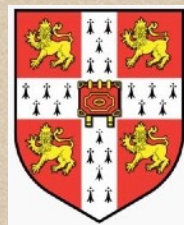


Numerical and analytic modeling of boson stars

Ulrich Sperhake



DAMTP, University of Cambridge

UKNR 2025
University of Birmingham
15-16 Dec 2025



1. Background

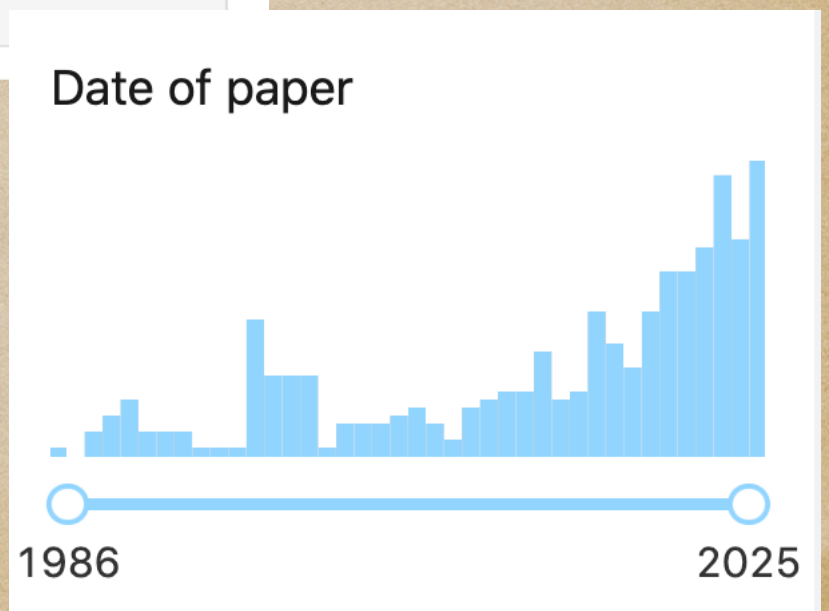
The popularity of boson stars



The popularity of boson stars



find t boson stars



The idea of boson stars

- "Gravitational-electromagnetic entities" or Geons

- Wheeler 1955

PHYSICAL REVIEW

VOLUME 97, NUMBER 2

JANUARY 15, 1955

Geons*

JOHN ARCHIBALD WHEELER

Palmer Physical Laboratory, Princeton University, Princeton, New Jersey

(Received September 8, 1954)

Associated with an electromagnetic disturbance is a mass, the gravitational attraction of which under appropriate circumstances is capable of holding the disturbance together for a time long in comparison with the characteristic periods of the system. Such

of equations of self-consistent geon; mass and radius values. 4. Transformations and interactions of electromagnetic geons; evaluation of refractive index barrier penetration integral for spherical geon; photon-photon collision processes as additional

- Energy = mass gravitates → Compact (equilibrium?) objects
- Geons are not equilibrium configurations
- Dark matter candidates: QCD axions, ALPs, dark photons,...
- Complex fields (scalar, vector,...)
→ Genuine equilibrium states; $T_{\alpha\beta}$ stationary!
- First shown for scalar fields → "Boson stars"
- Feinblum & McKinley PR 168 (1968), Kaup PR 172 (1968),
Ruffini & Bonazzola PR 187 (1969)

A boson star zoo

- Mini BSs (no self-interaction) Kaup PR (1968) and others
- "Solitonic" BSs (self-interacting scalar field) → more compact
Colpi+ PRL (1986), Lee PRD (1987), ...
- Proca stars Brito+ Phys.Lett.B (2016)
- ℓ -boson stars (multiple scalar fields) Alcubierre+ CQG (2018)
- Multi-oscillating BSs Choptuik+ PRL (2019)
- Thin-shell BSs Collodel & Doneva 2203.08203
- Axion stars e.g. Helfer et al JCAP 2017
- Higher-spin fields Jain & Amin 2109.04892
- Multi-field BSs Sanchis-Gual+ PRL (2021)
- May condense from local over-densities Widdicombe+ JCAP (2018)

Focus here: Single-scalar BSs

A boson-star précis

- BS = Classical limit of a quantum field
Ground state = Boson-Einstein condensate
- Only known scalar field: Higgs, $\mu = 125 \text{ GeV}$
- Units $G = c = \hbar = 1$, scale determined by μ [inverse length]
E.g. $\mu = 1.33 \times 10^{-10} \text{ eV} \Rightarrow \mu^{-1} = M_{\odot} = 1.48 \text{ km}$
- Astrophysical stars need ultralight fields $\rightarrow$ Beyond SM
axion like, dark matter candidates: $10^{-11} \dots 10^{-20} \text{ eV}$
- Properties: Compactness 0 to $> \text{NSs}$, any Mass
- BSs can have light rings $\rightarrow$ Black-hole mimickers; stable???

Reviews: Liebling & Palenzuela LRR 2023, Visinelli 2109.05481,

Mielke Fund.Theor.Phys. 2016, Schunck & Mielke GRG 1999, ...

Formalism and basic features

- GR + minimally coupled complex scalar field φ

$$S = \int \sqrt{-g} \left\{ \frac{1}{16\pi G} R - \frac{1}{2} [g^{\mu\nu} \nabla_\mu \bar{\varphi} \nabla_\nu \varphi + V(\varphi)] \right\} dx^4$$

$$T_{\alpha\beta} = \partial_{(\alpha} \bar{\varphi} \partial_{\beta)} \varphi - \frac{1}{2} g_{\alpha\beta} [g^{\mu\nu} \partial_\mu \bar{\varphi} \partial_\nu \varphi + V(\varphi)]$$

- Potential; analogous to EOS:

$$V_{\text{mas}}(\varphi) = \mu^2 |\varphi|^2 + \lambda |\varphi|^4, \quad V_{\text{sol}}(\varphi) = \mu^2 |\varphi|^2 \left(1 - 2 \frac{|\varphi|^2}{\sigma_0^2} \right)^2$$

- Mini Boson Stars: $\lambda \rightarrow 0, \quad \sigma_0 \rightarrow \infty$

- Ansatz $\varphi(t, r) = A(r) e^{i\epsilon\omega t + \delta\Phi}, \quad \varphi(t, r, \theta, \phi) = A(r, \theta) e^{i(\epsilon\omega t + m\phi + \delta\Phi)}$

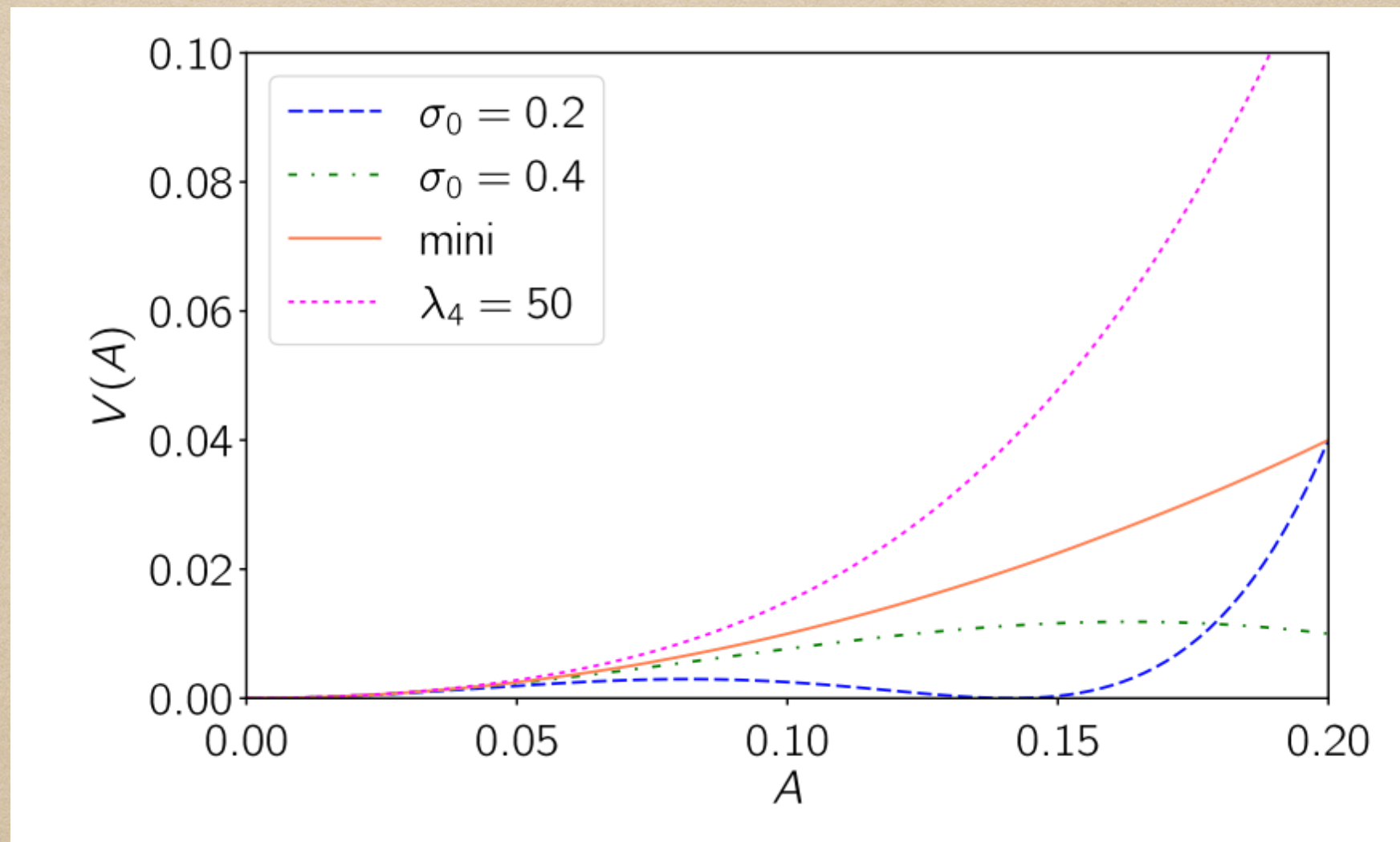
$\delta\Phi$ = phase offset, ϵ = BS vs anti-BS

- Regular solutions only for countably infinite values

$$\omega_0 < \omega_1 < \omega_2 < \dots \quad (\text{ground state, excited states})$$

Formalism and basic features

- Scalar potentials



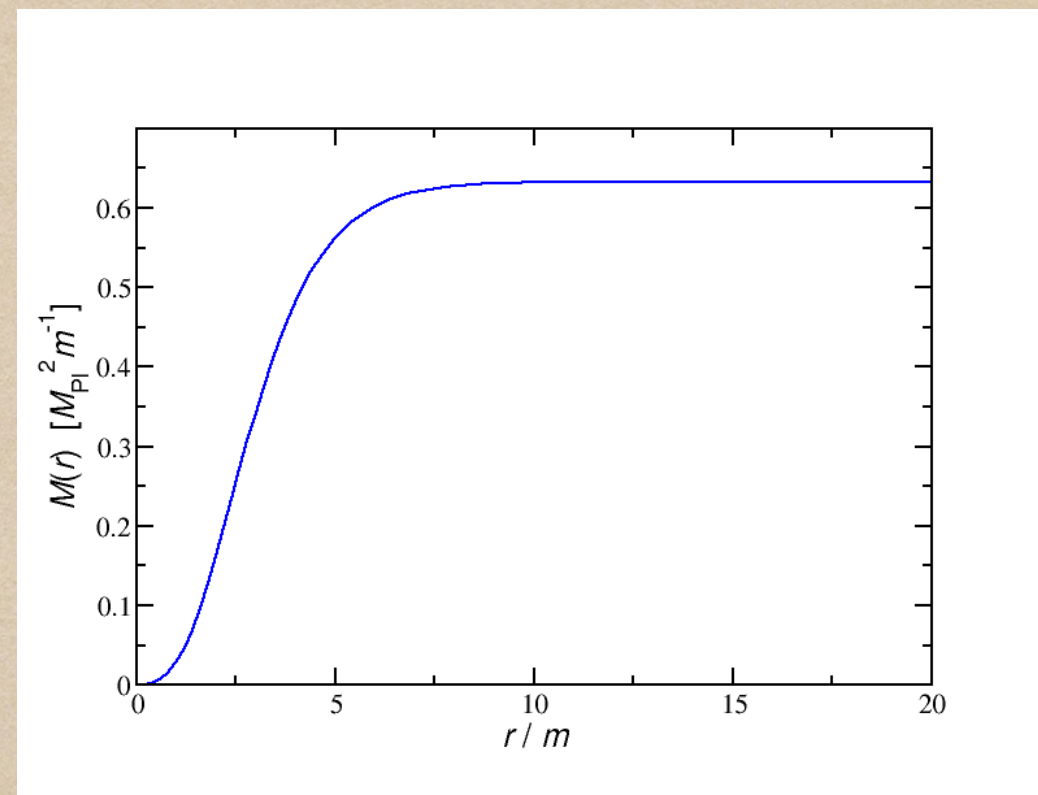
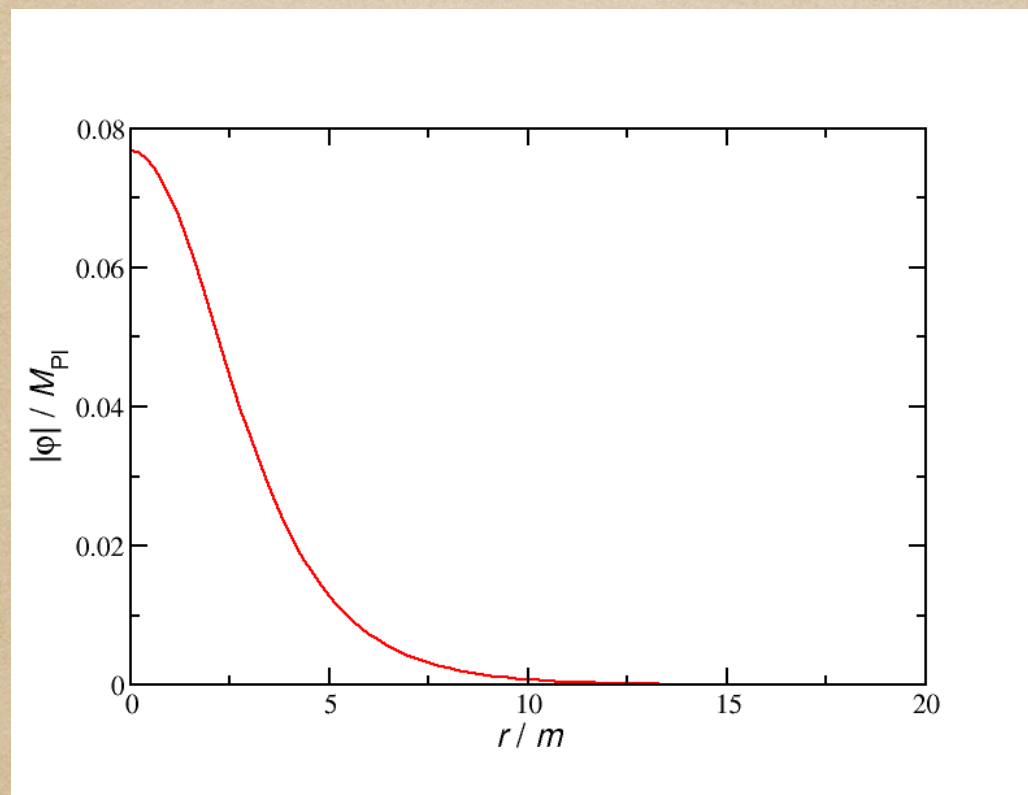
Evstafyeva et al PRD 2023

- Massive Potential: Stiff EOS
- Solitonic Potential: Soft EOS

Formalism and basic features

- E.g. Maximal-mass, spherically symm. mini boson star (Kaup limit)

