i\bar{\Psi}\gamma^\mu(\partial_\mu - ia_\mu)\Psi + \frac{1}{4\pi}\epsilon^{\mu\nu\rho}A_\mu\partial_\nu a_\rho$$

$$\text{Fermion} \longleftrightarrow \text{Monopole operator} \int f = 4\pi$$

Parity anomaly means that theory really makes sense only on boundary of topological insulator

Some notation

$$S_{\text{scalar}}[\phi; A] = \int d^3x |(\partial_\mu - iA_\mu)\phi|^2 + \dots$$

$$\text{and } Z_{\text{scalar}}[A] = \int \mathcal{D}\phi \exp\left(iS_{\text{scalar}}[A]\right)$$

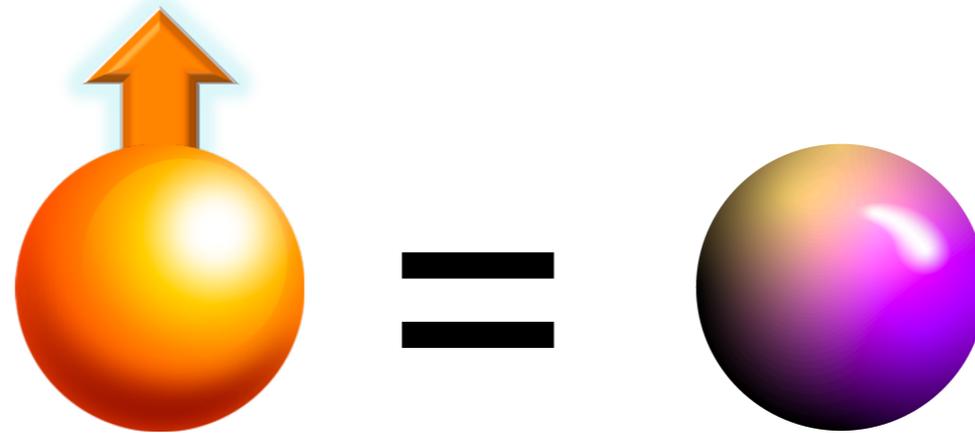
$$S_{\text{fermion}}[\psi; A] = \int d^3x i\bar{\psi}\gamma^\mu(\partial_\mu - iA_\mu)\psi + \dots$$

$$\text{and } Z_{\text{fermion}}[A] = \int \mathcal{D}\psi \exp\left(iS_{\text{fermion}}[A]\right)$$

$$S_{CS}[a] = \frac{1}{4\pi} \int d^3x \epsilon^{\mu\nu\rho} a_\mu \partial_\nu a_\rho$$

$$S_{BF}[a; A] = \frac{1}{2\pi} \int d^3x \epsilon^{\mu\nu\rho} a_\mu \partial_\nu A_\rho$$

Something Simple...Flux Attachment



Boson + flux

Fermion

Wilczek, Jain and many others

Relativistic Flux Attachment

$$S_{\text{scalar+flux}} = \int d^3x \left| (\partial_\mu - ia_\mu)\phi \right|^2 + \frac{1}{4\pi} \epsilon^{\mu\nu\rho} a_\mu \partial_\nu a_\rho + \frac{1}{2\pi} \epsilon^{\mu\nu\rho} A_\mu \partial_\nu a_\rho$$

When we turn off the background field A , the equation of motion for a attaches one unit of flux

$$\rho_{\text{scalar}} + \frac{f}{2\pi} = 0$$

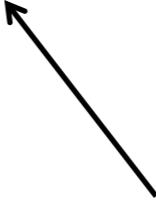
This is a fermion

A Seed Duality

Theory A: $S_{\text{scalar+flux}} = S_{\text{scalar}}[\phi; a] + S_{CS}[a] + S_{BF}[a; A]$

Theory B: $S = S_{\text{fermion}}[A] - \frac{1}{2}S_{CS}[A]$

ensures Hall conductivities agree



Polyakov '88 and many others

(More recently, many works on non-Abelian versions of this duality, notably by Minwalla et. al. and Aharony et.al.)

