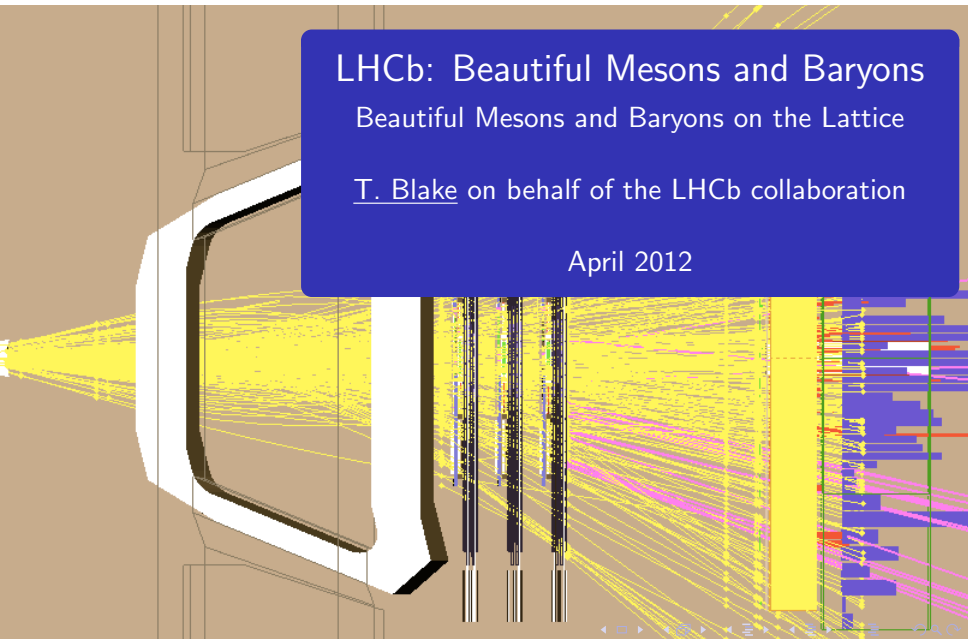


# LHCb: Beautiful Mesons and Baryons

Beautiful Mesons and Baryons on the Lattice

T. Blake on behalf of the LHCb collaboration

April 2012



- Will try to discuss briefly recent LHCb results for:
  1. Rare decay processes.
  2.  $B_s$  mixing observables.
  3. Mass measurements and spectroscopy.

$$B_{(d,s)}^0 \rightarrow \mu^+ \mu^-$$

$$B_s \rightarrow \mu^+ \mu^- \text{ and } B_d \rightarrow \mu^+ \mu^-$$

- Sensitive to contributions from scalar + pseudo-scalar sector.

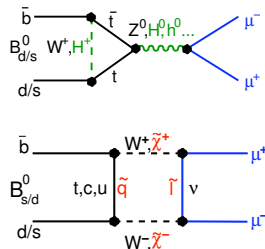
→ Interesting to probe NP models with extended Higgs sector, e.g. MSSM, 2HDM, ...

e.g. in MSSM, branching fraction scales approximately as  $\tan^6 \beta / M_A^4$

- More generally:

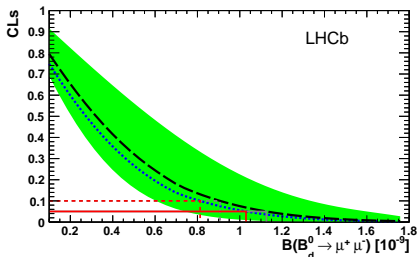
$$\mathcal{B}(B_q^0 \rightarrow \mu^+ \mu^-) \simeq \frac{G_F^2 \alpha^2 M_{B_q^0}^3 f_{B_q^0}^2 \tau_{B_q^0}}{64 \pi^3 \sin^4 \theta_W} |V_{tb} V_{tq}^*|^2 \left( 1 - \frac{4m_\mu^2}{M_{B_q^0}^2} \right)^{1/2} \times$$

$$\left[ M_{B_q^0}^2 \left( 1 - \frac{4m_\mu^2}{M_{B_q^0}^2} \right) |C_S|^2 + \left( M_{B_q^0} C_P + \frac{2m_\mu}{M_{B_q^0}} C_{10} + \frac{2m_\mu}{M_{B_q^0}} C'_{10} \right)^2 \right]$$