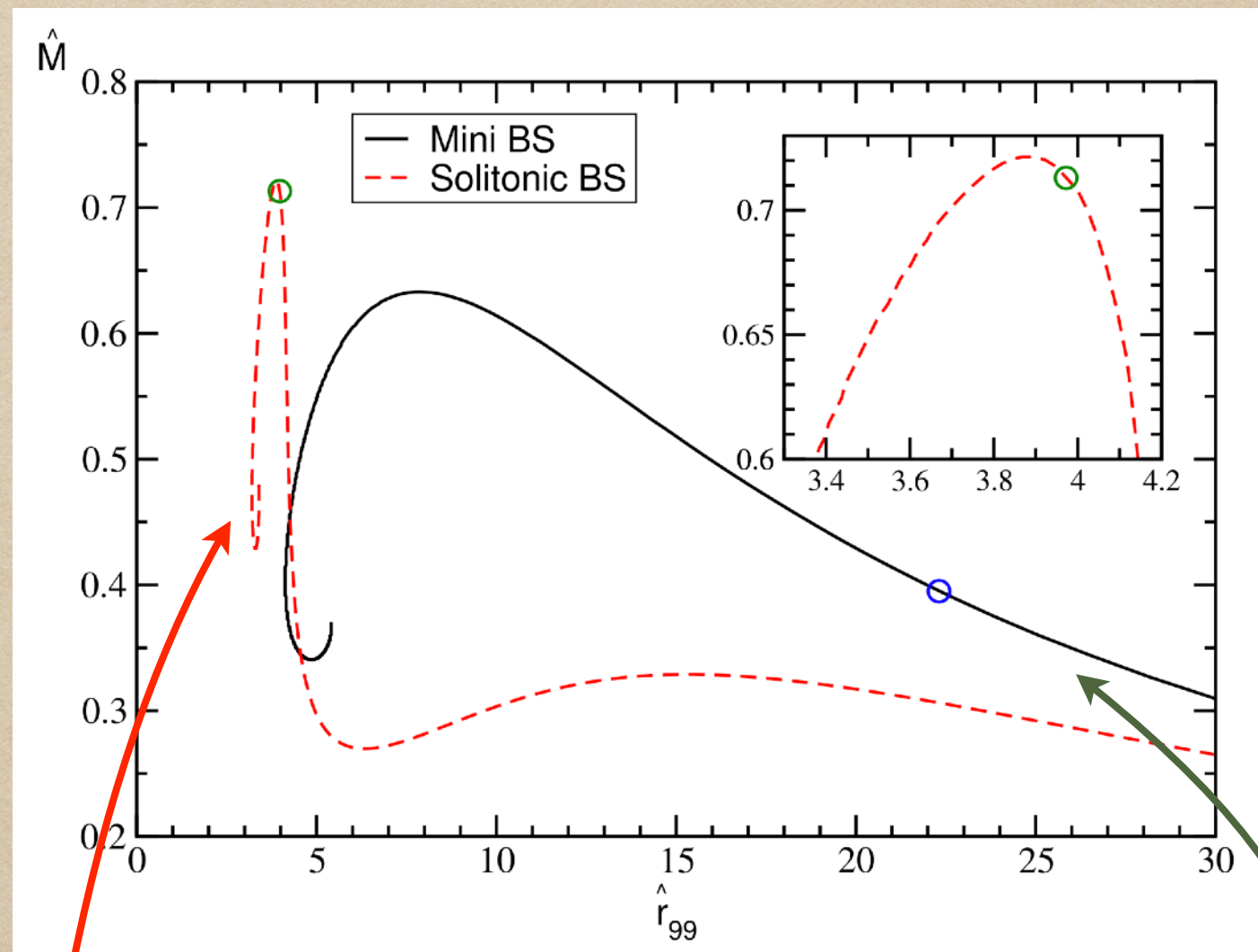
$$\omega_0 = 0.853 m, \quad M = 0.633 M_{\text{Pl}}^2 / m$$



- Excited states unstable
collapse to BH, dispersion or migration to stable ground-state BS
Balakrishna, Seidel, Suen PRD (1998)

Formalism and basic features

- Mass-Radius curves similar to Tolman-Oppenheimer-Volkoff stars



unstable

stable

- Radial stability similar to TOV stars Marks 2508.11757

Spinning Boson Stars

- Scalar BSs cannot spin perturbatively Kobayashi+ PRD (1994)
- Spinning scalar BSs exist with but have quantized spin
Schunck & Mielke Phys.Lett.A (1998)
- Spinning scalar BSs may be unstable in contrast to
spinning Proca stars! Sanchis-Gual+ PRL (2019)

Possibly due to toroidal structure: scalar field vanishes at origin

- What happens in scalar BS inspiral and merger?
 - Kerr BH
 - Non-spinning BS; angular momentum shed
 - Total dispersal
 - Spinning BS with exact angular momentum?

2. A primer of boson-star dynamics



Theory and Numerical Modelling

- Massive complex scalar field + GR

$$S = \int \frac{\sqrt{-2}}{2} \left\{ \frac{R}{8\pi G} - [g^{\mu\nu} \nabla_\mu \bar{\varphi} \nabla_\nu \varphi + V(\varphi)] \right\} d^4x$$

⇒ Einstein-Klein-Gordon equations

- Space-time formulations: GHG, CCZ4, BSSNOK

- Techniques like for BH binaries. We use:

GRChombo Radia et al CQG 2021

Lean US PRD 2006

- Typical grid setups:

$dx = \frac{M}{48} \cdots \frac{M}{32}$, domain size $\sim 1024 M$, 8 refinement levels

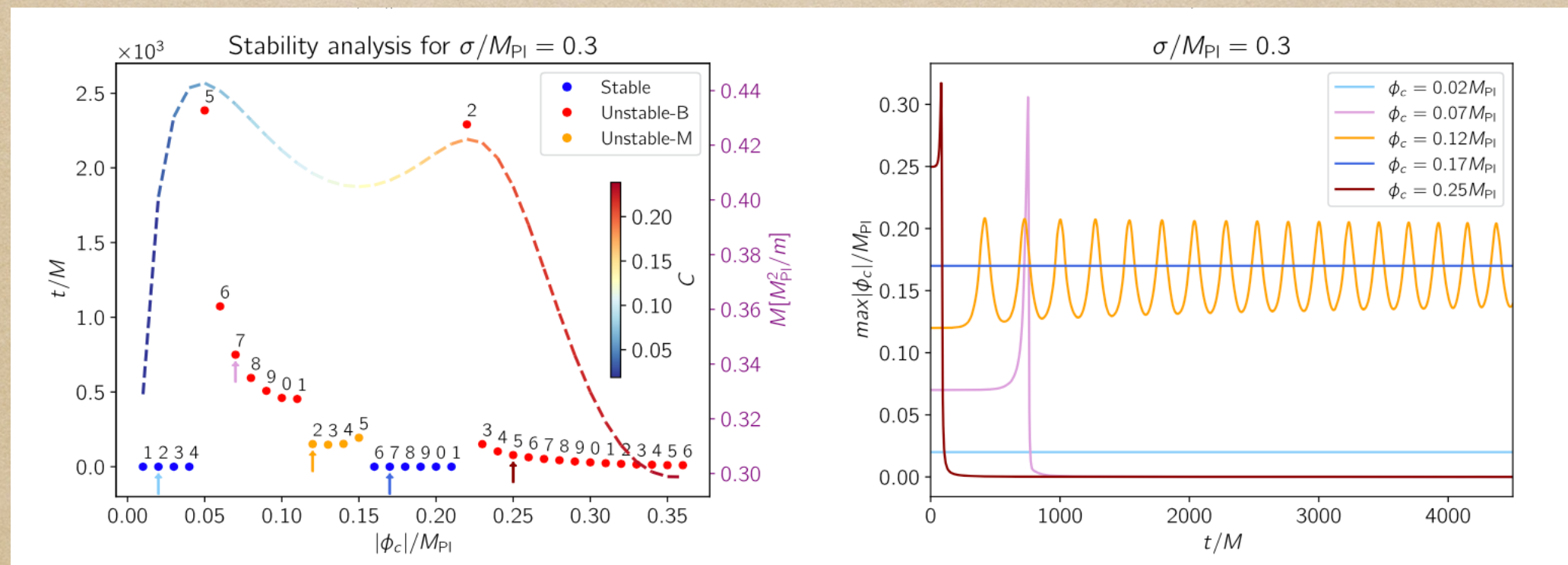
- Improve binary initial data: Constant Volume Element (CVE)

Helper et al PRD 2019, CQG 2021, Evstafyeva et al CQG 2023,

Atteneder et al PRD 2024

Stability of nonspinning BSs

- Select potential; here solitonic
- Maxima in $M(|\phi_c|)$ curve (typically) signals swap in stability
- Radial stability similar to TOV stars Marks 2508.11757

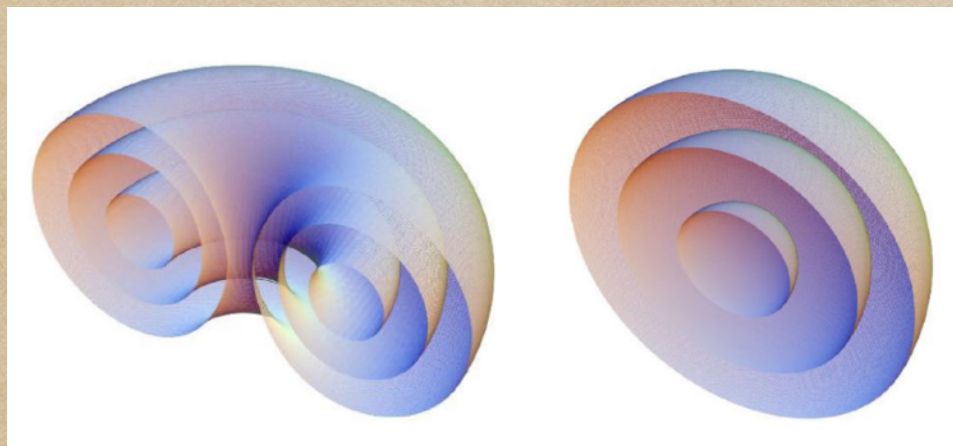


Ge et al. 2410.23839

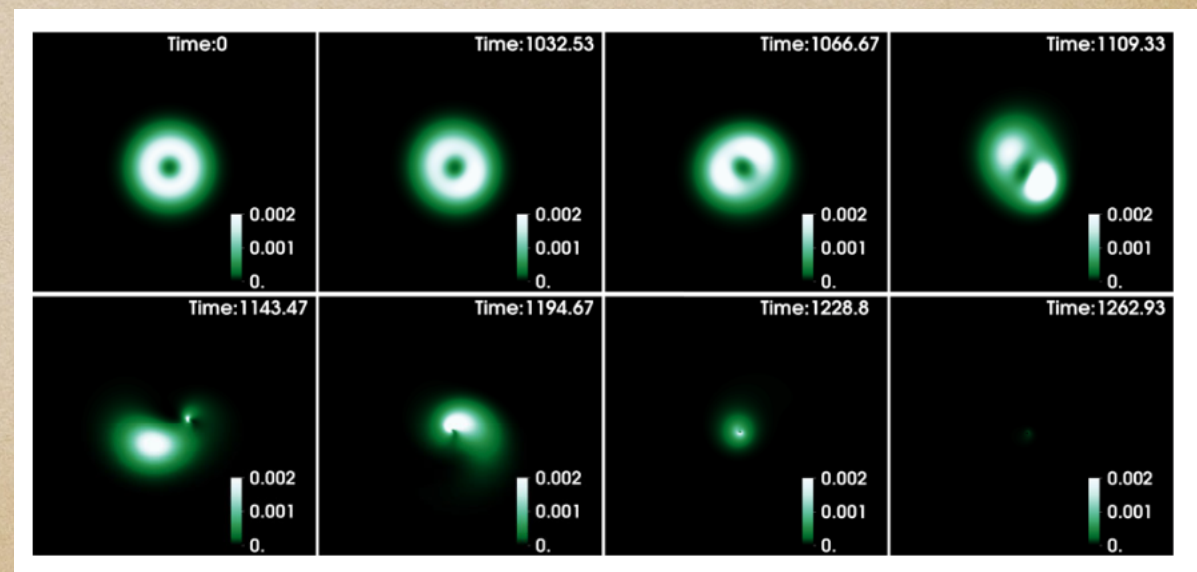
- Evolution of BSs:
 - BH collapse
 - Migration to stable BS
 - Evaporation to infinity (partial or whole)

Formation and Stability of spinning BSs

- Evolve scalar or vector cloud; no self-interaction!
 - Cloud collapses due to gravitational cooling
 - Scalar case: Instability shedding all angular momentum
 - Vector case: No instability, spinning Proca star forming
- Evolve spinning equilibrium stars
 - Scalar case: Instability → Formation of spinning BH
 - Vector case: Stable evolutions even for explicit perturbations
- May be due toroidal/spherical energy profile

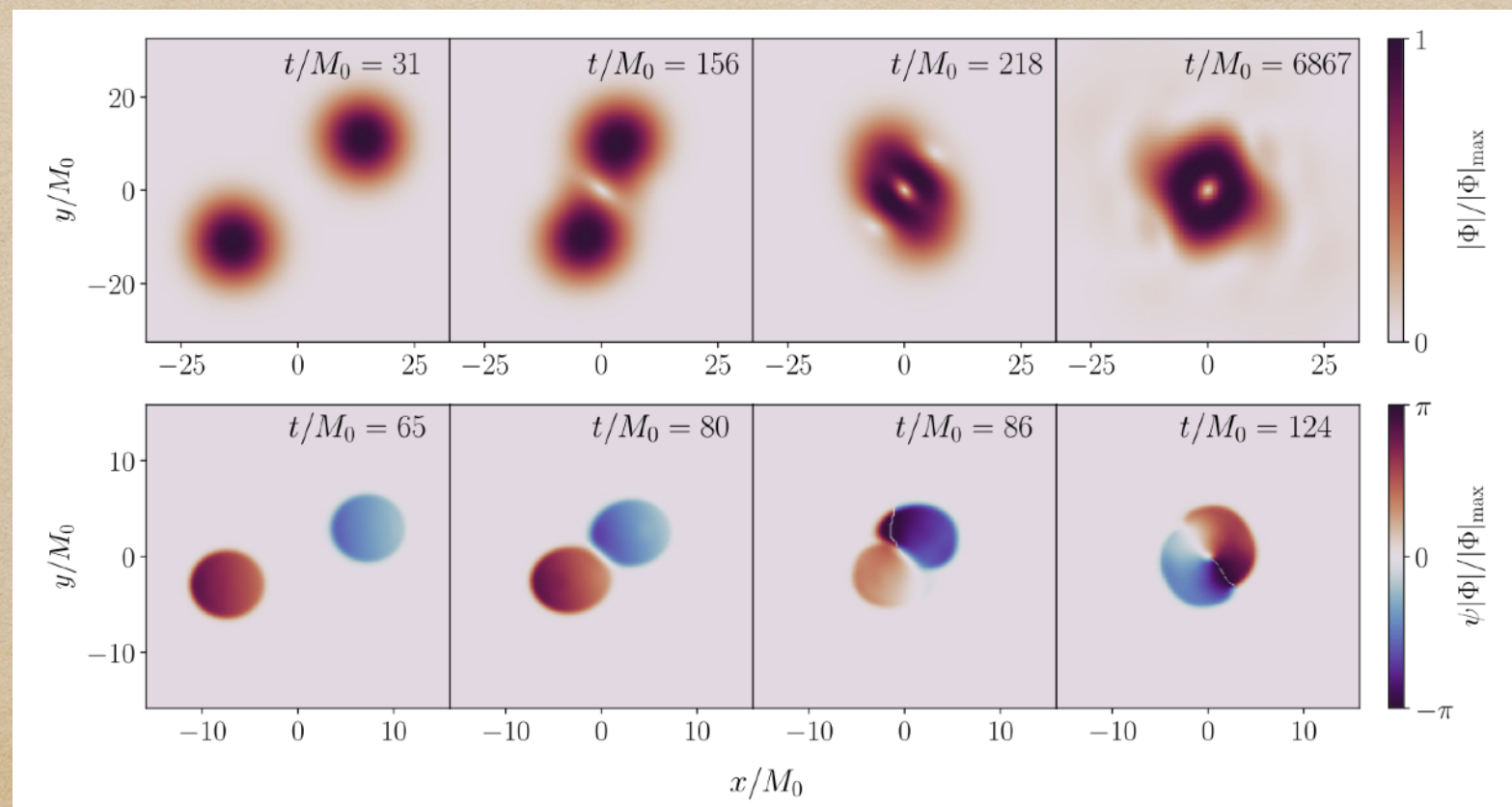


Sanchis-Gual et al PRL 2019



Spinning scalar BSs with self-interaction

- Nonrotating ($\delta\Phi = 0 \dots \pi$, $M/r \approx 0.037$) scalar BS binary on eccentric orbits Siemonsen & East PRD 2023
- Massive and solitonic potentials
- Potential and angular momentum strongly affects GW signal
- Threshold dephasing $0 < \delta\Phi_{\text{cr}} < \pi$ required for forming spinning BS
- Dephasing, scalar interaction, vortex condition (J/M^2) important



Black-hole mimickers

- Ultracompact objects with light ring but no horizon
- Observational evidence for BHs
 - Accretion
 - GWs + QNM ringdown

Talks Seppe Staelens
Gereth Marks!

