# Long-Lived Inverse Chirp Signals from Core-Collapse in Massive Scalar-tensor Gravity Ulrich Sperhake

C Moore, M Agathos, R Rosca, D Gerosa, C Ott



DAMTP, University of Cambridge LIGO Scientific Collaboration LIGO-P1700218



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# Do we need a theory beyond GR?

When asked what he would do if Eddington's mission failed...



Then I would feel sorry for the good Lord. The theory is correct anyway.

(Albert Einstein)

izguotes.com

But we have reasons to search for "beyond GR"

Q Renormalization: Requires, e.g., higher curvature terms.
 → GR is low-energy limit of more fundamental theory
 Q Dark energy: Why is Λ so small and why ρ<sub>dark</sub> ~ ρ<sub>mat</sub>
 Q Dark matter: "Neptune" or "Vulcan" ?

# Scalar tensor theory of gravity

- Scalars appear naturally in extra-dimensional theories
- Scalars prominent in cosmology
- ST theory well-posed; fairly well understood mathematically
- No-hair theorems limit potential of black-hole spacetimes
  - ⇒ Matter: Neutron stars, core-collapse
- Best example of smoking gun to date:
  - Spontaneous scalarization Damour & Esposito-Farese PRL 1993
- Collapse studies in massless case

Novak PRD 1998/1999 Novak & Ibanez ApJ 2000, Gerosa+ CQG 2016

### Core-collapse scenario to 0th order

- Massive stars:  $M_{\rm ZAMS} = 8 \dots 100 \ M_{\odot}$
- Core compressed from  $\sim 1500 \text{ km}$  to  $\sim 15 \text{ km}$  $\sim 10^{10} \text{ g/cm}^3$  to  $\gtrsim 10^{15} \text{ g/cm}^3$
- Released gravitational energy:  $\mathcal{O}(10^{53})$  erg
   ~ 99 % in neutrinos, ~  $10^{51}$  erg in outgoing shock, explosion
- Explosion mechanism: still uncertainties...
- Failed explosions lead to BH formation
   "Collapsar": possible engine for long-soft GRBs
- All of this handled for us by Woosley & Heger Phys.Rept. 2007
   → Initial pre-collapse profile

# Theoretical framework

Einstein frame: conformal metric

$$\bar{g}_{\mu\nu} = F(\varphi) \, g_{\mu\nu}$$

Action

$$S = \frac{1}{16\pi} \int dx^4 \sqrt{-\bar{g}} \left[ \bar{R} - 2\bar{g}^{\mu\nu} \partial_\mu \varphi \,\partial_\nu \varphi - 4V(\varphi) \right] + S_m [\psi_m, \bar{g}_{\mu\nu}/F(\varphi)]$$

- Energy momentum tensor:  $T_{\alpha\beta} = \rho h u_{\alpha} u_{\beta} + P g_{\alpha\beta}$
- Spherical symmetry:  $d\bar{s}^2 = \bar{g}_{\mu\nu}dx^{\mu}dx^{\nu} = -F\alpha^2 dt^2 + FX^2 dr^2 + r^2 d\Omega^2$

$$u^{\alpha} = \frac{1}{\sqrt{1 - v^2}} [\alpha^{-1}, \ vX^{-1}, \ 0, \ 0]$$

- Equations (gravity):  $\partial_r \alpha = \dots$ ,  $\partial_r X = \dots$  $\partial_t \partial_t \varphi = \dots$
- Equations (matter):  $(\rho, h, v) \leftrightarrow (D, S^r, \tau) \Rightarrow HRSC$ GR1D code O'Connor & Ott CQG 2009

# Equation of state

- Pressure: "cold" + "thermal" contribution: P = P<sub>c</sub> + P<sub>th</sub>
   Hybrid EOS for cold part: P<sub>c</sub> =  $\begin{cases}
   K_1 \rho^{\Gamma_1} & \text{if } \rho \leq \rho_{\text{nuc}} \\
   K_2 \rho^{\Gamma_2} & \text{if } \rho > \rho_{\text{nuc}}
   \end{cases}$
- Internal energy from 1st law:  $\epsilon_c = \begin{cases} \frac{K_1}{\Gamma_1 1} \rho^{\Gamma_1 1} & \text{if } \rho \leq \rho_{\text{nuc}} \\ \frac{K_2}{\Gamma_2 1} \rho^{\Gamma_2 1} + E_3 & \text{if } \rho > \rho_{\text{nuc}} \end{cases}$
- Thermal pressure:  $P_{\rm th} = (\Gamma_{\rm th} 1)\rho(\epsilon \epsilon_{\rm th})$
- Parameters:  $\Gamma_1 = 1.3$ ,  $\Gamma_2 = 2.5$ ,  $\Gamma_{th} = 1.35$

