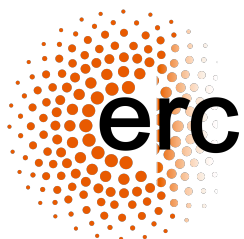


What is the Gauge Group of the Standard Model?

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The Standard Model

$$\tilde{G} = U(1) \times SU(2) \times SU(3)$$

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All standard model fields are invariant under \mathbf{Z}_6 generated by

$$\eta = e^{2\pi i Y/6} \otimes \begin{pmatrix} -1 & & \\ & -1 & \\ & & -1 \end{pmatrix} \otimes \begin{pmatrix} \omega & & \\ & \omega & \\ & & \omega \end{pmatrix} \quad \omega^3 = 1$$

i.e.

$$\begin{array}{lll} & e_R : & (\mathbf{1}, \mathbf{1})_{-6} \\ l_L : & (\mathbf{2}, \mathbf{1})_{-3} & d_R : (\mathbf{1}, \mathbf{3})_{-2} \quad H : (\mathbf{2}, \mathbf{1})_3 \\ q_L : & (\mathbf{2}, \mathbf{3})_{+1} & u_R : (\mathbf{1}, \mathbf{3})_{+4} \end{array}$$

The Standard Model

The gauge group of the standard model is

$$G = \frac{U(1) \times SU(2) \times SU(3)}{\Gamma}$$

with $\Gamma = \mathbf{Z}_6, \mathbf{Z}_3, \mathbf{Z}_2$ or $\mathbf{1}$

Which describes our world? How can we tell?

A Warm Up: $SU(N)$ vs $SU(N)/\mathbf{Z}_N$

Following Aharony, Seiberg and Tachikawa, 2013

Line Operators in $SU(N)$ and $SU(N)/\mathbf{Z}_N$

- Wilson line operators: probe electric particles
 - Labeled by transformation under center of $SU(N)$

$$z^e = 0, 1, \dots, N - 1$$

- 't Hooft line operators: probe magnetic particles
 - Also labeled by transformation under center