Related closely to **light rings** rather than event horizon

- Theorem: Smooth compact objects with light ring have at least **two** light rings, one them **stable**. Cunha, Berti & Herdeiro PRL 2017
- Stable LRs = attractor for (massless) particles → Instability?
Keir CQG 2016, Cardoso et al PRD 2014, Cunha 2503.00117
- NR simulations of ultra compact boson stars inconclusive:
 - Migration or BH collapse: Cunha et al PRL 2023
 - Long-term stable BS evolutions: Marks et al PRL 2025,
Evstafyeva et al 2508.11527, Siemonsen PRL 2024, Staelens 2508.19460

Ergoregion Instability

- Ergoregion = All dragged in rotation sense
- Horizonless compact objects with ergoregion
 - Ergoregion Instability: Exponential growth of negative energy modes through energy emission to ∞

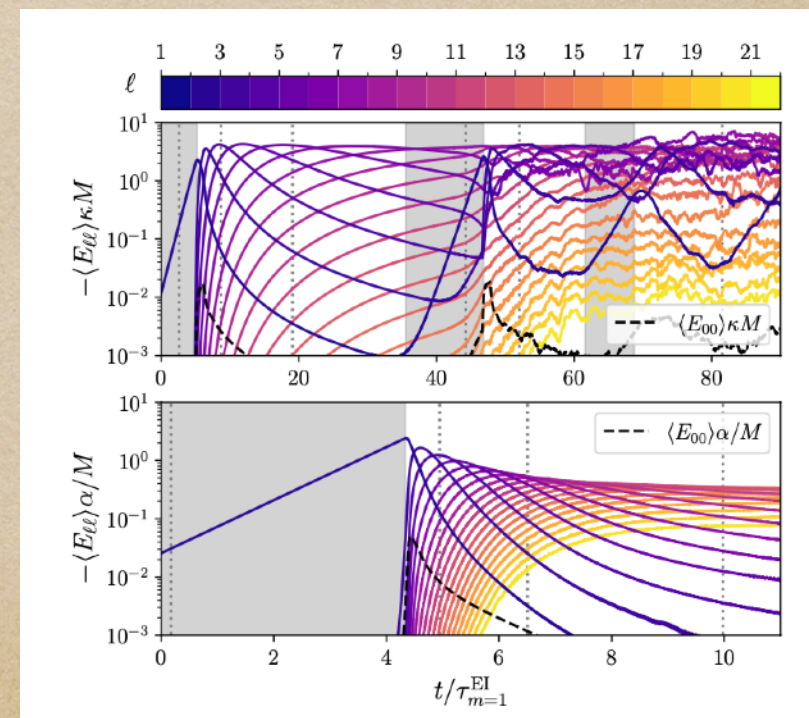
Friedmann C.Ma.Ph. 1978, Vilenkin Ph.Lett.B 1978, Moschidis C.Ma.Ph. 2018

- Cf. R-Mode Instability of Neutron stars, but unlike Superradiance
- NR simulations of rapidly spinning BSs and toy problem:

- e -folding time can be $10^4 M$
- $l = m = 1$ retrograde mode grows
- Then turbulent cascade
- (short) NR runs: No instability

Siemonsen 2510.07467, 2510.07468

Siemonsen & East 2512.10526

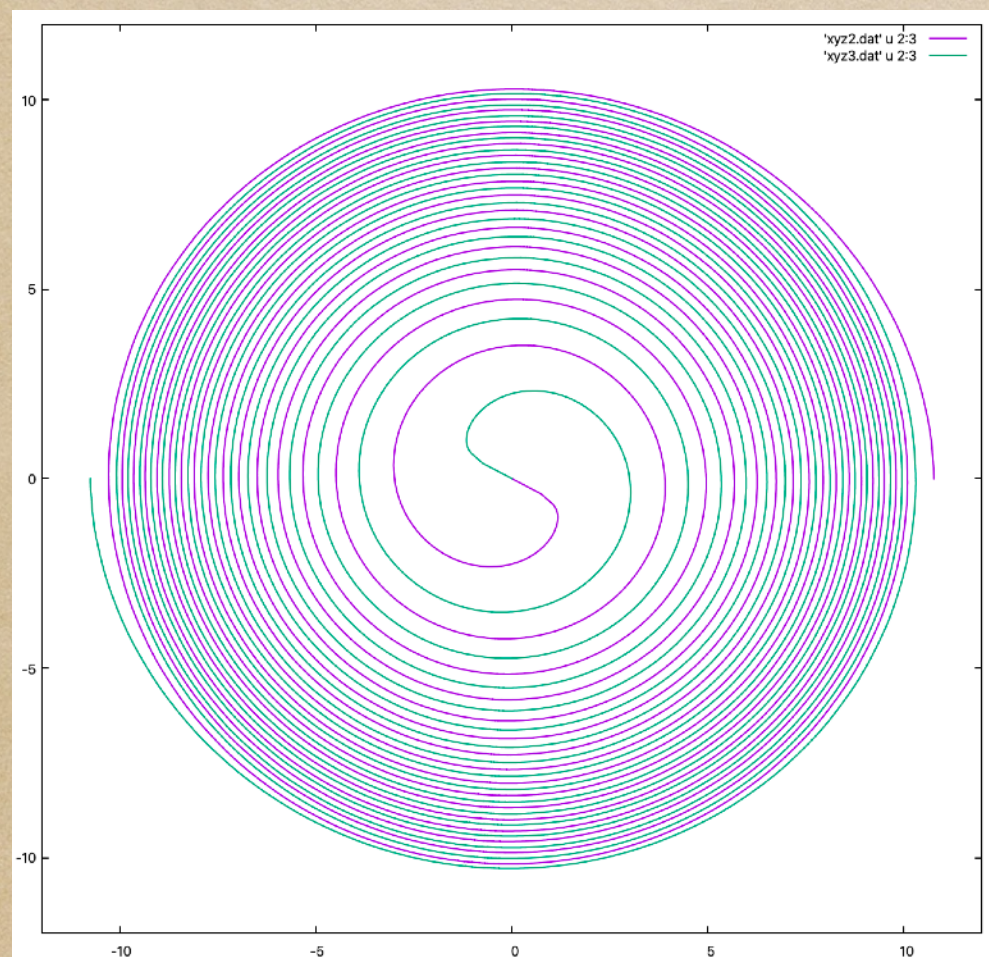


BS binaries

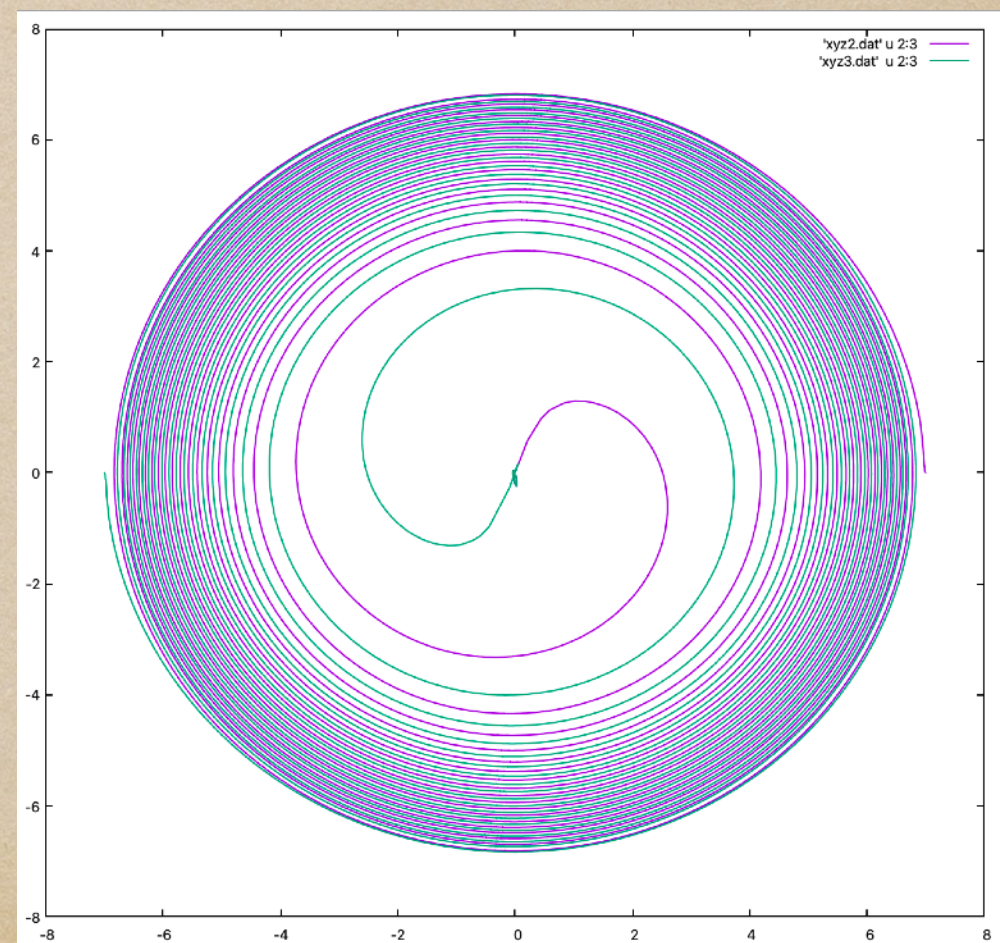
- Phase error $\approx 0.1 \dots 0.2$
- Amplitude error $\lesssim 3\%$
- Eccentricity $\approx 0.002 \dots 0.005$

Evstafyeva et al 2406.02715

compact BS binary: A17-d14

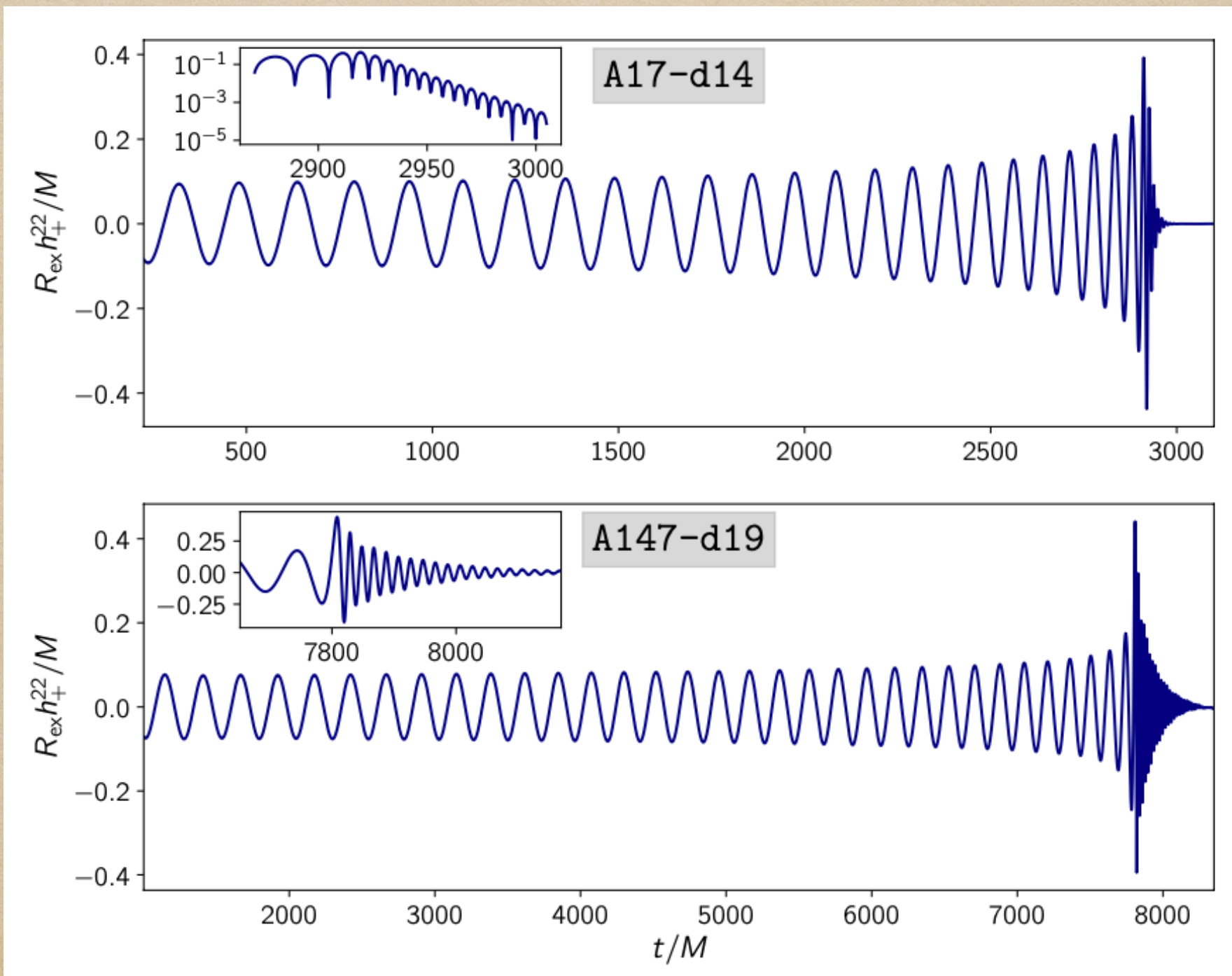


fluffy BS binary A147-d19



BS binaries

- GW strain from compact and fluffy BSs

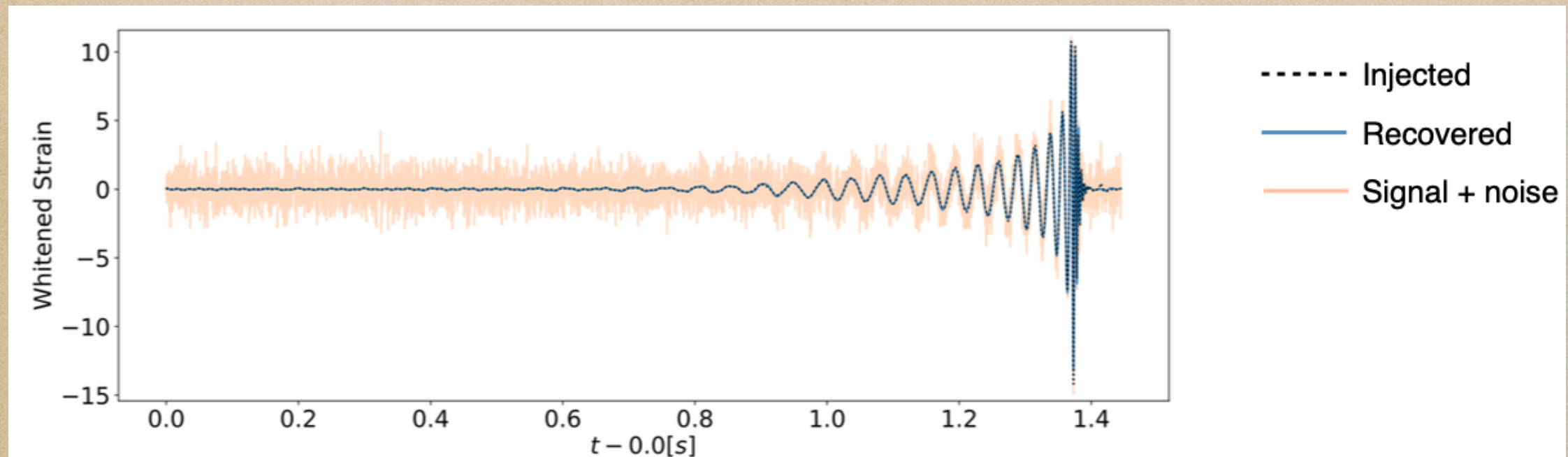


Waveform approximants

- Parameter estimation performed with Bilby Ashton et al 2019
- IMRPhenomXP: Frequency domain Pratten et al 2021
 - Quasi-circular, spin-precessing black-hole binaries
 - Quadrupole modes
- IMRPhenomPv2_NRTidal: Frequency domain Dietrich et al 2017, 2019
 - Quasi-circular, spin-precessing neutron-star binaries
 - Tidal deformability parameters $\Lambda_{A,B}$
- We have tested more with similar results.