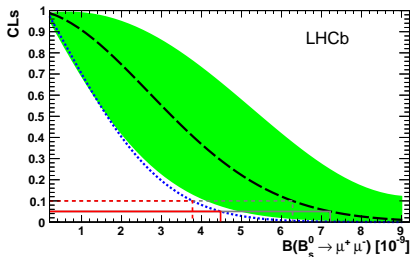


- Set limit on the branching fraction using the CLs technique, dividing data into bins of BDT response and mass.
- BDT response and mass line-shape of the signal calibrated from data using  $B \rightarrow hh'$  and  $J/\psi/\psi(2S)/\Upsilon(1S) \rightarrow \mu^+ \mu^-$  decays.

$$\mathcal{B}(B_d \rightarrow \mu^+ \mu^-) < 1.0 \times 10^{-9} \text{ (95\% C.L.)}$$



$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 4.5 \times 10^{-9} \text{ (95\% C.L.)}$$



- Expected limit  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 7.2 \times 10^{-9}$  (bkg + SM)

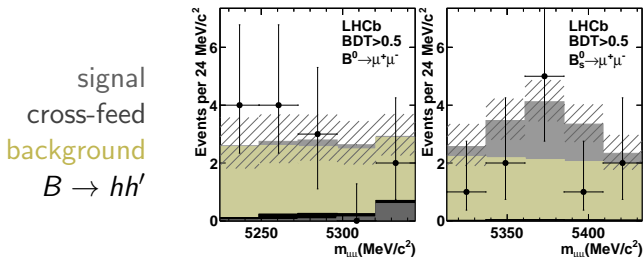
c.f.  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 22 \times 10^{-9}$  [ATLAS-CONF-2012-010]

$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 7.7 \times 10^{-9}$  [CMS-BPH-11-020]

- LHCb sets a limit for  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 4.5 \times 10^{-9}$  (95% C.L.)
- c.f. SM expectation of  $(3.2 \pm 0.2) \times 10^{-9}$ .
- Best fit branching fraction estimated to be:

$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) = (0.8^{+1.8}_{-1.3}) \times 10^{-9}$$

using a simultaneous maximum likelihood fit to the  $\mu^+ \mu^-$  mass distribution in the 8 BDT bins.



- CDF:  $(1.0^{+0.8}_{-0.6}) \times 10^{-8}$

- Experimental limits are now close to the SM prediction.
- May become increasingly important to understand the uncertainty on (and the central value for) the SM prediction.
- Several predictions available:

$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) = (3.2 \pm 0.2) \times 10^{-9} \quad (f_{B_s} = 225 \pm 4 \text{ MeV}).$$

$$\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) = (3.7 \pm 0.4) \times 10^{-9} \quad (f_{B_s} = 250 \pm 12 \text{ MeV}).$$

- For details see A. Buras [[PLB 566 \(2003\)](#)], with  $f_{B_s}$  from J. Laiho et al. [[PRD 81 \(2010\)](#)] and C. McNeile et al. [[arXiv:1110.4510](#)].

$b \rightarrow s\ell^+\ell^-$  processes



- Measure  $\mathcal{B}(B_s \rightarrow \phi \mu^+ \mu^-)$  normalised using  $B_s \rightarrow J/\psi \phi$  decays.
- With simple cut based analysis and strong PID requirements, observe:

$77 \pm 10$   $B_s \rightarrow \phi \mu^+ \mu^-$  candidates.

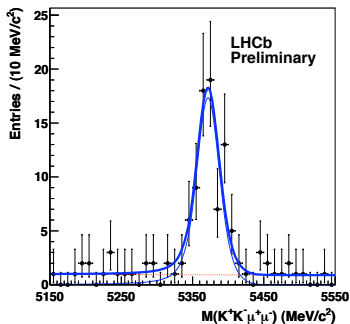
- Giving:

$$\mathcal{B}(B_s \rightarrow \phi \mu^+ \mu^-) = (0.78 \pm 0.10(\text{stat}) \pm 0.06(\text{syst}) \pm 0.28(\mathcal{B})) \times 10^{-6}$$

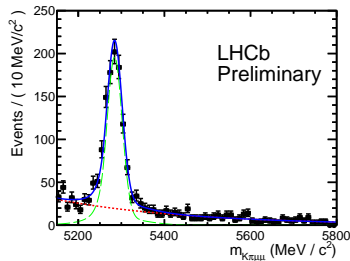
(LHCb Preliminary)

- Lower than but still statistically compatible with the CDF result [PRL 107 (2011)].