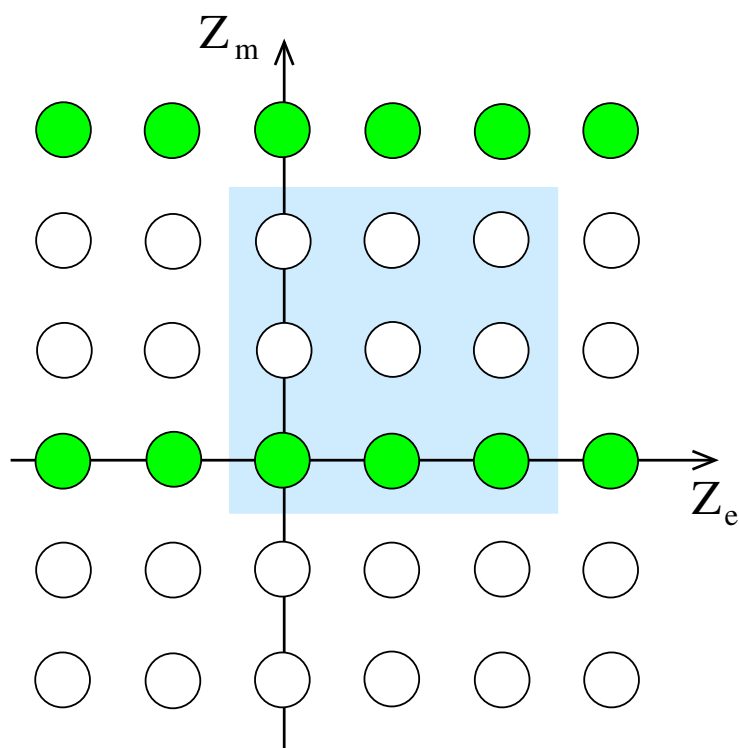
$$z^m = 0, 1, \dots, N - 1$$

- Dirac quantisation requires

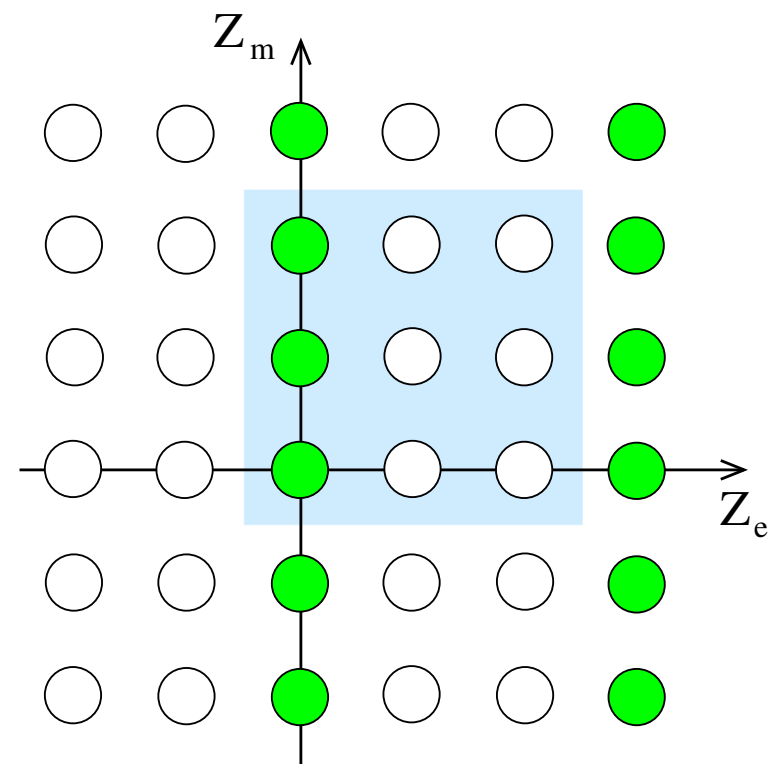
$$z^e z'^m - z^m z'^e = 0 \mod N$$

Line Operators in $SU(N)$ and $SU(N)/\mathbf{Z}_N$

$SU(3)$



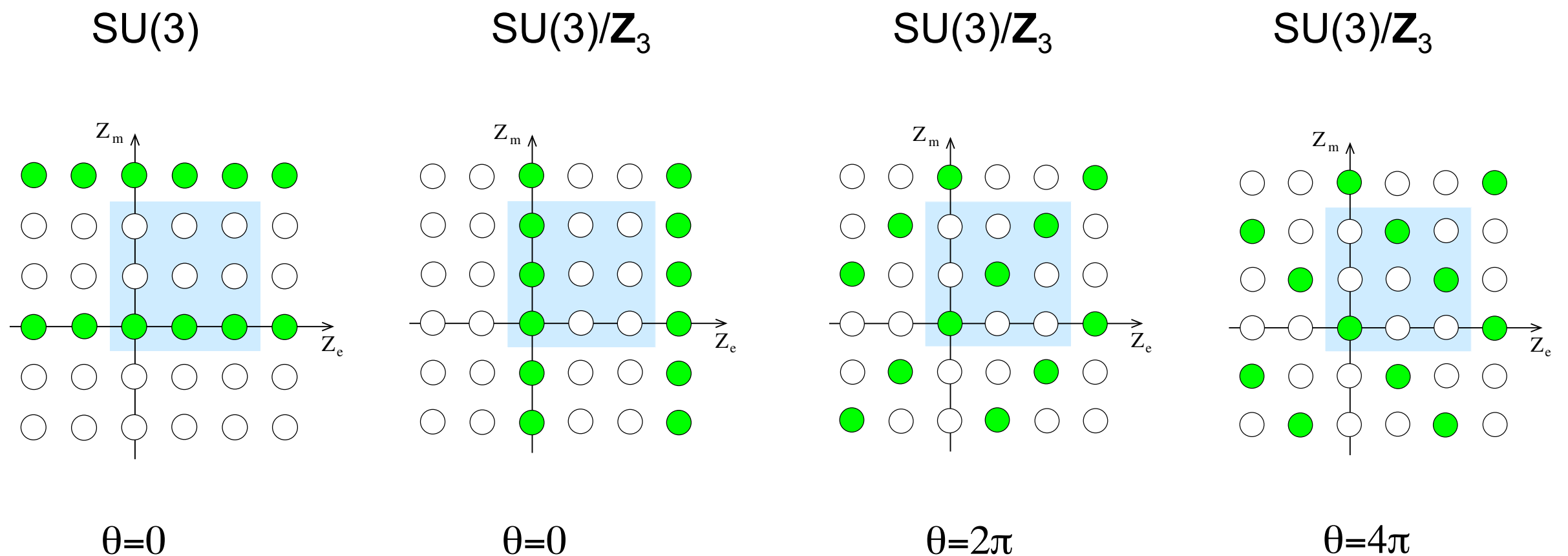
$SU(3)/\mathbf{Z}_3$



Periodicity of Theta Angle for $SU(N)$ and $SU(N)/\mathbf{Z}_N$

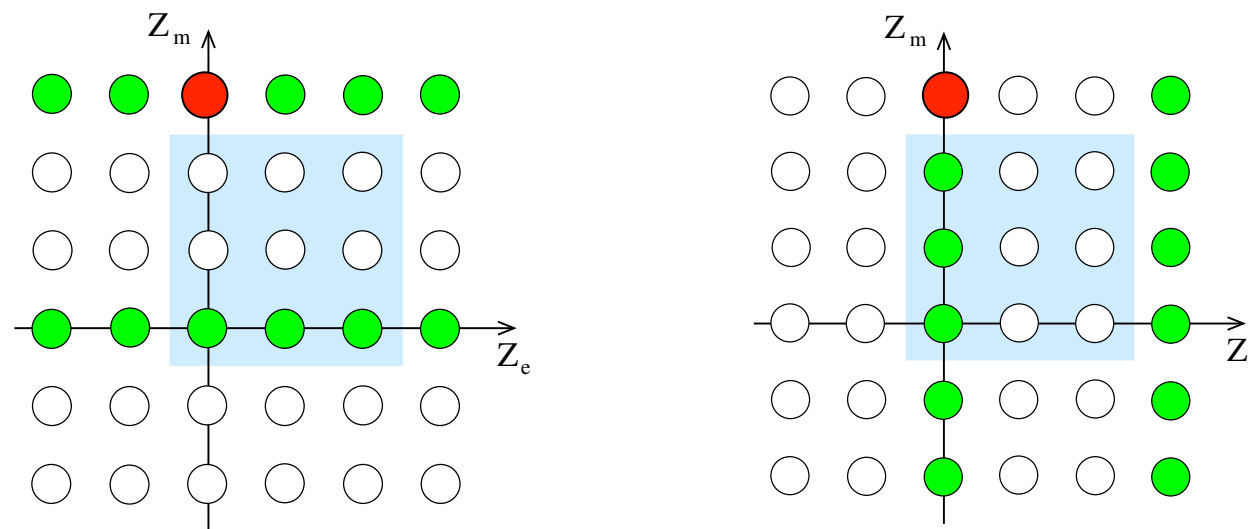
- For $SU(N)$, θ has periodicity 2π
- For $SU(N)/\mathbf{Z}_N$, θ has periodicity $2\pi N$