Compact BSs using IMRPhenomXP

- Injections often recovered but with **biased** parameters!
- Example A17-d15 with $M_{\text{tot}} = 77 M_{\odot}$, $d_L = 200 \text{ Mpc}$ in the analysis



- **Recovered:** $M_1 = 37.8 \pm 1.1 M_{\odot}$, $M_2 = 25.4 \pm 1.2 M_{\odot}$, $d_L = 236 \pm 20 \text{ Mpc}$,
 $a_1 \approx 0.95$, $a_2 \approx 0.15$

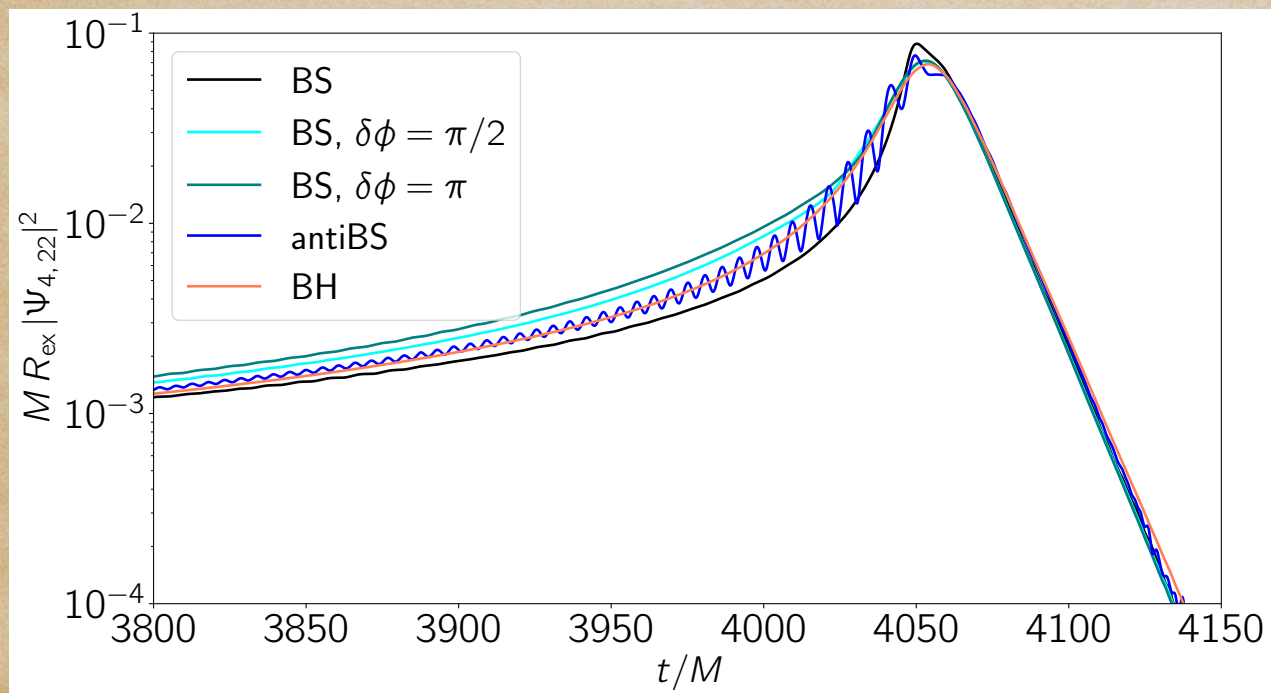
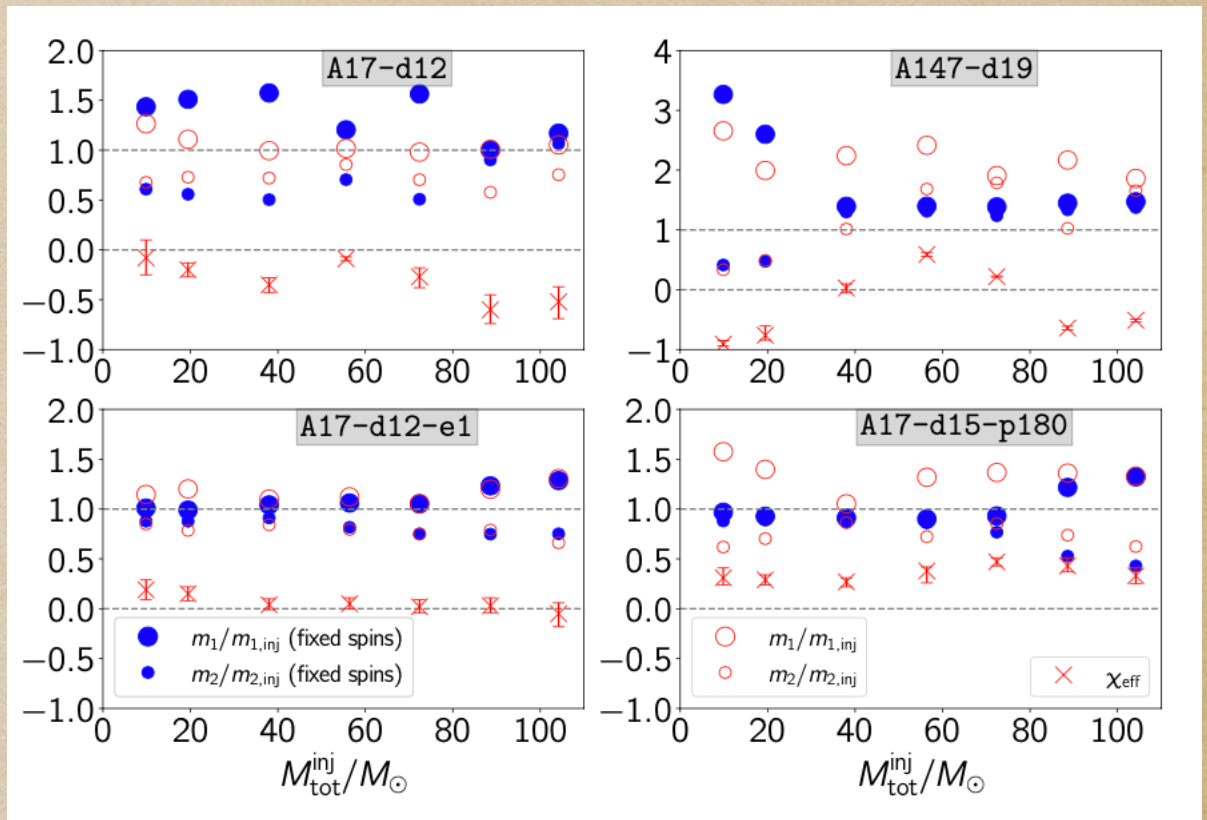
Recovered SNR $\approx$ injected SNR

$$\log \mathcal{B}_N^S = 5392$$

- Parameter bias **not** random!

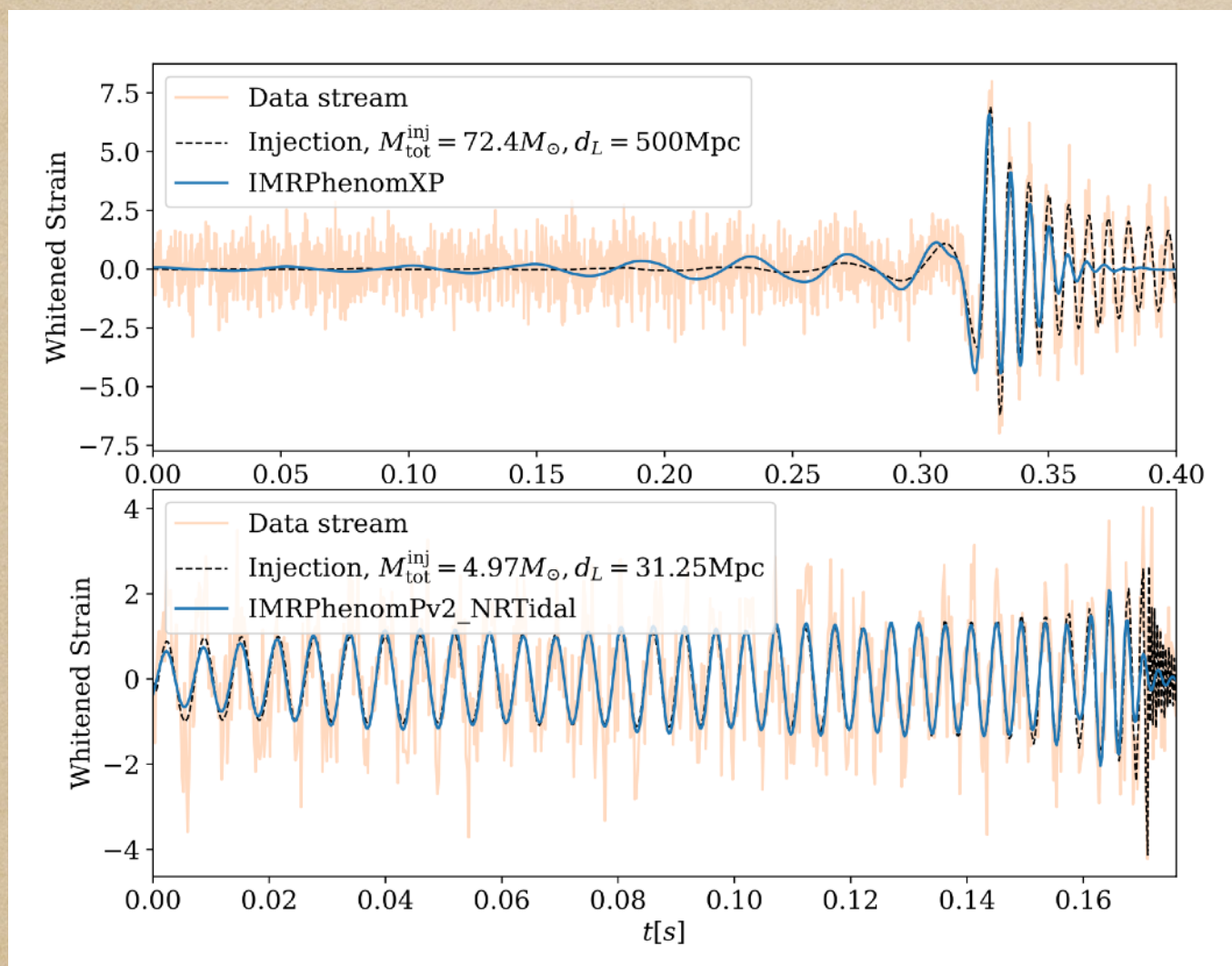
Understanding the PE bias

- PE bias
- VS
- Chirp steepness



Recovery of “Fluffy” BS binaries

- Parameter estimation always erratic for *fluffy* BBS
- BH approximants may capture inspiral or merger but never both!
- Residual often not compatible with Gaussian noise

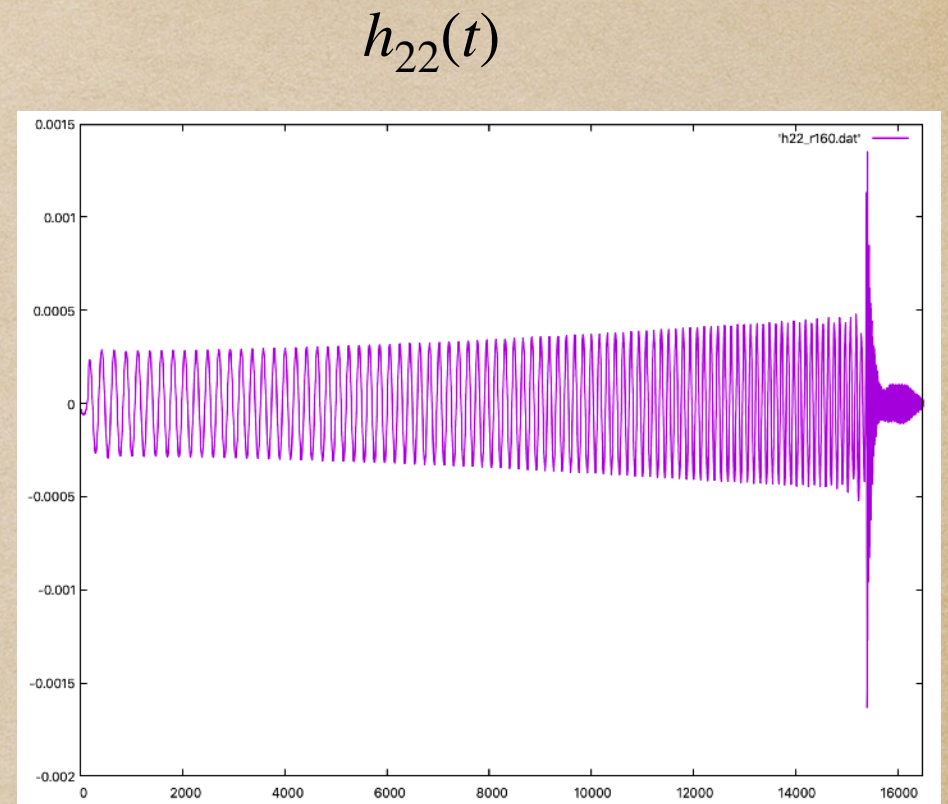
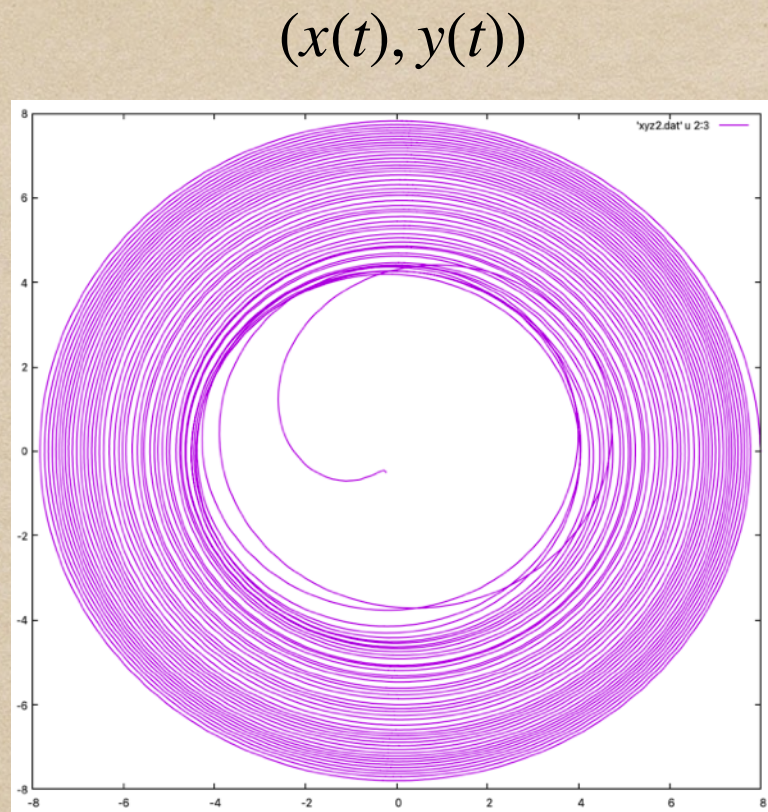


cf. also

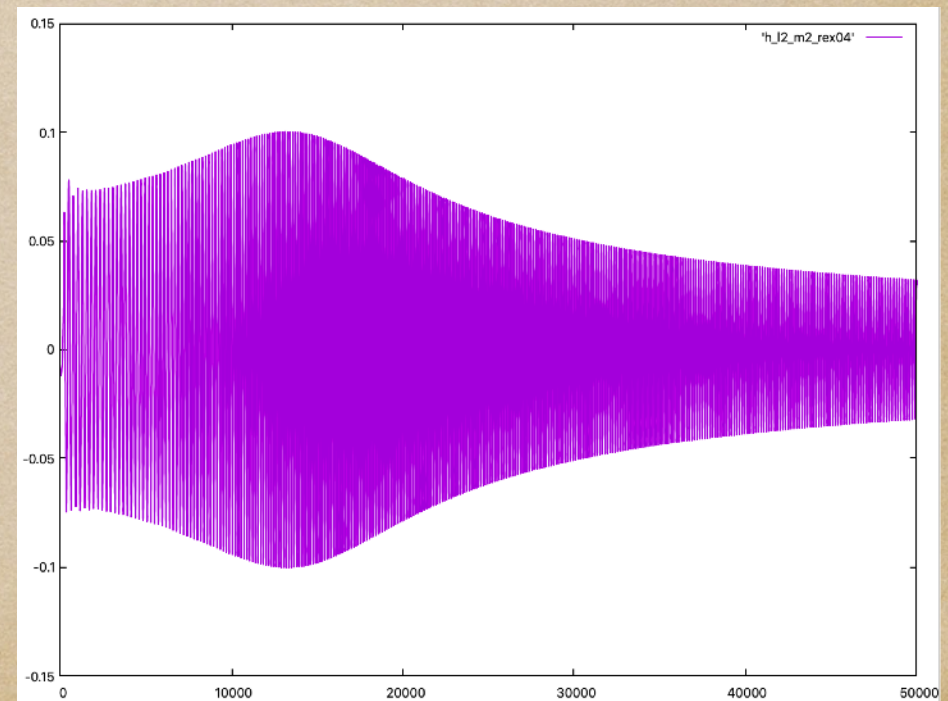
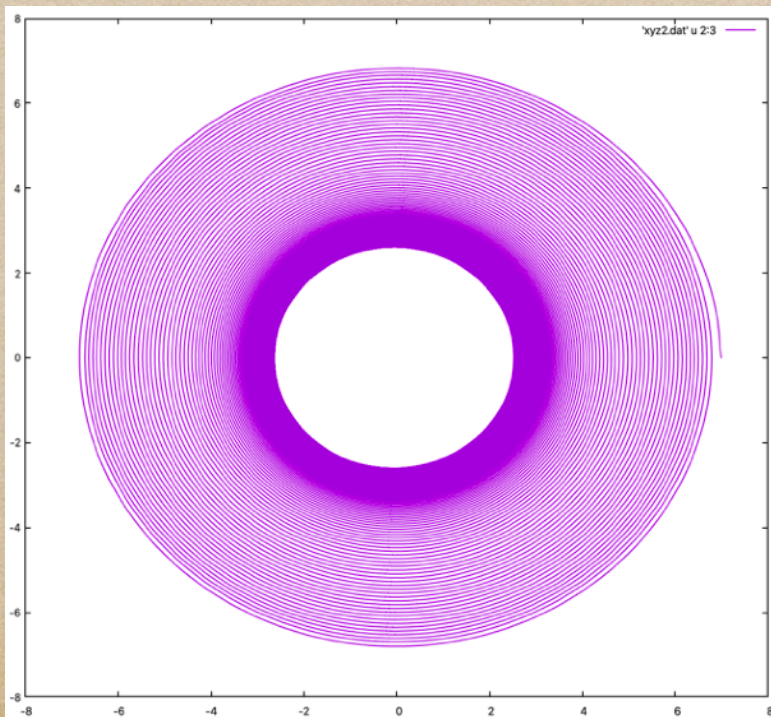
Krishnendu et al 2511.17341