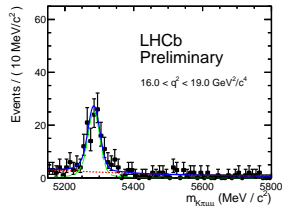
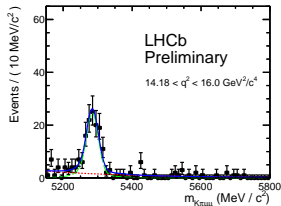
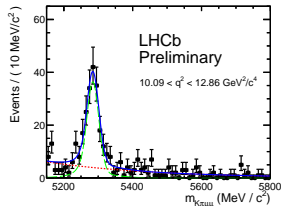
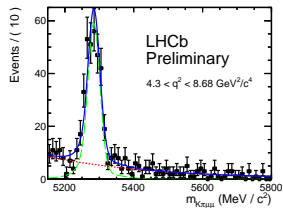
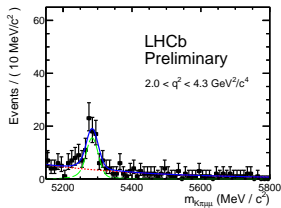
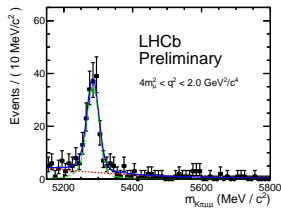
NB Can get to  $\sim 10\%$  normalising to  $B_d$  modes.



- Gives access to  $C_7^{(\prime)}$ ,  $C_9^{(\prime)}$  and  $C_{10}^{(\prime)}$ .
- Sensitivity through branching fraction measurements and angular observables. Angular analysis can be sensitive to the chiral structure of the NP (e.g.  $S_3/A_T^2$ ).
- Decay described by three angles and dimuon invariant mass squared ( $\theta_\ell, \theta_K, \phi, q^2$ ).
- Analysis based on  $1 \text{ fb}^{-1}$ .
- Observe 900 candidates ( $BABAR + Belle + CDF \sim 600$ ).



# Yield in $q^2$ -bins



# Anatomy of the $B_d \rightarrow K^{*0} \mu^+ \mu^-$ decay

- Anatomy of the decay was explained in detail by Christoph yesterday. Described by 12 angular terms:

$$\frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\phi dq^2} \propto [J_1^s + J_1^c + (J_2^s + J_2^c) \cos 2\theta_\ell + J_3 \sin^2 \theta_\ell \cos 2\phi + J_4 \sin 2\theta_\ell \cos \phi + J_5 \sin \theta_\ell \cos \phi + J_6 \cos \theta_\ell + J_7 \sin \theta_\ell \sin \phi + J_8 \sin 2\theta_\ell \sin \phi + J_9 \sin^2 \theta_\ell \sin 2\phi]$$

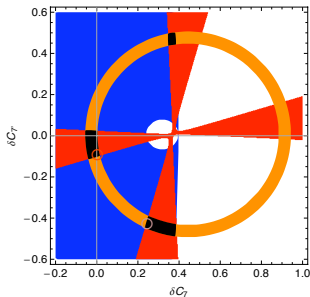
- By folding in  $\phi$  can cancel terms in the angular expression:
  - Leaving  $J_1^{s,c}$ ,  $J_2^{s,c}$ ,  $J_3$ ,  $J_6$  and  $J_9$ .
- Relate  $J_1^{s,c}$  &  $J_2^{s,c}$  to  $F_L$  and  $F_T = 1 - F_L$  in the massless case.
- Some possibility to access  $J_4$ ,  $J_5$ ,  $J_7$  and  $J_8$  using other folding techniques. See A. Bharucha & W. Reece [[arXiv:1002.4310](https://arxiv.org/abs/1002.4310)].

- Take advantage of a symmetry of the system, by folding the distribution in  $\phi$ , to reduce the number of free parameters.
- $\hat{\phi} = \phi + \pi$  if  $\phi < 0$  and  $\hat{\phi} = \phi$  if  $\phi > 0$ , leading to:

$$\frac{1}{\Gamma} \frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\hat{\phi} dq^2} = \frac{9}{16\pi} \left[ F_L \cos^2\theta_K + \frac{3}{4}(1 - F_L)(1 - \cos^2\theta_K) + \right. \\ F_L \cos^2\theta_K (2\cos^2\theta_\ell - 1) + \\ \frac{1}{4}(1 - F_L)(1 - \cos^2\theta_K)(2\cos^2\theta_\ell - 1) + \\ S_3(1 - \cos^2\theta_K)(1 - \cos^2\theta_\ell)\cos 2\hat{\phi} + \\ \frac{4}{3}A_{FB}(1 - \cos^2\theta_K)\cos\theta_\ell + \\ \left. A_{Im}(1 - \cos^2\theta_K)(1 - \cos^2\theta_\ell)\sin 2\hat{\phi} \right]$$