This can be seen by looking at the Witten effect on line operators



Emergent Magnetic \mathbf{Z}_N Symmetry

Both $SU(N)$ and $SU(N)/\mathbf{Z}_N$ confine through condensation of magnetic monopoles



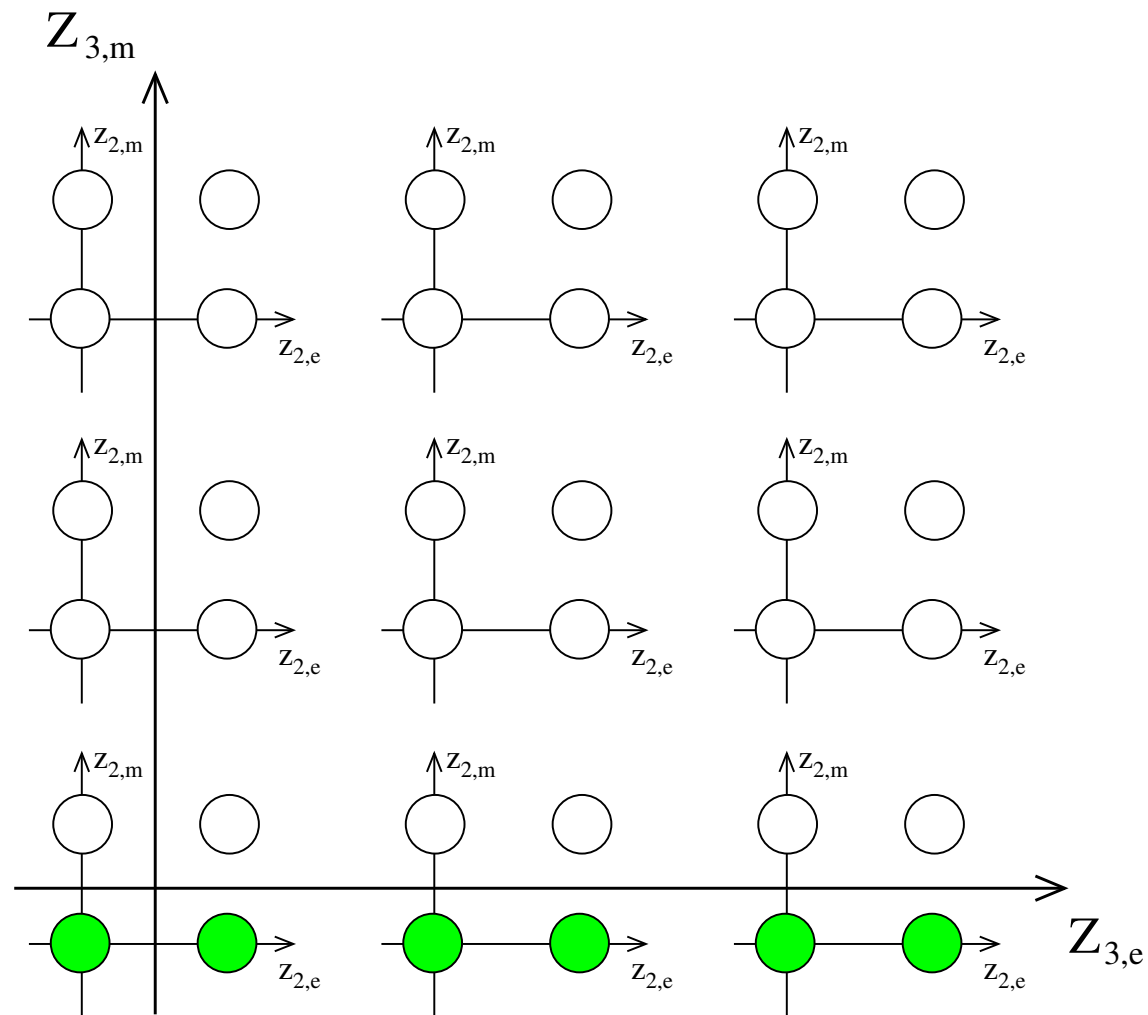
- For $SU(N)/\mathbf{Z}_N$ this is *not* the minimal monopole.
- There is an emergent \mathbf{Z}_N magnetic gauge symmetry.
 - This changes the local dynamics when the theory is compactified on, say, \mathbf{S}^1