Some "weird" waveform samples from fluffy BSs

$$q = 1,$$
$$\delta\Phi = 120^\circ$$



$$q = 1,$$
$$\delta\Phi = 180^\circ$$

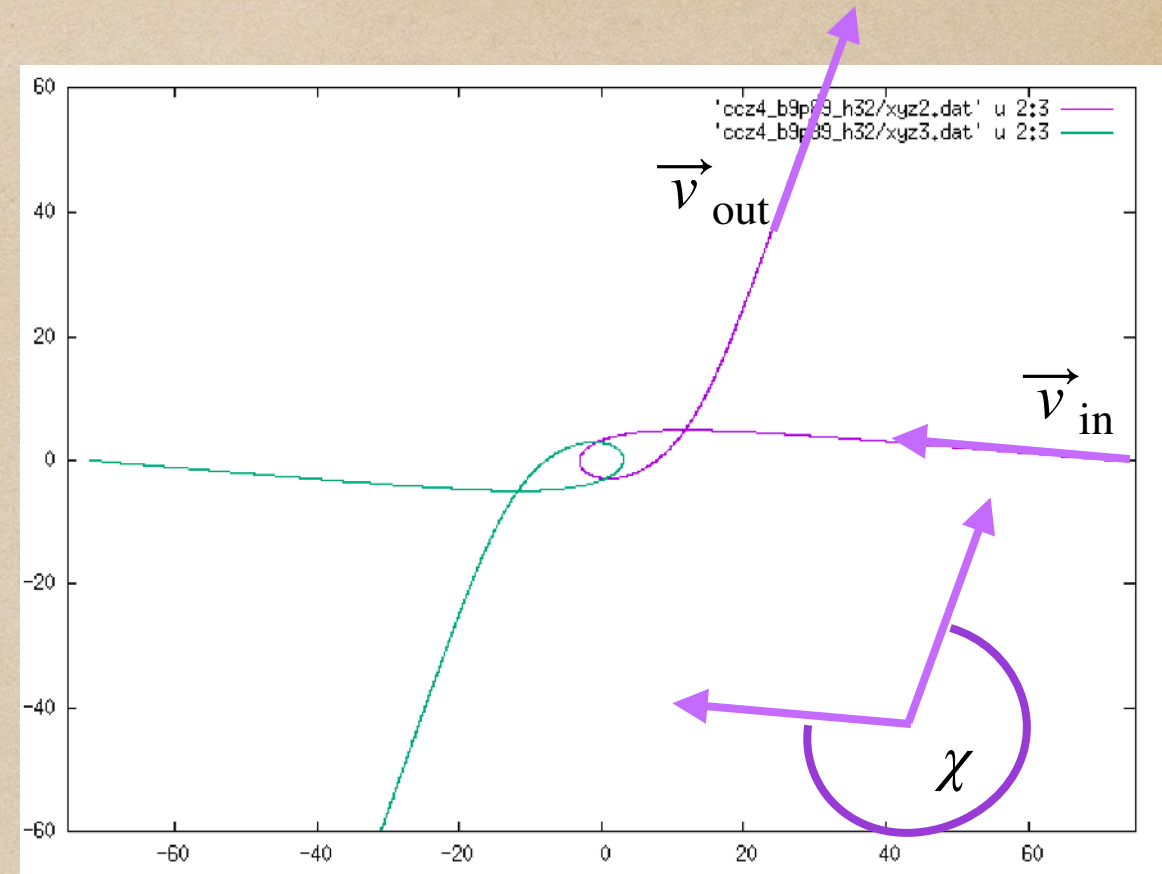


3. EOB modeling of BS scattering

Analytic Modeling of BS Binaries

Damour, Jain & US 2512.00945

- Consider 2 scattering BSs
- NR simulations with Lean;
error $\Delta\chi \approx 1.6^\circ$
- 3 contributions:
 - Point particle BH
 - Tidal "tid"
 - Scalar interaction "scl"



- PM scattering angle: $\chi(\gamma, j) = \chi^{\text{BH}}(\gamma, j) + \chi^{\text{tid}}(\gamma, j) + \chi^{\text{scl}}(\gamma, j)$

EOB potential: $w(\gamma, \bar{r}) = w^{\text{BH}}(\gamma, \bar{r}) + w^{\text{tid}}(\gamma, \bar{r}) + w^{\text{scl}}(\gamma, \bar{r})$

- EOB map: $\pi + \chi(\gamma, j) = 2j \int_0^{\bar{u}_{\text{max}}(\gamma, j)} \frac{d\bar{u}}{\sqrt{\gamma^2 - 1 + w(\gamma, \bar{u}) - j^2 \bar{u}^2}}, \quad \bar{u} = \frac{1}{\bar{r}}$

Damour 1710.10599

The analytic potential

- BH:

- χ^{BH} known up to 4PM Kälin, Liu & Porto 2008.06047

- w^{BH} up to 4PM Damour & Rettegno 2211.01399

- Numerically fitted 5PM term Rettegno et al 2307.06999

- Tidal:

- at 6PM and 7PM using scattering angles from

- Kälin, Liu & Porto 2008.06047, Bini, Damour & Geralico 2001.00352,

- Cheung & Solon 2006.06665

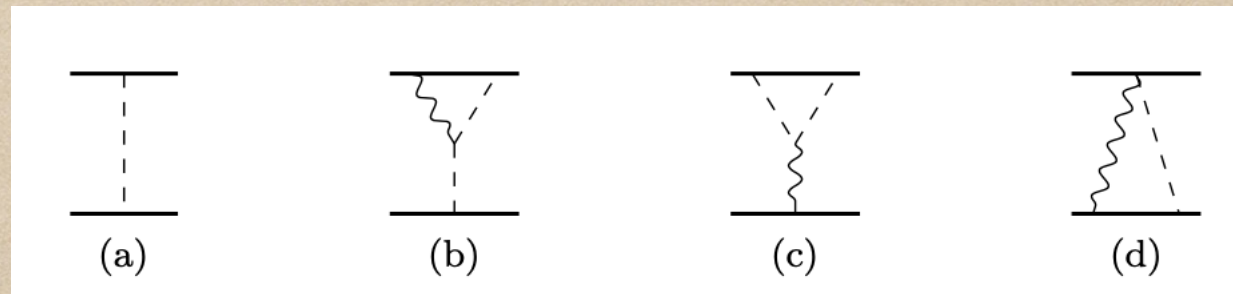
- Scalar: Previously unknown...

The scalar potential contribution

- EFT action for 2BSs as world lines $\mathbf{z}_A(\tau_A)$, $A = 1, 2$, with source

$$s_A(\tau_A) = c_A e^{i\epsilon_A \omega_A \tau_A + i\Phi_A}, \quad c_A = \text{scalar charges}$$

- Perturbative analysis $\rightarrow$ Feynman diagrams



- Leading Order from (a): $\chi_{\text{LO}}^{\text{scl}}(\gamma, j) = \frac{8\pi M h \tilde{\mu}}{m_1 m_2} \frac{c_1 c_2}{\gamma^2 - 1} K_1(\tilde{\mu} |b|) \cos(\Phi_2 - \Phi_1)$,

where $h = \sqrt{1 + 2\nu(\gamma - 1)}$, $\tilde{\mu}^2 = \mu^2 - \frac{2\omega_1 \omega_2 \gamma - \omega_1^2 - \omega_2^2}{\gamma^2 - 1}$

plus a negligible nonconservative part for $\Phi_1 \neq \Phi_2$

$$\rightarrow w_{\text{LO}}^{\text{scl}}(\gamma, \vec{r}) = \frac{8\pi c_1 c_2}{G m_1 m_2} \frac{e^{-GMh\tilde{\mu}\vec{r}}}{\vec{r}} \cos(\Phi_2 - \Phi_1)$$

... continued

- Higher-order contribution $w_{\text{HO}}^{\text{scl}}$ not fully known;

we assume a factorized form

$$w_{\text{HO}}^{\text{scl}} = w_{\text{g-dressing}}^{\text{scl}} \times w_{\text{LO}}^{\text{scl}}, \quad w_{\text{g-dressing}}^{\text{scl}} = e^{GM((2\omega^2 - \mu^2)/\tilde{\mu}) \log(2\gamma_E GM\tilde{\mu}\tilde{r})}$$

- Input needed:

- total energy E

- total angular momentum $J = 1.006 J_{\text{NR}}$

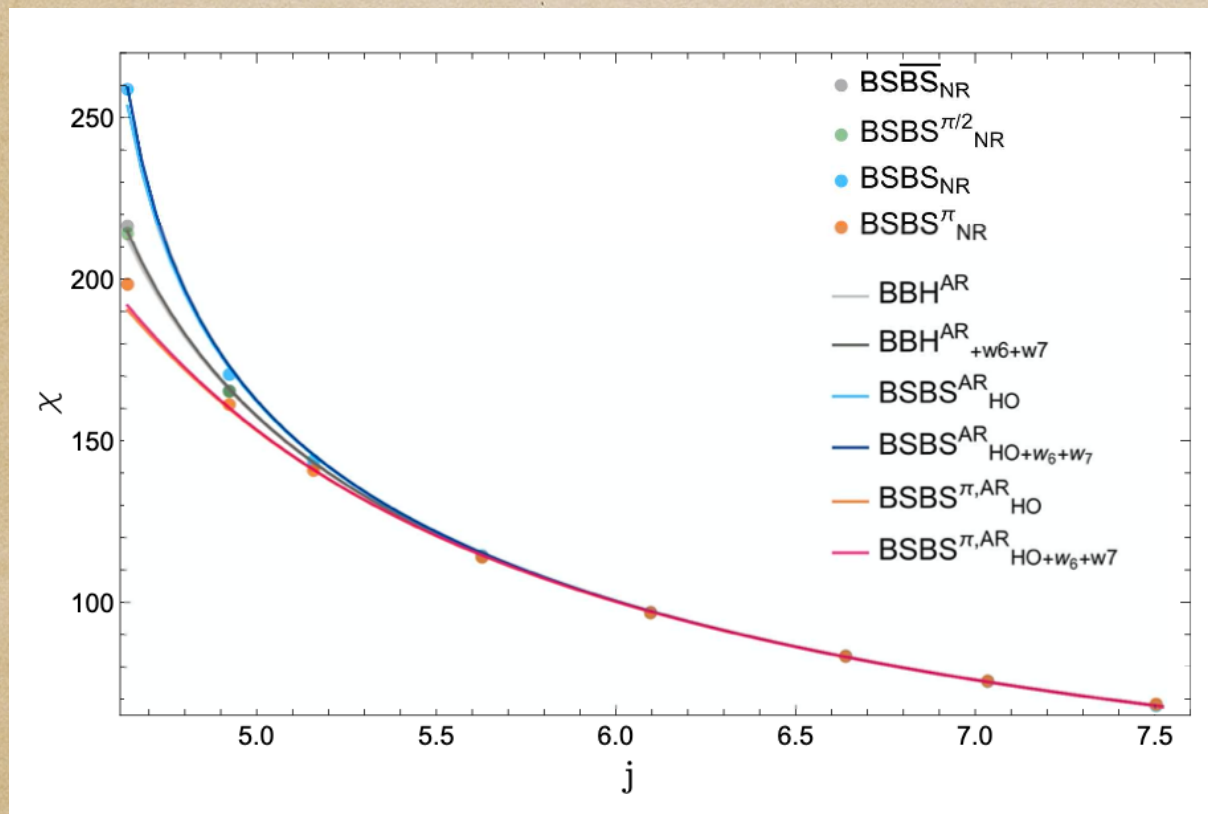
- tidal deformability parameters $\Lambda_1 = \Lambda_2 \approx 7$

Sennett et al 1704.08651

- scalar charges c_1, c_2 from BS asymptotics in isotropic radius

Results

- Equal-mass BS binaries with $\sigma_0 = 0.2$, $|\varphi_c| = 0.17$
- Four binaries: BSBS, BS $\overline{\text{BS}}$, $\Phi_2 - \Phi_1 = \frac{\pi}{2}$, $\Phi_2 - \Phi_1 = \pi$



- Scalar corrections dominant
- Tidal effects still important for small impact parameter
- Scalar contribution attractive (repulsive) for $\Phi_2 - \Phi_1 = 0$ ($\Phi_2 - \Phi_1 = \pi$)
- BS $\overline{\text{BS}}$ close to BH binaries

5. Mixed binaries

BH-BS Binaries

- "Piercing papers"

Cardoso et al PRD 2022, Zhong et al PRD 2023

BH passing through a fluffy mini or solitonic BS:

$\gtrsim 90\%$ of BS mass accreted, even for light BHs (test particle limit)

- Head-on collisions and orbital evolutions of BH-BS binaries

Marks, Staelens & Sperhake in prep.