 $K_1 = 4.9345 \times 10^{14} \text{ [cgs]}, \quad \rho_{\text{nuc}} = 2 \times 10^{14} \text{ g cm}^{-3}$  $K_2, \quad E_3 \quad \text{from continuity at} \quad \rho = \rho_{\text{nuc}}$ 

# The coupling function and potential

Coupling function, potential:

$$F(\varphi) = e^{-2\alpha_0\varphi - \beta_0\varphi^2}$$

$$V(\varphi) = \frac{1}{2}\mu^2\varphi^2$$



• Only for  $\mu \lesssim 10^{-19} \,\mathrm{eV}$  !! Here:  $\mu[\mathrm{eV}] \in [10^{-15}, 10^{-12}]$ Ramazanoglu & Pretorius PRD 2016

Free parameters:  $\mu, \alpha_0, \beta_0, \Gamma_1, \Gamma_2, \Gamma_{th}$ 

### Convergence test

• For  $\mu = 10^{-14} \text{ eV}$ ,  $\alpha_0 = 10^{-4}$ ,  $\beta_0 = -20$  $\Gamma_1 = 1.3$ ,  $\Gamma_2 = 2.5$ ,  $\Gamma_{\text{th}} = 1.35$ 

• Using  $N_1 = 5000$ ,  $N_2 = 10000$ ,  $N_3 = 20000$  points



• Discretization error:  $\sim 5\%$ 



 $r\varphi \gg$  massless case; fairly insensitive to parameters; dispersion!

# Waveforms ``far from" the source

- LIGO will observe the
   above scalar profiles
   after they propagate
   to large distances
- In the massless case this is almost trivial  $\varphi(t;r) = \frac{1}{r}\varphi(t-r;r_{extract})$
- In the massive case
   things are more
   complicated: signals
   propagate with
   dispersion



# Waveforms ``far from" the source Far from the source, scalar dynamics are governed by the flat-space Klein-Gordon wave equation ∂<sup>2</sup><sub>t</sub>φ - ∇<sup>2</sup>φ + ω<sup>2</sup><sub>\*</sub>φ = 0 Easier to work with the radially rescaled field σ ≡ rφ As the signal propagates outwards: Low frequencies are suppressed High frequency power spectrum is unaffected Signal spreads out in time High frequencies arrive earlier than low frequencies

- Signal becomes increasingly oscillatory



# Waveforms ``far from" the source

Signals become more oscillatory as they propagate outwards

- In the large-distance limit the stationary phase approximation applies  $\rightarrow$  analytic expression for the time domain signal
  - Signals have a characteristic "inverse chirp" lasting many years

SPA frequency as function of time (Inverse Chirp)  $F(t) = \frac{\omega_*}{2\pi} \frac{1}{\sqrt{1 - (d/t)^2}}$ 

Distance to source  $d = 10 \,\mathrm{kpc}$ 



# Detection with LIGO-Virgo

GWs from core-collapse in ST gravity may fall into 3 classes:

- Burst signals: For light scalars (μ < 10<sup>-20</sup> eV) and short
   distances (10 kpc), the pulse does not disperse significantly;
   will look like a < 1 s burst</li>
- Continuous wave signal: for heavier scalars, long dispersion turns pulse into a quasi-monochromatic signal
   capture using standard directed CW searches, assuming EM counterpart; e.g. SN1987A, Kepler1604
- Stochastic background:
  - Many quiet sources + very long duration (superposed)
  - Cosmological redshift + mass variation  $\rightarrow$  smeared low-f cutoff
  - Characteristic "bump" in background, peaking at  $\sim \omega_*$
  - Well in reach for aLIGO/AdVirgo stochastic searches

### Conclusions

- We have simulated stellar core collapse in massive ST theory
- Explored combined parameter space of EOS and ST theory parameters
- Spontaneous scalarization occurs as in massless case, but
   effect can be more dramatic because the scalar mass "screens"
   the effect of the scalar, allowing larger values of α<sub>0</sub>, β<sub>0</sub> to be
   compatible with binary pulsar observations
- Signals propagate with dispersion, signals can last for years to centuries at kpc distances
- Signals can show up in LIGO/Virgo burst, CW or stochastic searches