Line Operators in the Standard Model

The spectrum of line operators depends on the choice of $\Gamma = \mathbf{Z}_6, \mathbf{Z}_3, \mathbf{Z}_2$ or $\mathbf{1}$

Line Operators

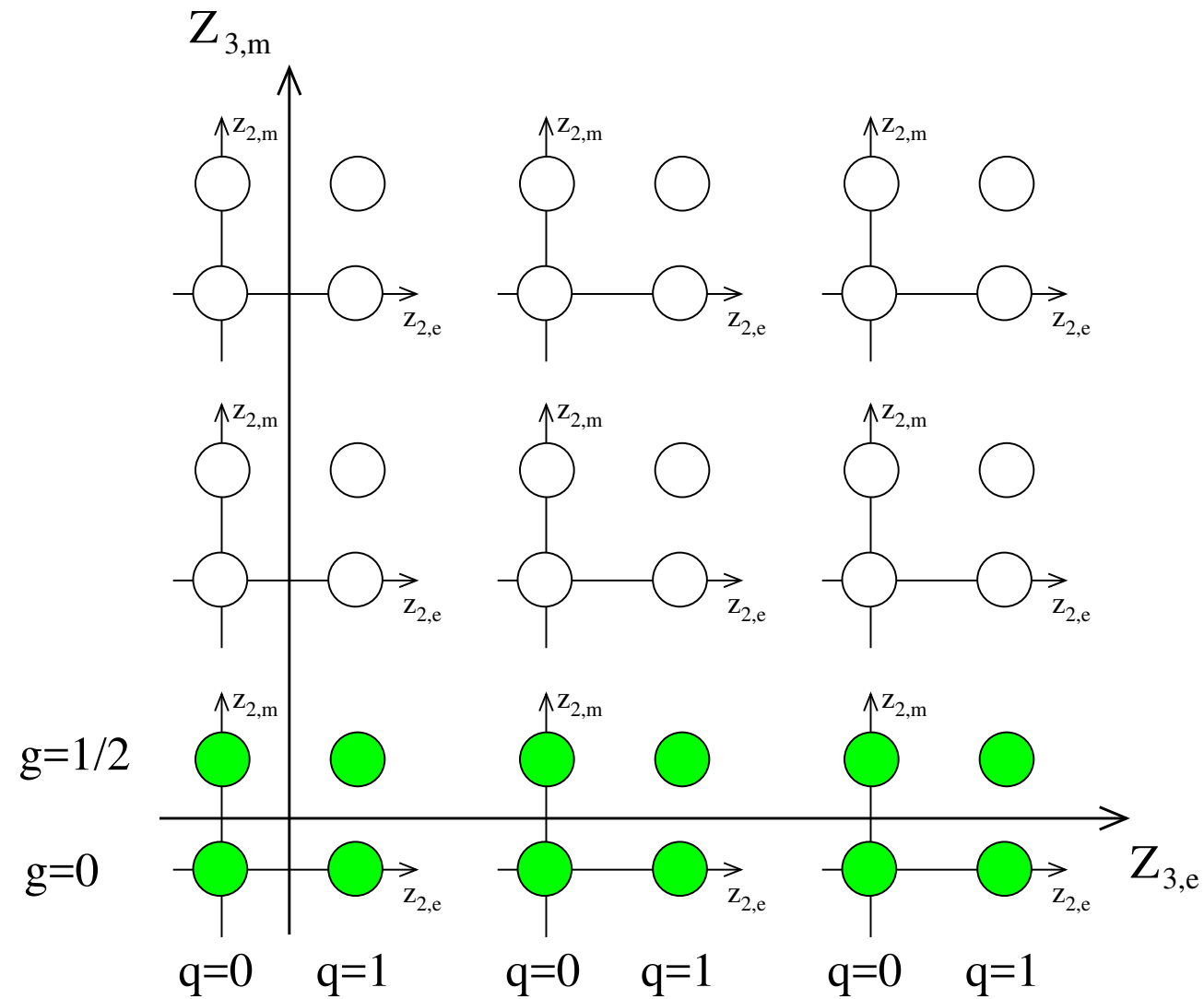
$\Gamma = 1$:



Together with U(1) line operators with $(q,g) = (1,0)$ and $(0,1)$

Line Operators

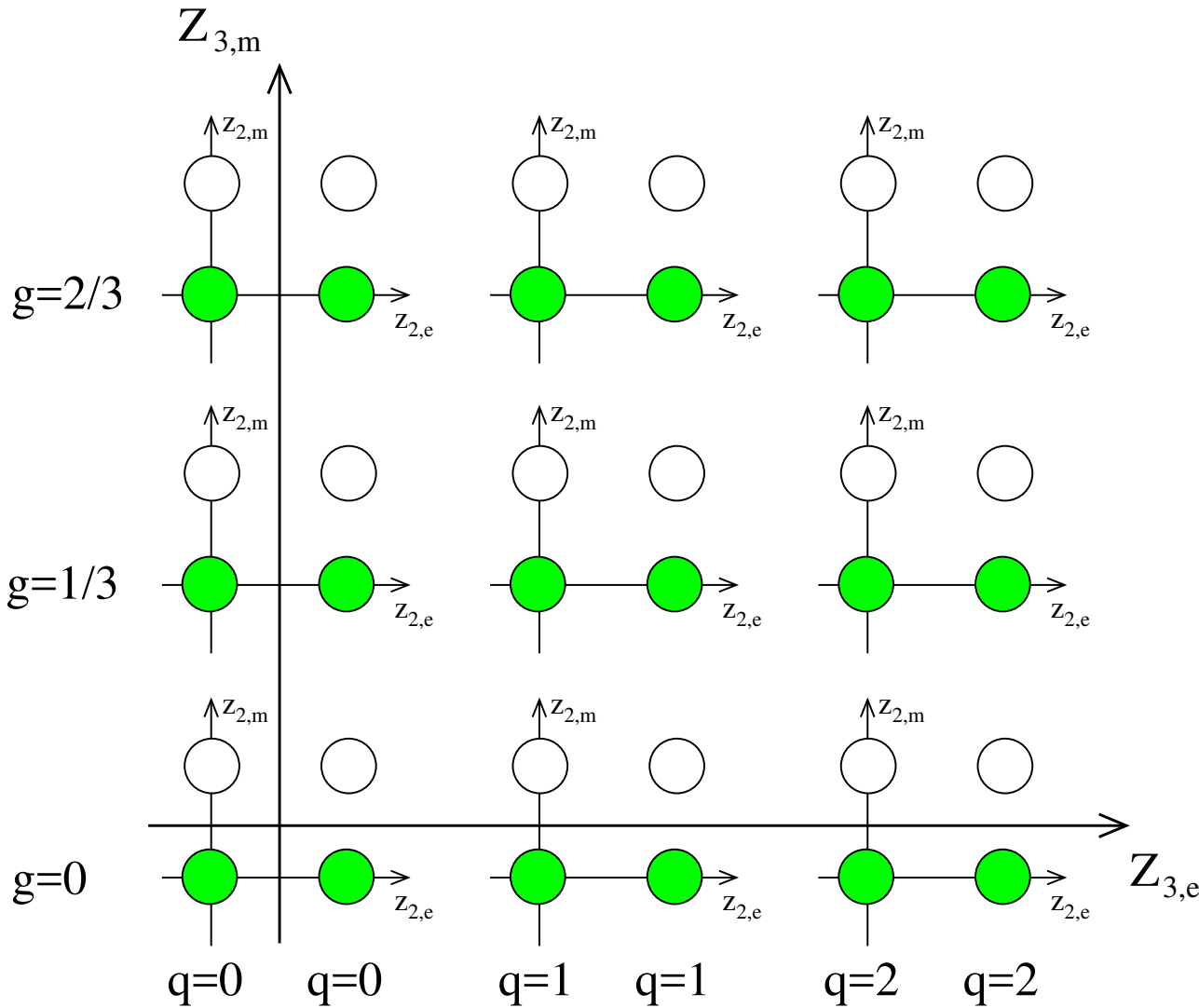
$$\underline{\Gamma = \mathbf{Z}_2:}$$



Together with $U(1)$ line operators with $(q,g) = (2,0)$ and $(0,1)$

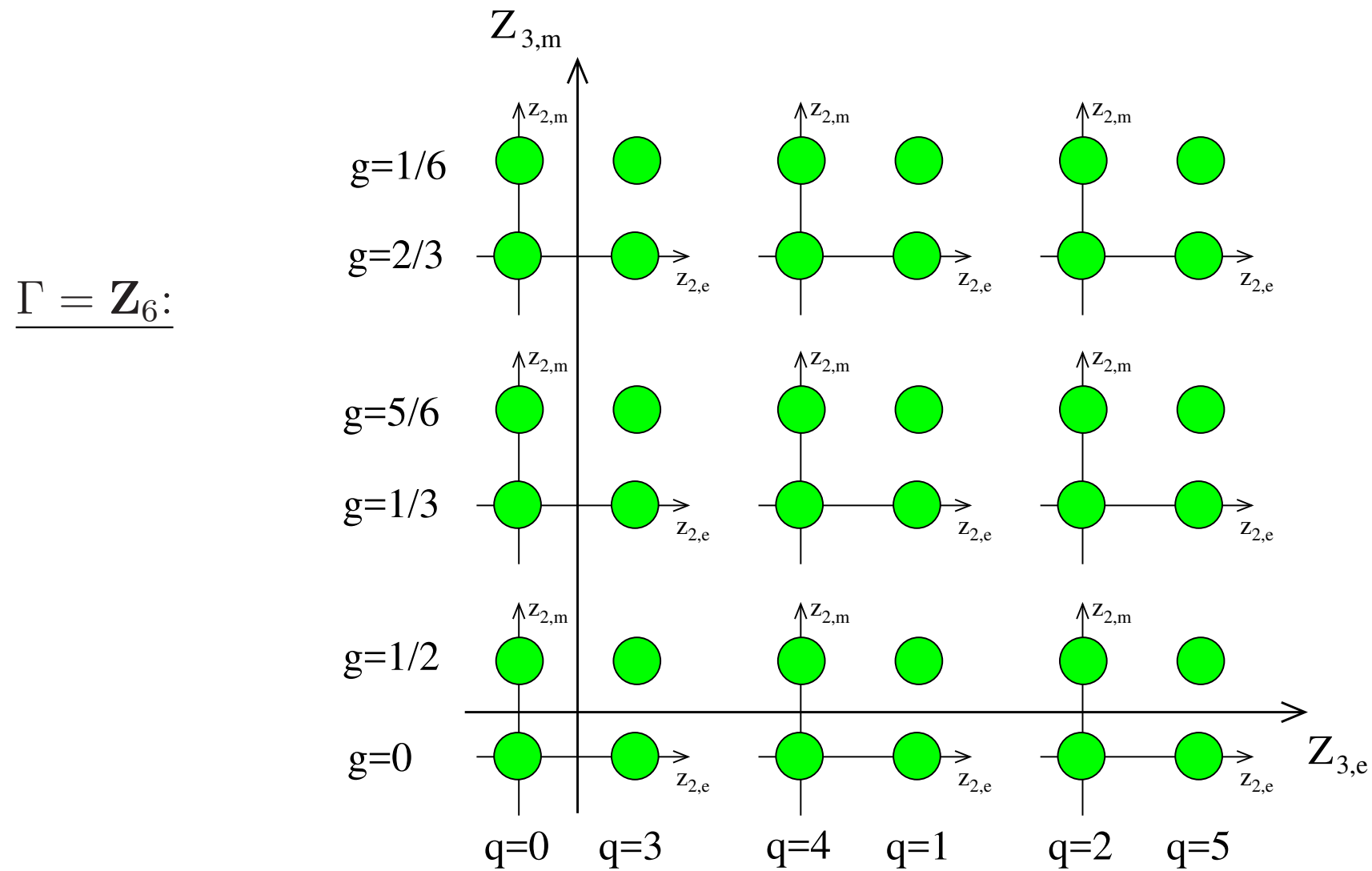
Line Operators

$$\Gamma = \mathbf{Z}_3:$$



Together with $U(1)$ line operators with $(q,g) = (3,0)$ and $(0,1)$

Line Operators



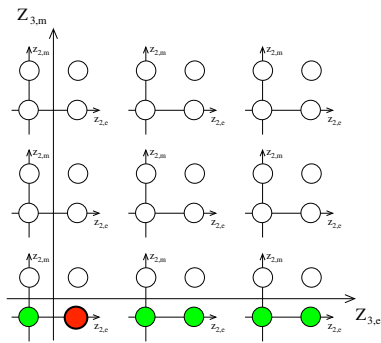
Together with $U(1)$ line operators with $(q, g) = (6, 0)$ and $(0, 1)$

Electroweak Symmetry Breaking

$$H : (\mathbf{2}, \mathbf{1})_3 \text{ condenses and } SU(2) \times U(1) \rightarrow U(1)_{em}$$