- GRChombo code

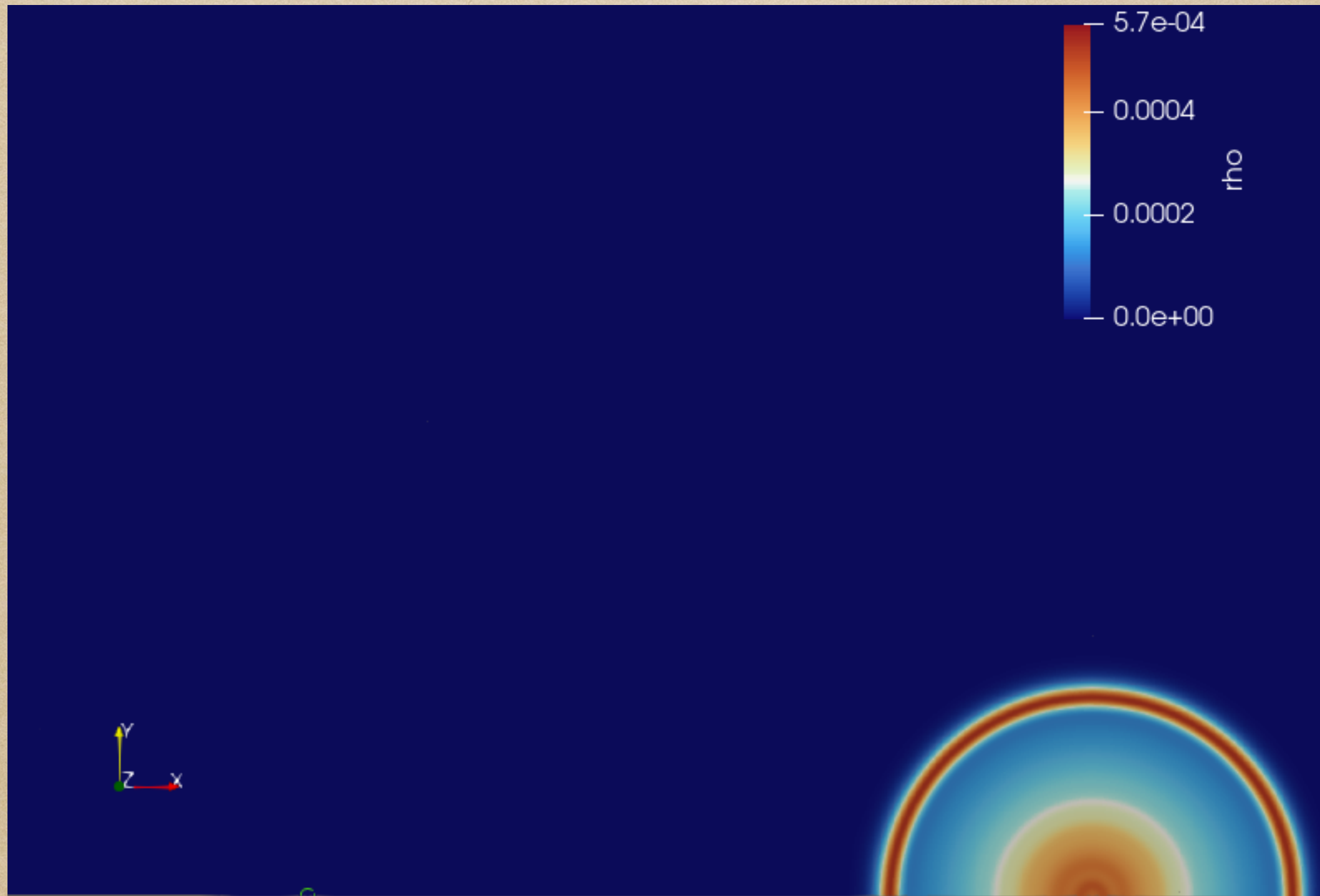
- CCZ4 formulation

- Constant-Volume-Element initial data

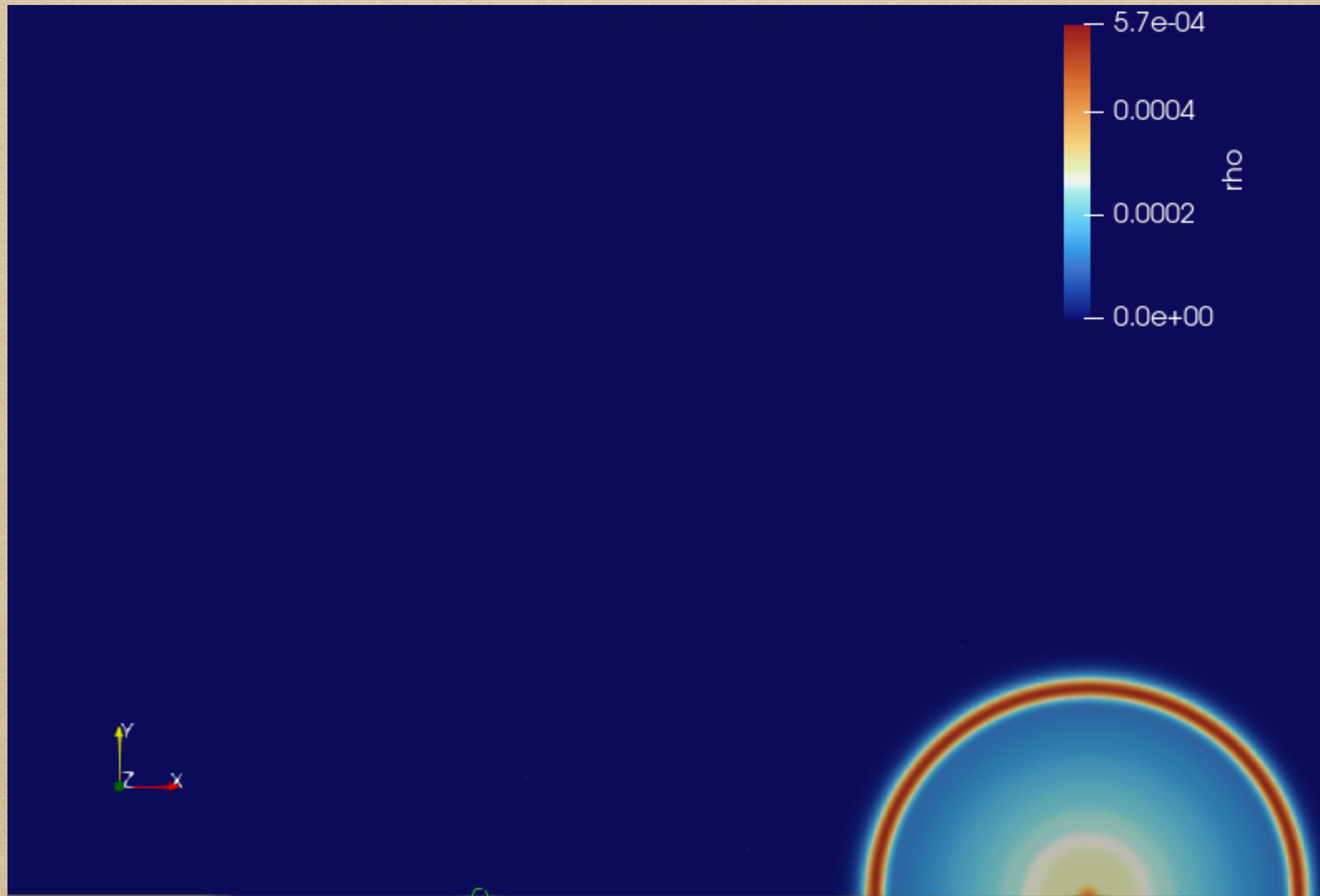
- Example: Thin-Shell BS Head-on collision with BH

Credit to Gareth Marks!!!