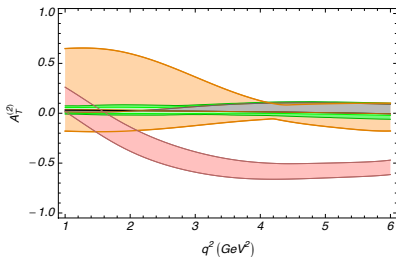
# New physics sensitivity to $\mathcal{C}_7^{(\prime)}$

- $\mathcal{C}_7$  and  $\mathcal{C}_7'$  are constrained by  $b \rightarrow s\gamma$  processes. Even in the SM-like allowed region can still have large sensitivity to  $\mathcal{C}_7'$  through  $A_T^2$ .
- Where  $S_3$  is related to theoretically clean observable  $A_T^2$  through  $S_3 = \frac{1}{2}(1 - F_L)A_T^2$ .



$A_I(K^{*0}\gamma)$ ,  $S_{K^{*0}\gamma}$ ,  $B(b \rightarrow s\gamma)$

S. Descotes-Genon et. al. [[arXiv:1104.334](https://arxiv.org/abs/1104.334)]



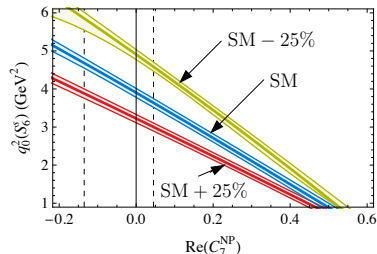
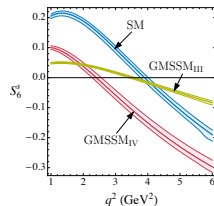
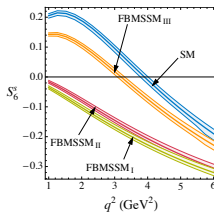
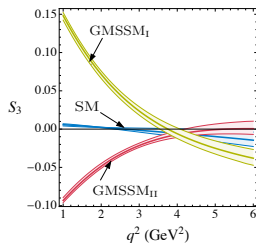
Non-SM like region

SM-like region

# Sensitivity to NP through $A_{FB}$ and $S_3$

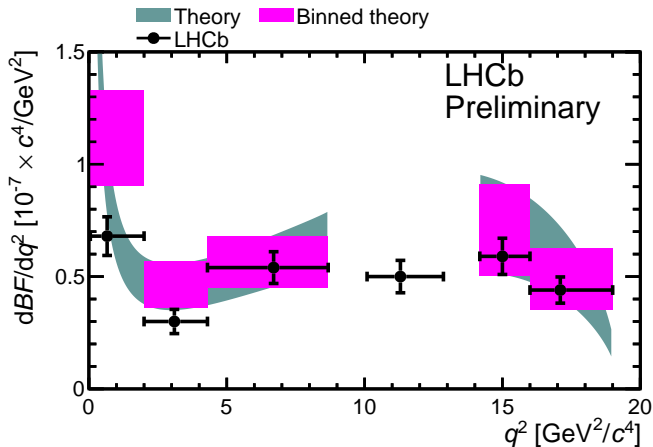
- Can be highly sensitive to NP contributions to  $C_7^{(\prime)}$ ,  $C_9^{(\prime)}$  and  $C_{10}^{(\prime)}$ .
- e.g. W. Almannshofer et. al. [[arXiv:0801.1214v5](https://arxiv.org/abs/0801.1214v5)] , where  $S_6 = -\frac{4}{3}A_{FB}$ .

Generic MSSM  
Flavour blind MSSM



SM  $\pm 25\%$  refers to  
a 25% change in  $C_9$ .

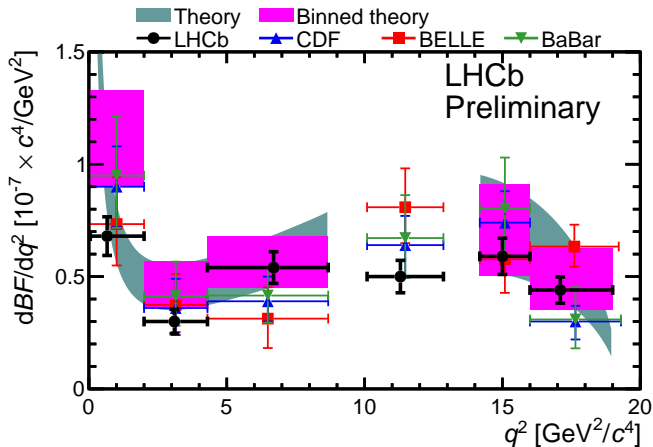
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ differential branching fraction