A Seed Duality

Theory A: $S_{\text{scalar+flux}} = S_{\text{scalar}}[\phi; a] + S_{CS}[a] + S_{BF}[a; A]$

Theory B: $S = S_{\text{fermion}}[A] - \frac{1}{2}S_{CS}[A]$

$$Z_{\text{scalar+flux}}[A] = \int \mathcal{D}a Z_{\text{scalar}}[a] \exp\left(iS_{CS}[a] + iS_{BF}[a; A]\right) = Z_{\text{fermion}}[A] e^{-\frac{i}{2}S_{CS}[A]}$$

Basic Idea

$$Z_{\text{scalar+flux}}[A] = \int \mathcal{D}a \, Z_{\text{scalar}}[a] \exp\left(iS_{CS}[a] + iS_{BF}[a; A]\right) = Z_{\text{fermion}}[A] e^{-\frac{i}{2}S_{CS}[A]}$$

Manipulate path integral to derive other dualities

- c.f. Kapustin and Strassler '99 for supersymmetric theories

Example 1: Another Flux Attachment

$$Z_{\text{scalar+flux}}[A] = \int \mathcal{D}a \, Z_{\text{scalar}}[a] \exp\left(iS_{CS}[a] + iS_{BF}[a; A]\right) = Z_{\text{fermion}}[A] e^{-\frac{i}{2}S_{CS}[A]}$$

Promote A to dynamical gauge field. Right-hand side becomes

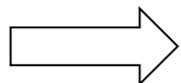
$$Z_{\text{fermion+flux}}[C] = \int \mathcal{D}A \, Z_{\text{fermion}}[A] \exp\left(-\frac{i}{2}S_{CS}[A] - iS_{BF}[A; C]\right)$$

Do the same on left-hand side.

new background field

$$\int \mathcal{D}A \, Z_{\text{scalar+flux}}[A] e^{-iS_{BF}[A; C]} = Z_{\text{scalar}}[C] e^{iS_{CS}[C]}$$

integrate out A ; equation of motion
also eliminates a through $da = dC$



$$Z_{\text{fermion+flux}}[C] = Z_{\text{scalar}}[C] e^{iS_{CS}[C]}$$

Chen, Fisher, Wu '93

Example 2: Bosonic Particle-Vortex Duality

$$Z_{\text{fermion+flux}}[C] = Z_{\text{scalar}}[C] e^{iS_{CS}[C]} \quad \Longrightarrow \quad Z_{\text{fermion+flux}}[C] e^{-iS_{CS}[C]} = Z_{\text{scalar}}[C]$$

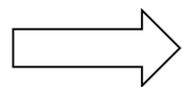
Promote C to a dynamical gauge field. Right-hand side becomes

$$Z_{\text{scalar-QED}}[A] = \int \mathcal{D}C \, Z_{\text{scalar}}[C] e^{iS_{BF}[C;A]}$$

On the left-hand-side something nice happens. After integrating out C , we find

$$\int \mathcal{D}a \, Z_{\text{fermion}}[a] \exp\left(\frac{i}{2}S_{CS}[a] - iS_{BF}[a;A] + iS_{CS}[A]\right)$$

But this is just our original “fermion+flux” partition function (or, more precisely, a time-reversed version)



$$Z_{\text{scalar-QED}}[A] = Z_{\text{scalar}}[A]$$

Example 3: Fermionic Particle-Vortex Duality

$$Z_{\text{scalar+flux}}[C] = Z_{\text{fermion}}[C] e^{-\frac{i}{2}S_{CS}[C]} \implies Z_{\text{scalar+flux}}[C] e^{+\frac{i}{2}S_{CS}[C]} = Z_{\text{fermion}}[C]$$

Again, promote C to a dynamical gauge field, but now restricted to have even flux sector. Left-hand side is

$$Z_{\text{QED}}[A] = \int \mathcal{D}A Z_{\text{fermion}}[A] e^{\frac{i}{2}S_{BF}[a;A]}$$

On the left-hand-side something nice happens. After integrating out C , we find

$$\int \mathcal{D}a Z_{\text{scalar}}[a] \exp\left(-iS_{CS}[a] - iS_{BF}[a;A] - \frac{i}{2}S_{CS}[A]\right)$$

This is again the time-reversed version of our original scalar+flux. We can then use duality to derive



$$Z_{\text{QED}}[A] = Z_{\text{fermion}}[A]$$

Example 4: A Self-Dual Theory

- Take two copies of the scalar+flux = fermion duality.
- Gauge an overall symmetry to get $N_f=2$ QED

$$Z_{\text{QED}[N_f=2]}[A; C] = \int \mathcal{D}a \, Z_{\text{fermion}}[a + C] Z_{\text{fermion}}[a - C] e^{+iS_{BF}[a; A]}$$

Two background gauge fields: A for the topological current, and C for the Cartan of the $SU(2)$ flavour

The dual theory of 2 scalars takes the schematic form

$$\int \mathcal{D}\phi_1 \mathcal{D}\phi_2 \mathcal{D}c_+ \exp \left(iS[\phi_1, \phi_2, c_+; A + C] - \frac{i}{2} S_{BF}[c_+; A - C] \right)$$

But this is invariant under swapping A and C , together with time-reversal



$$Z_{\text{QED}[N_f=2]}[A; C] = \bar{Z}_{\text{QED}[N_f=2]}[C; A]$$



Time reversed theory

Summary

Many dualities follow from relativistic version of flux attachment