Denote electric and magnetic charges as Q and G (such that electron has $Q=-1$)

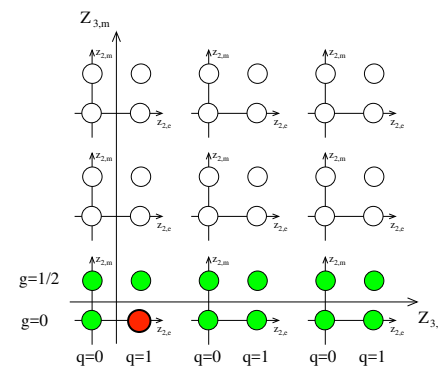
$$\Gamma = \mathbf{1}:$$



$$Q_{\min} = 1/6$$

$$G_{\min} = 6$$

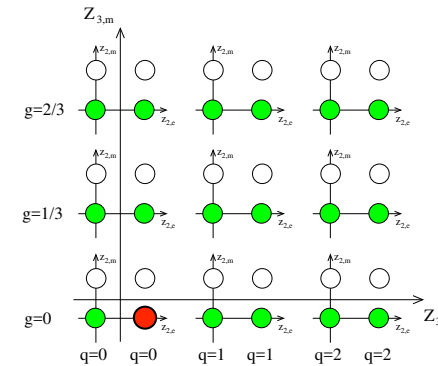
$$\Gamma = \mathbf{Z}_2:$$



$$Q_{\min} = 1/3$$

$$G_{\min} = 3$$

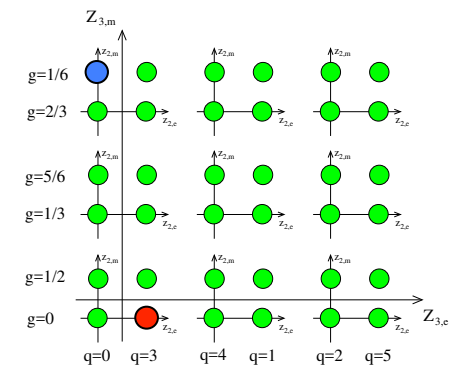
$$\Gamma = \mathbf{Z}_3:$$



$$Q_{\min} = 1/6$$

$$G_{\min} = 2$$

$$\Gamma = \mathbf{Z}_6:$$



$$Q_{\min} = 1/3$$

$$G_{\min} = 1$$

- Comments
- For \mathbf{Z}_3 and \mathbf{Z}_6 , the naïve Dirac quantisation, QG integer, does not hold
 - Magnetic monopole also carries colour magnetic charge
 - Black hole physics suggests all possible charges arise as dynamical objects

(Banks and Seiberg, 2010)

Theta Angles in the Standard Model

The periodicity of theta angles depends on the choice of $\Gamma = \mathbf{Z}_6, \mathbf{Z}_3, \mathbf{Z}_2$ or $\mathbf{1}$

Theta Angles

For $\Gamma = \mathbf{Z}_p$: $\theta_2, \theta_3 \in [0, 2\pi)$ and $\theta_Y \in [0, 2\pi p^2)$

Moreover, T-invariant states depend on the choice of Γ

$$\text{e.g: } \Gamma = \mathbf{1} \quad \Longrightarrow \quad \theta_2, \theta_3, \theta_Y = 0 \text{ or } \pi$$

$$\Gamma = \mathbf{Z}_6 \quad \Longrightarrow \quad \theta_2 = \theta_3 = 0 \text{ and } \theta_Y = 0 \text{ or } 36\pi$$

$$\theta_2 = 0 \text{ and } \theta_3 = \pi \text{ and } \theta_Y = 12\pi \text{ or } 48\pi$$

$$\theta_2 = \pi \text{ and } \theta_3 = 0 \text{ and } \theta_Y = 18\pi \text{ or } 54\pi$$

$$\theta_2 = \theta_3 = \pi \text{ and } \theta_Y = 30\pi \text{ or } 66\pi$$

Note: Anomalous $B+L$ means that we can rotate away one of the three theta angles.

$$\text{The surviving angles are } \theta_3 \text{ and } \theta_{em} = \frac{\theta_Y + 18\theta_2}{4}$$