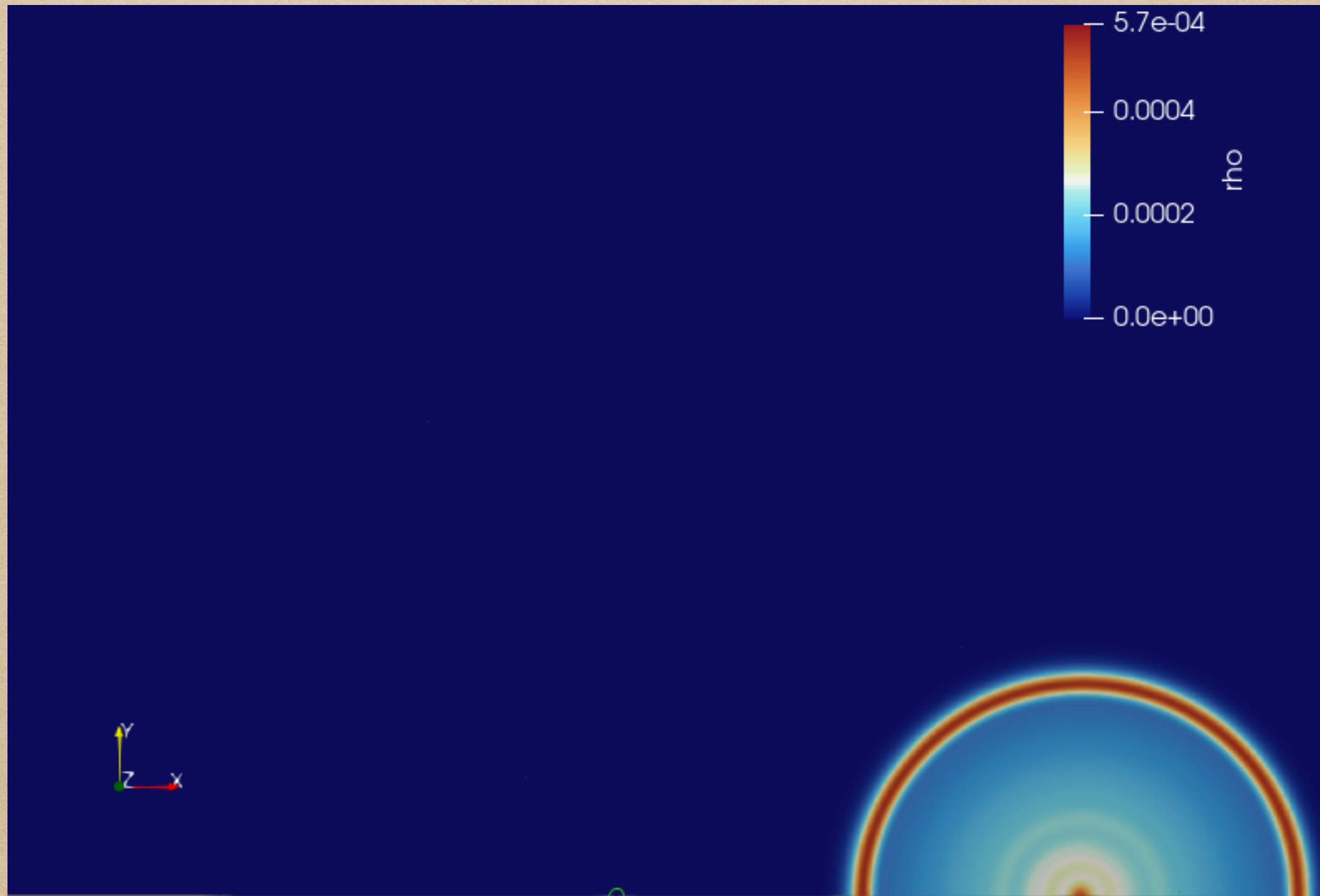
BH-BS Binaries



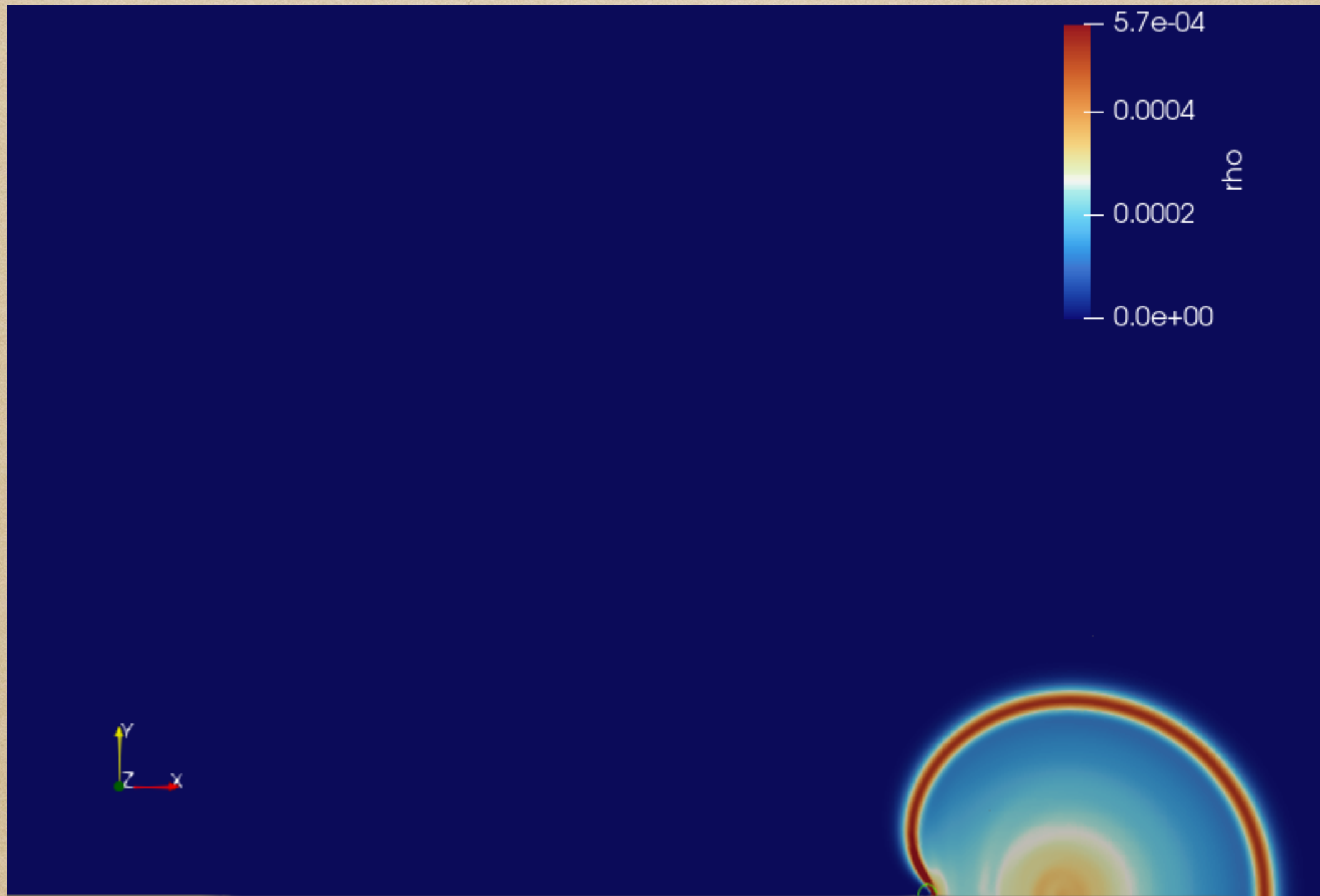
BH-BS Binaries



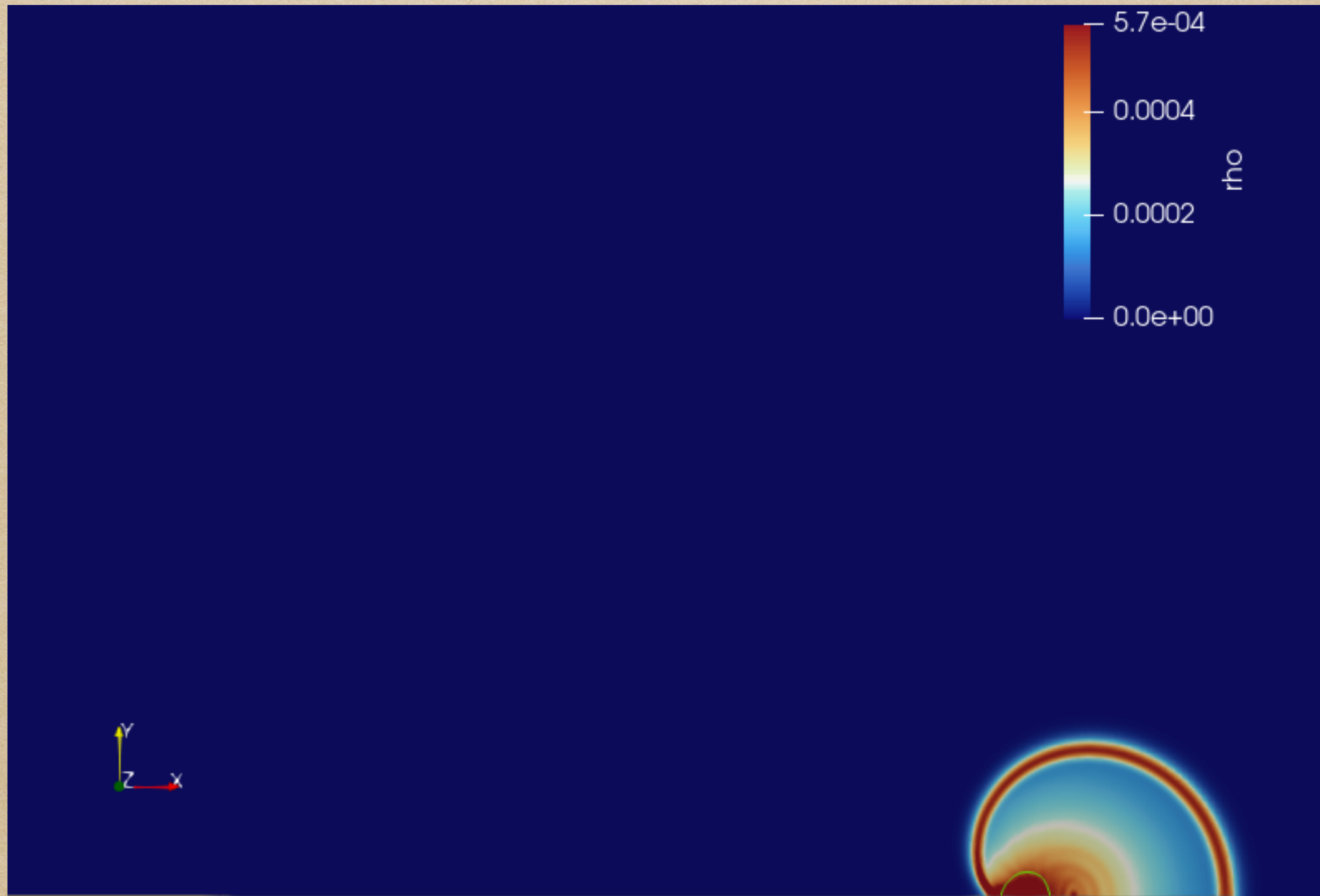
BH-BS Binaries



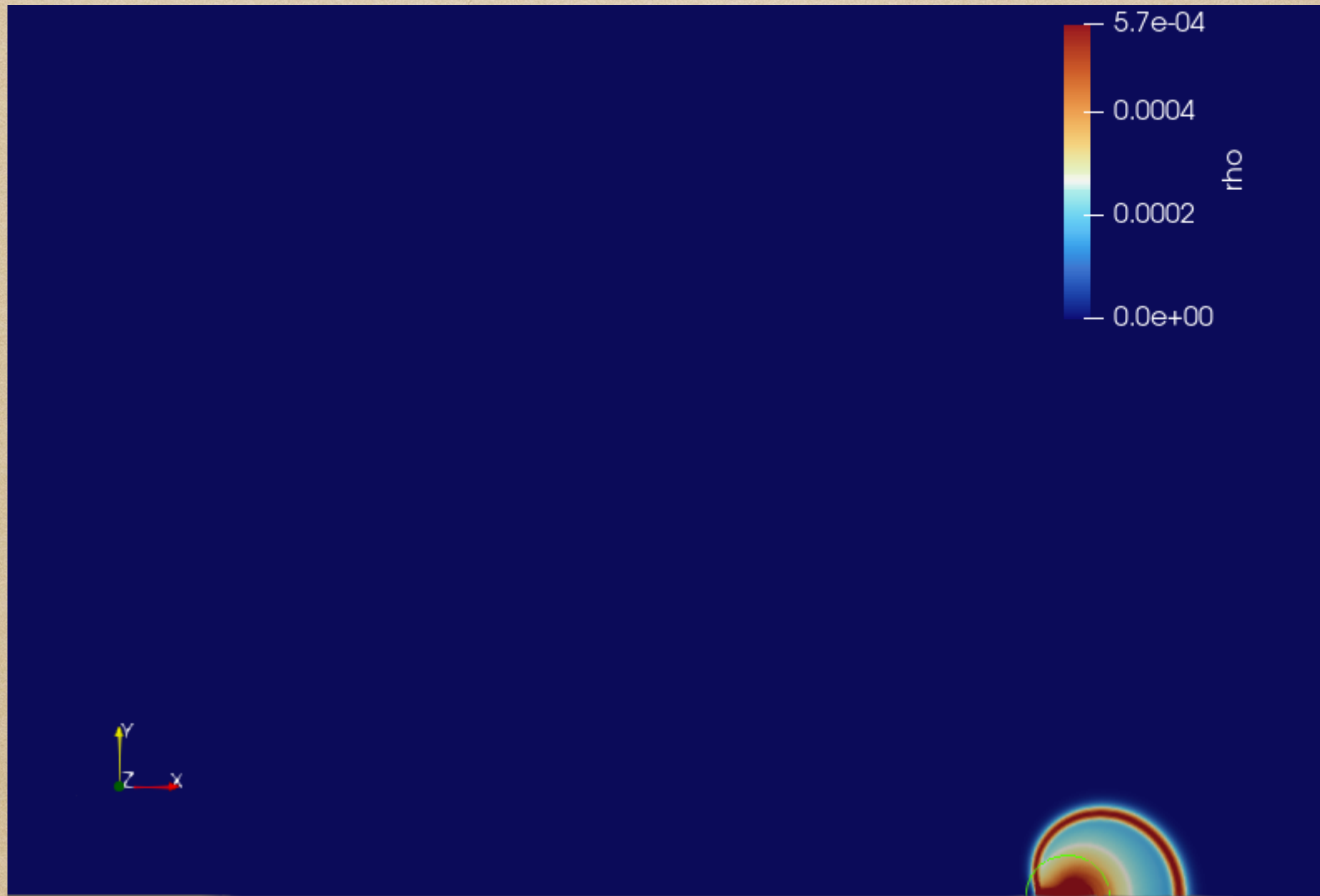
BH-BS Binaries



BH-BS Binaries



BH-BS Binaries



6. Summary

Conclusions

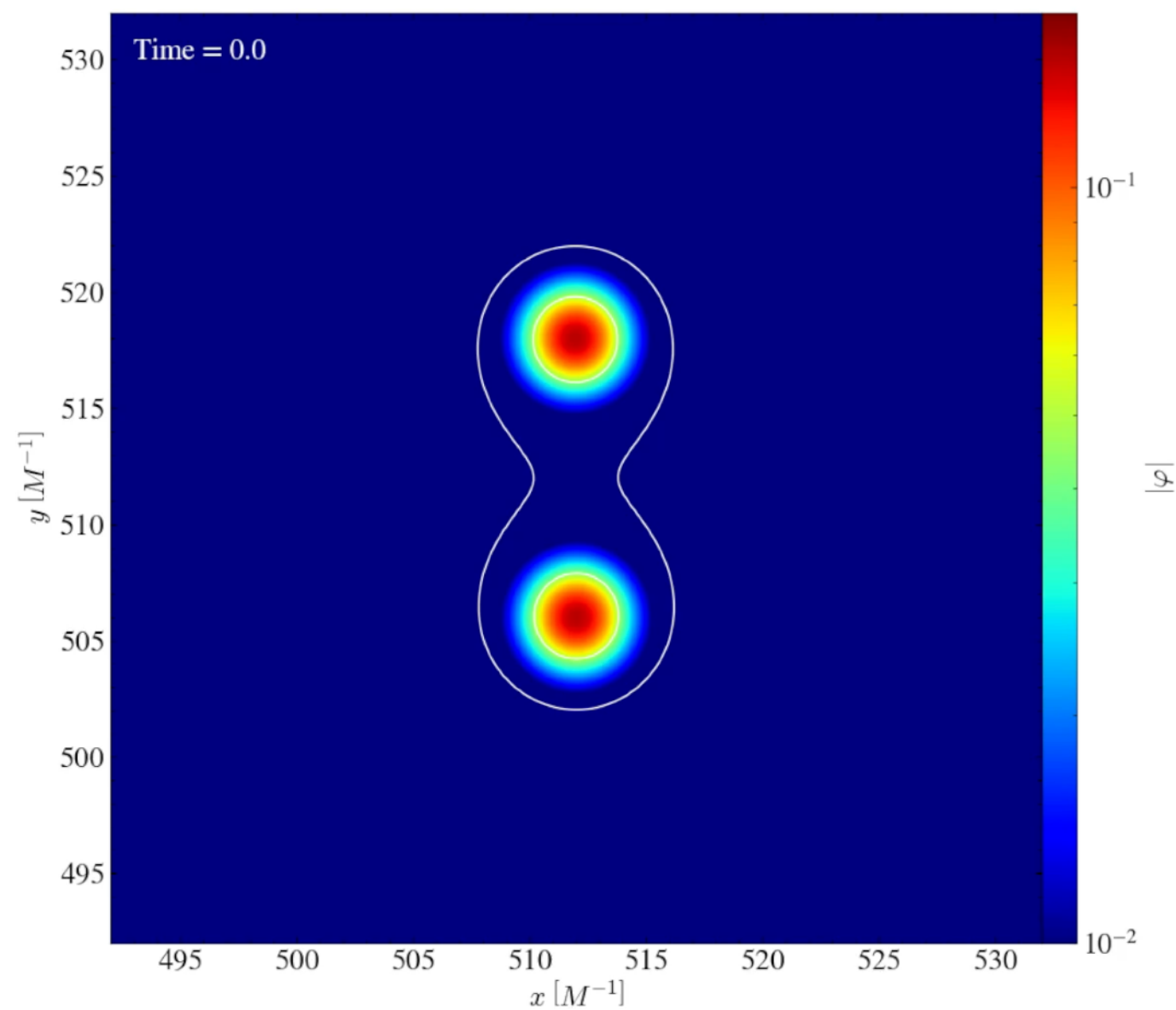
- BS highly amenable to NR! (Perhaps only ECO to be so...)
- Wide range of compactness and mass
- Stable and unstable branches similar to neutron stars
- Spin quantized; may form...
- Compact BBSs can look similar to BBHs → BS-BH degeneracy
- Three main contributions to GW forms: BH, Tidal, Field interaction
- Genuine BH Mimickers: Light ring + no horizon; seem to be stable
- Ergoregion instability

Next Challenges

- Need systematic NR waveforms
- Need for improved initial data!
- Need analytic waveform models
- Various stability questions: Spinning BSs, BH mimickers, ergoregion
- Mixed binary modeling in its infancy (foetus more like it...)

4. Extra slides

Example boson-star inspiral



Courtesy of T Evstafyeva

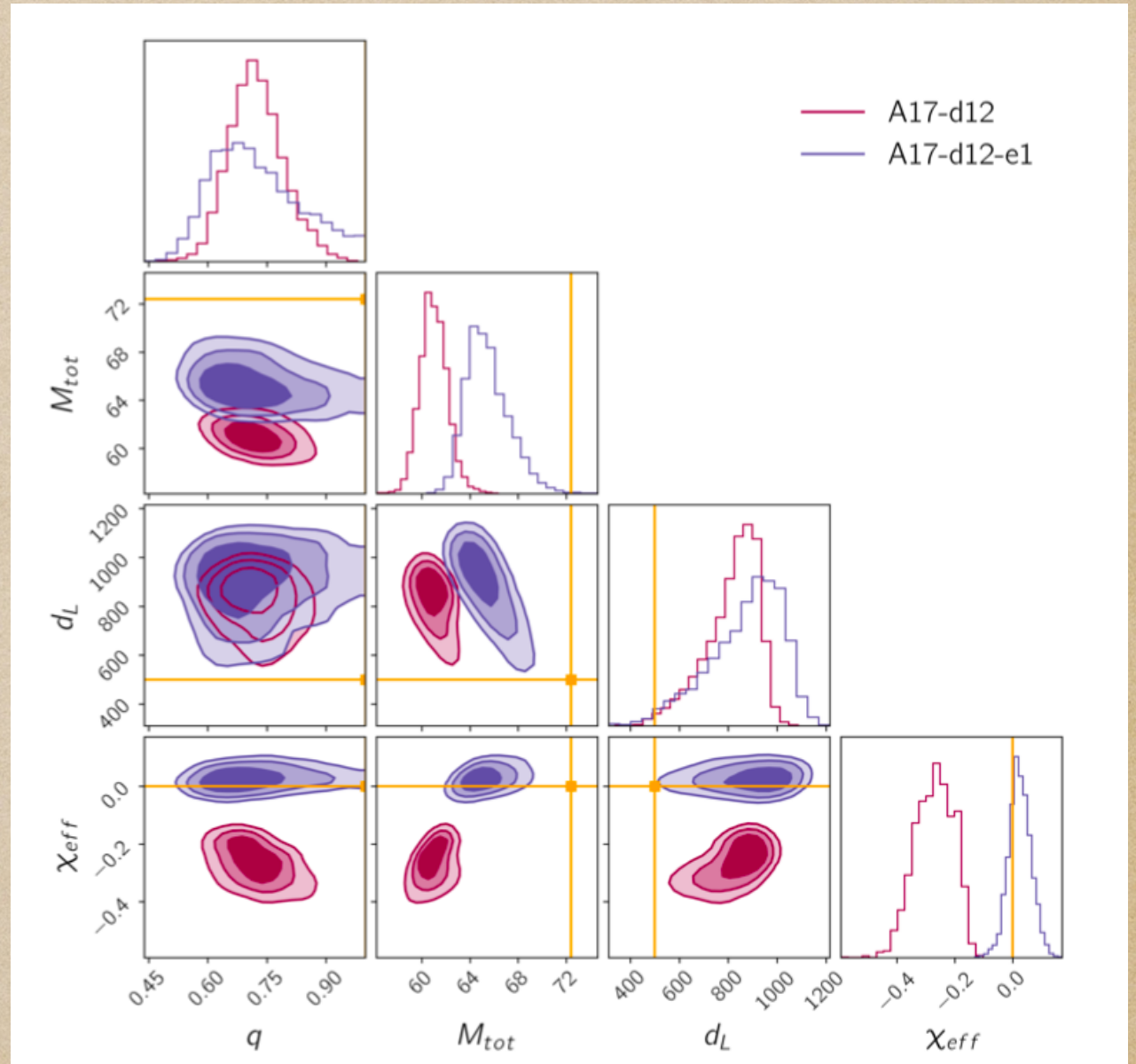
Results using IMRPhenomXP

- BBH approximants recover parameters best for *anti-BS* !!

Corner plot:

A17-d12-e1 vs. A17-d12

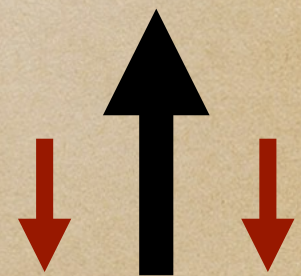
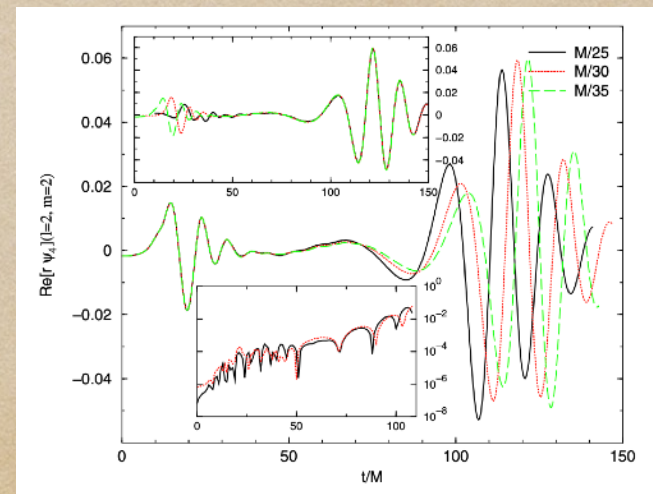
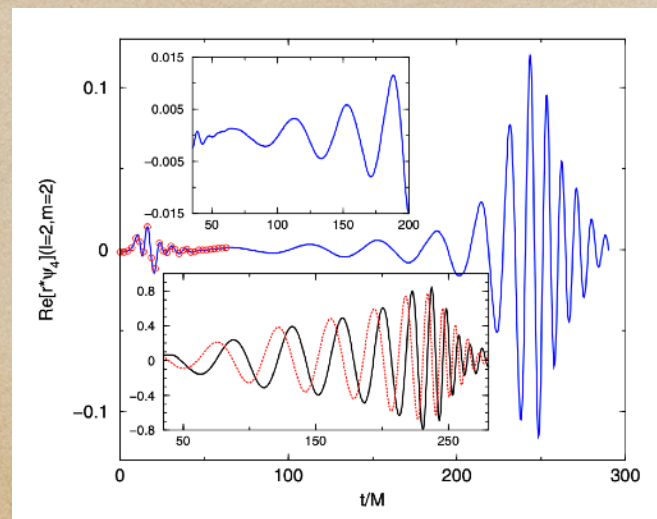
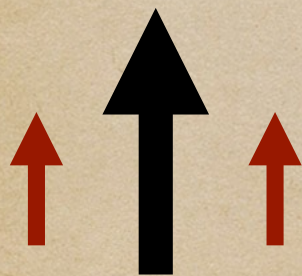
- These features can be explained with the chirp strength



Understanding the PE bias

- Main feature: Steepness of chirp
- For non-spinning BH binaries:
 - equal mass $\Rightarrow$ shallow chirp
 - unequal masses $\Rightarrow$ steep chirp (think of EMRIs)
- For spinning BH binaries:
 - aligned spins $\Rightarrow$ shallow chirp
 - anti-aligned spins $\Rightarrow$ steep chirp