The QCD Theta Angle

There is no observed CP violation from the strong force. Experiment bounds

$$\theta_3 \lesssim 10^{-10}$$

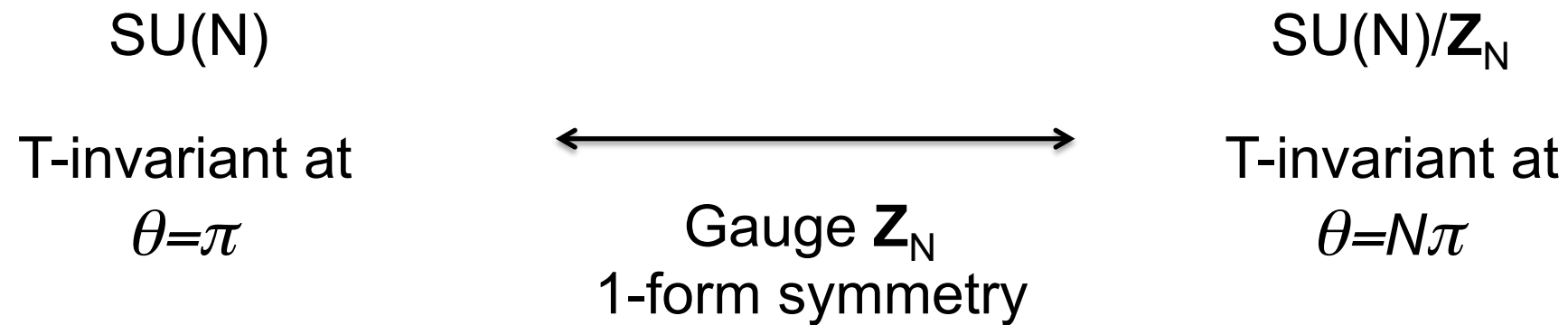
How do we know that θ_3 is not close to π ? The chiral Lagrangian tells us

$$\frac{m_{K^0}^2 - m_{K^+}^2 - m_{\pi^0}^2 + m_{\pi^+}^2}{m_\pi^2} \approx \frac{m_d \mp m_u}{m_d \pm m_u}$$

(Crewther, Di Vecchia, Veneziano and Witten 1979)

Is there a cleaner argument?

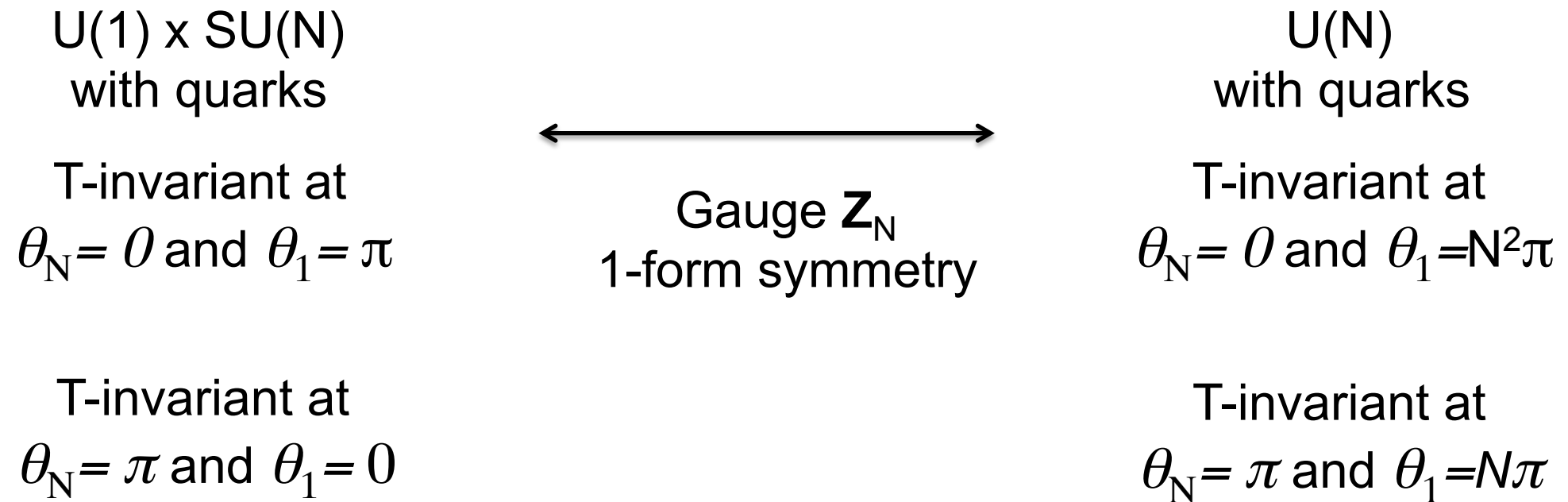
Yang-Mills at $\theta=\pi$



This is a mixed anomaly in time reversal and \mathbf{Z}_N center symmetry at $\theta=\pi$

't Hooft anomaly matching in infra-red \Rightarrow or • Time reversal is spontaneously broken
or • Theory is a TQFT
• Theory is not gapped

QCD at $\theta=\pi$



- There is again an anomaly when $\theta_N = \pi$
- There is also, now, a massless photon
 - This does not appear to be sufficient to absorb the anomaly

Further Questions

What are the physical implications of the different quotients?

How can we tell which gauge group describes our world?