The orbital 'Hang-up' effect Capanelli et al gr-qc/0601091



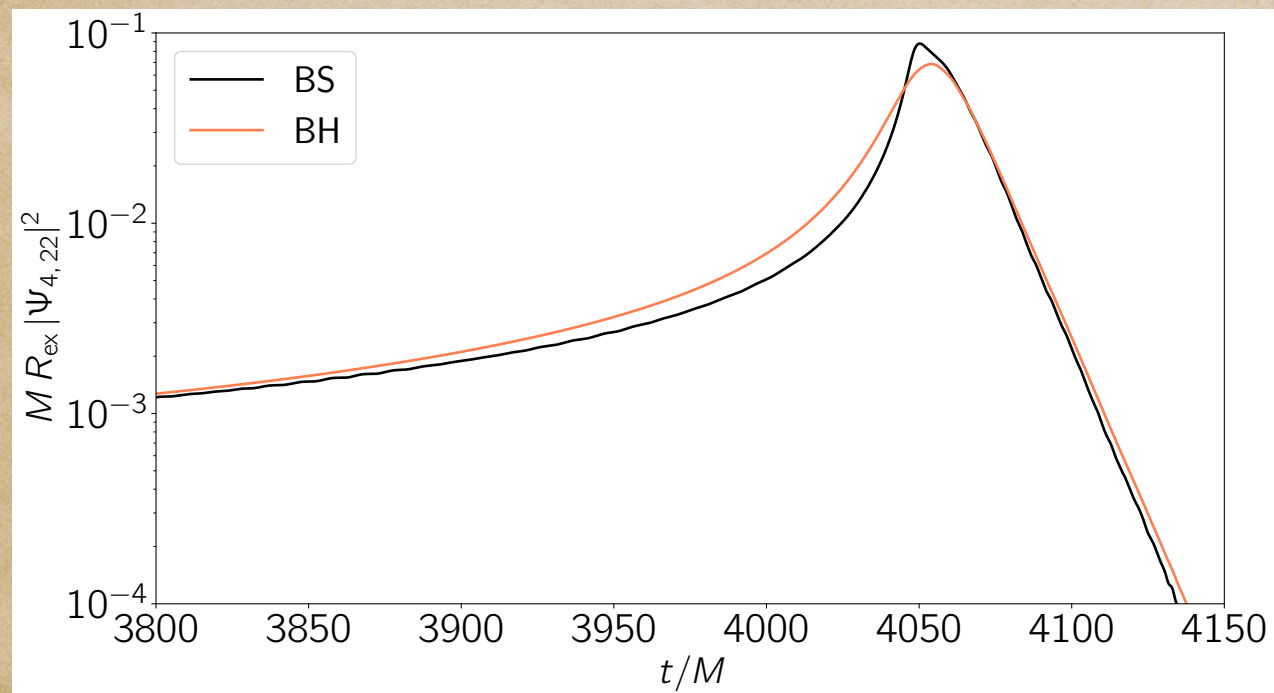
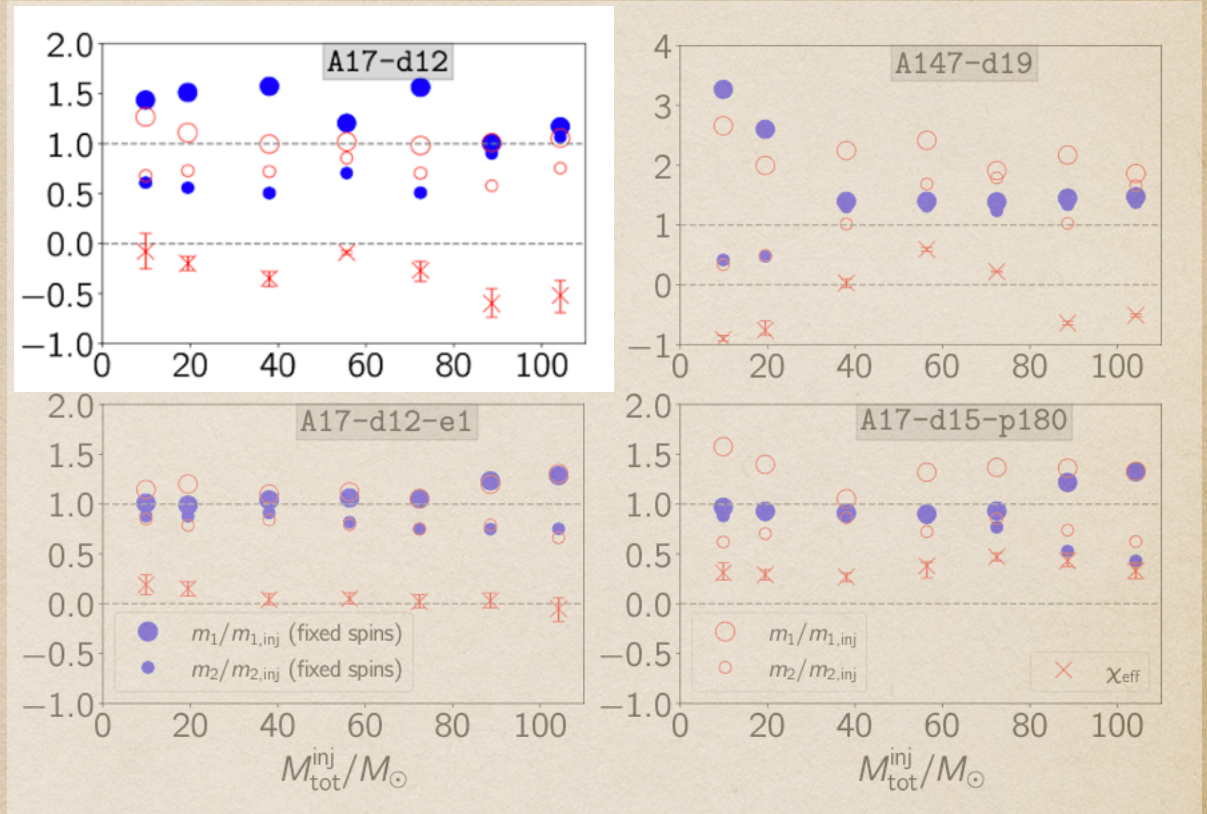
Understanding the PE bias: Standard BS

Fixed spins

BS chirp steeper

⇒ Like unequal-mass BHs

⇒ Bilby reports unequal masses



Variable spins

anti-aligned spins

⇒ Steeper chirp

⇒ Steep BS chirp also captured by anti-aligned spins

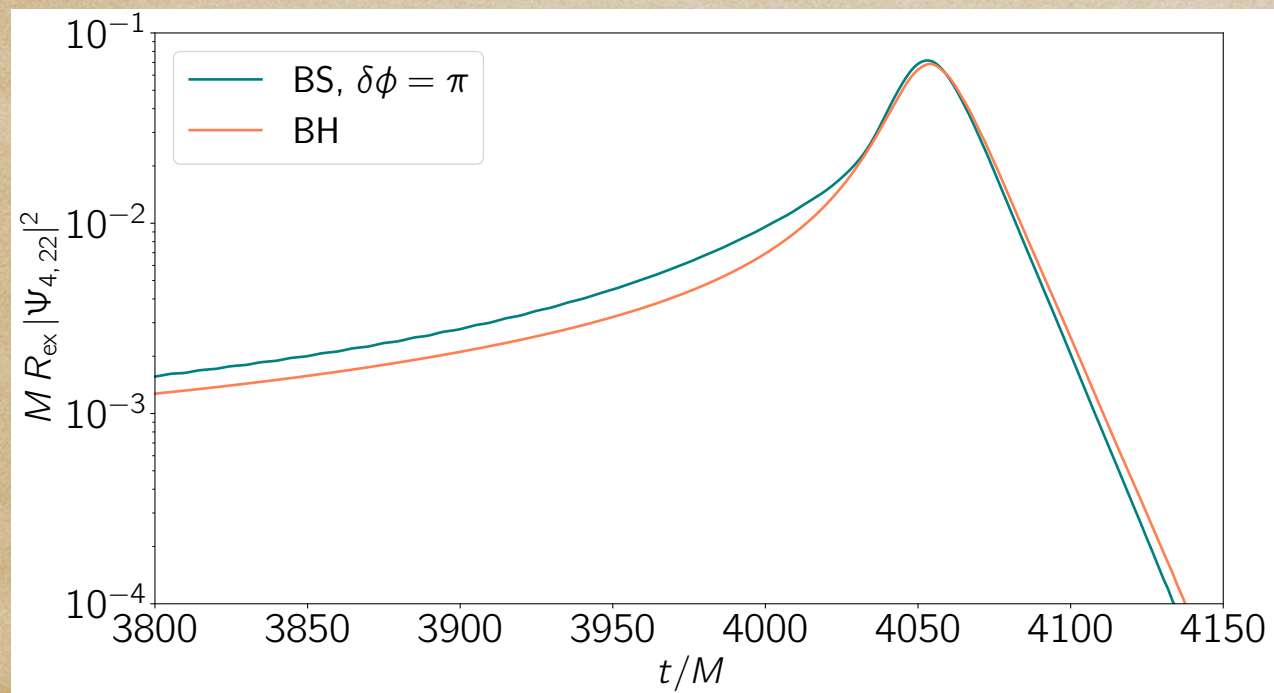
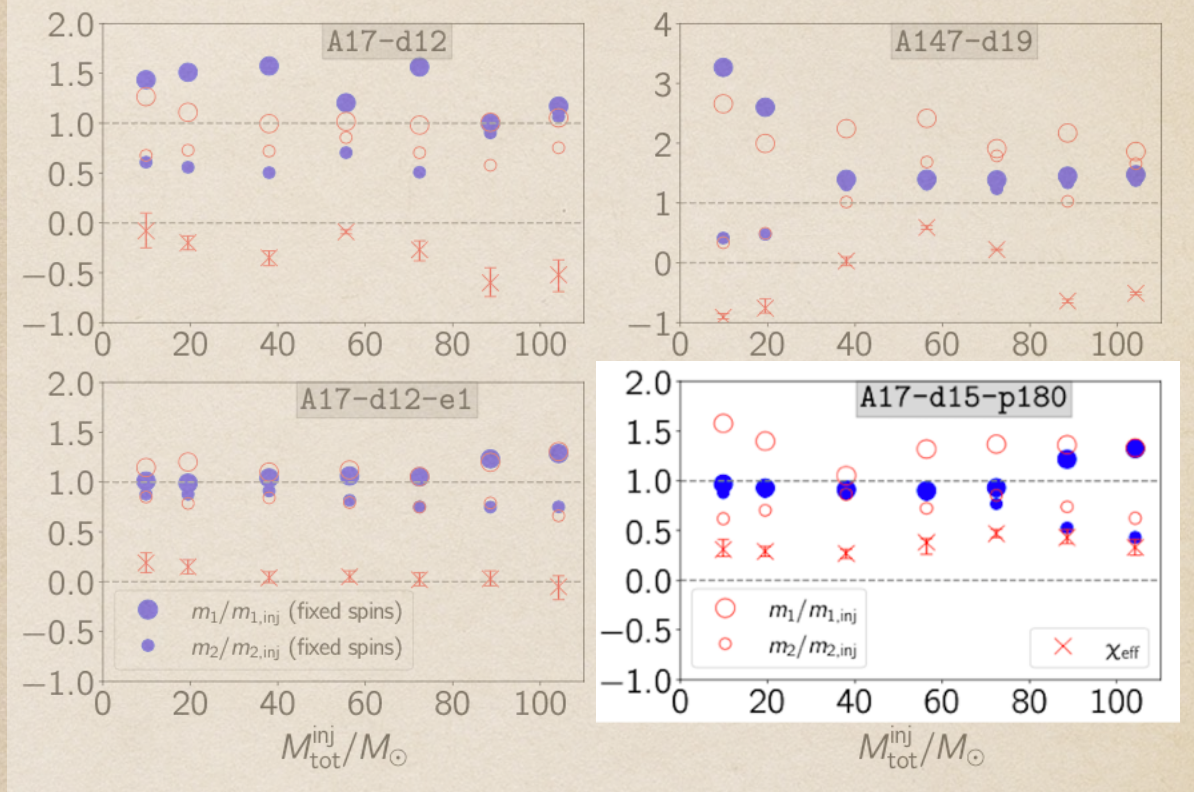
Understanding the PE bias: anti-phase

Fixed spins

BS chirp shallower

⇒ Best matched by $\sim$ equal mass BHs

⇒ Bilby reports $\sim$ equal masses



Variable spins

aligned spins

⇒ Shallow chirp

⇒ Bilby reports aligned spins
and allows unequal masses

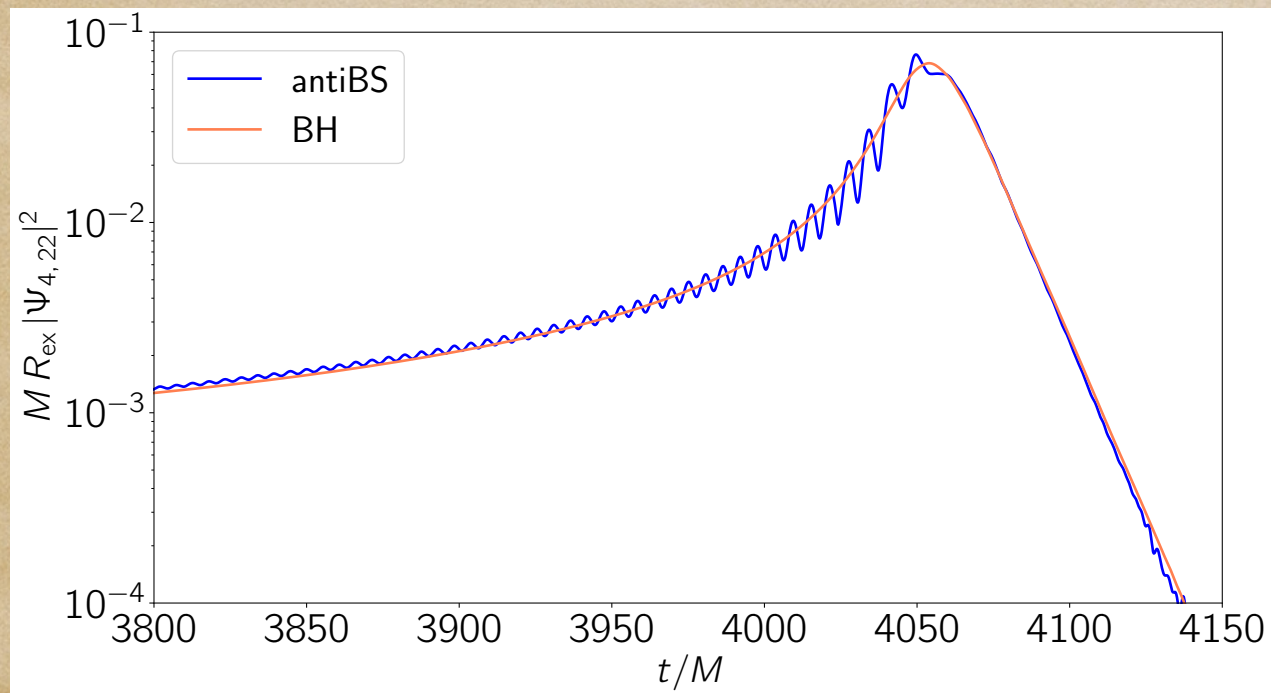
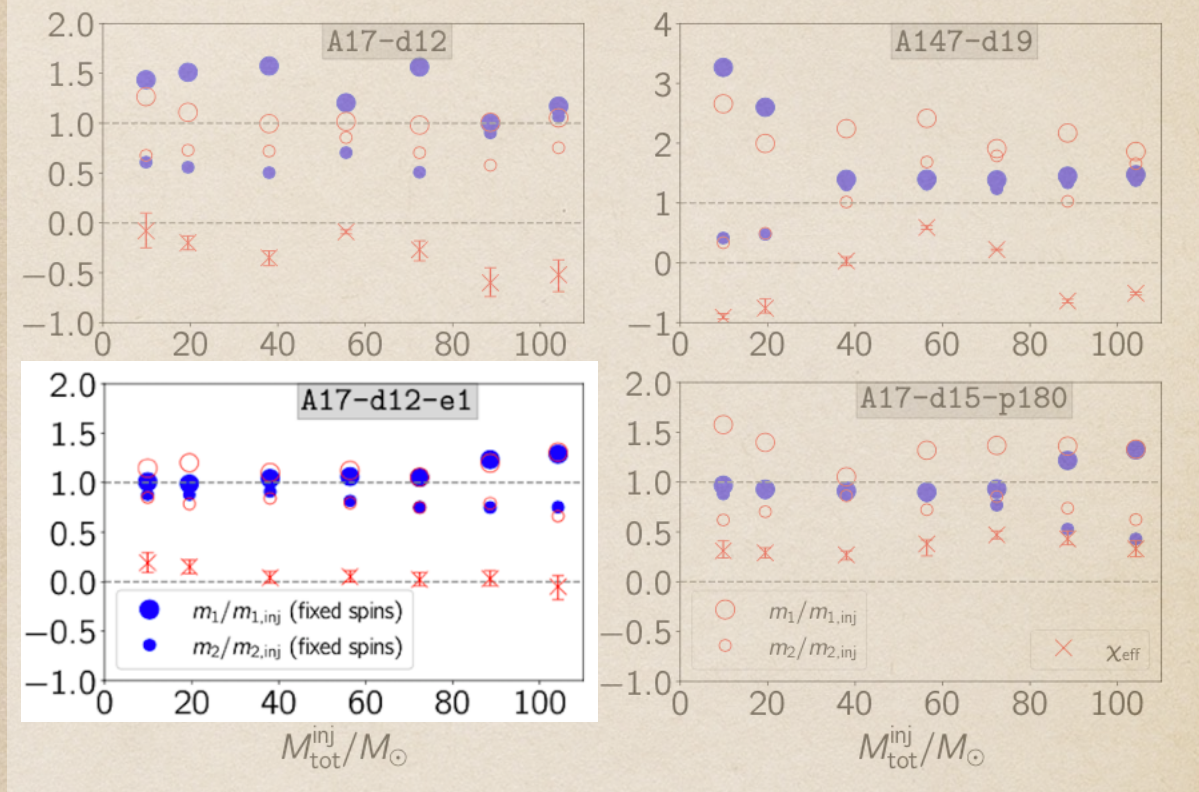
Understanding the PE bias: anti-BS

Fixed or variable spins

BS chirp similar to BHs

⇒ Comparable mass ratios

∧ Small spins



High-mass regime

LIGO mainly sees merger burst

⇒ Less reliable PE

Convergence