Theory prediction from C. Bobeth et al. [[arXiv:1105.0376](#)] (and references therein)



# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ differential branching fraction

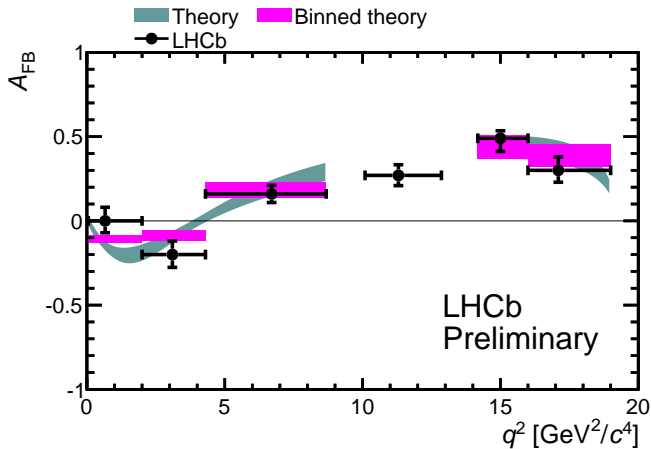


CDF, PRL 108 (2012)

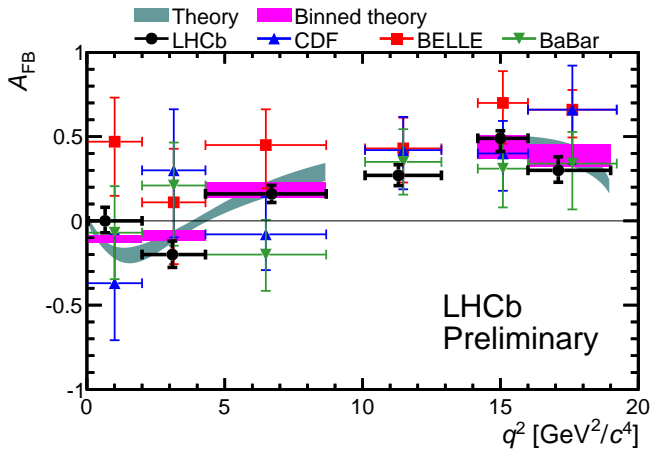
Belle, PRL 103 (2009)

BaBar prelim., Lake Louise 2012

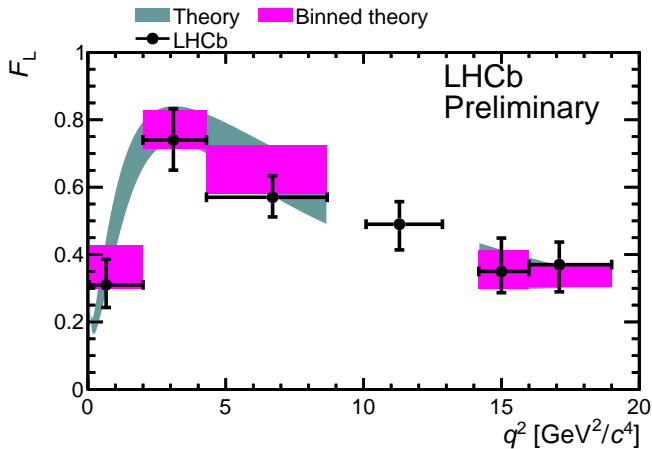
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ forward-backward asymmetry



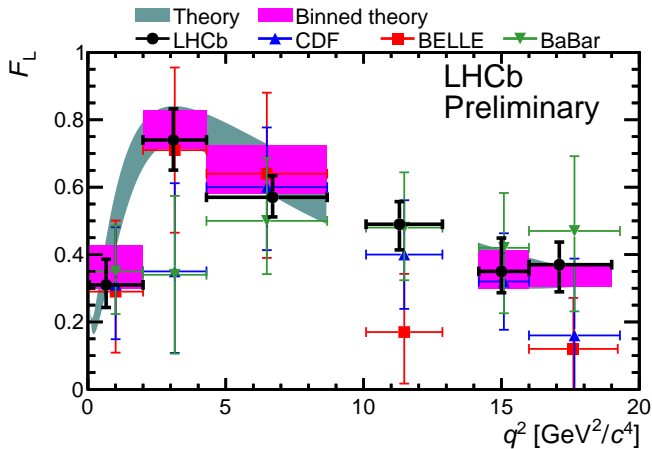
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ forward-backward asymmetry



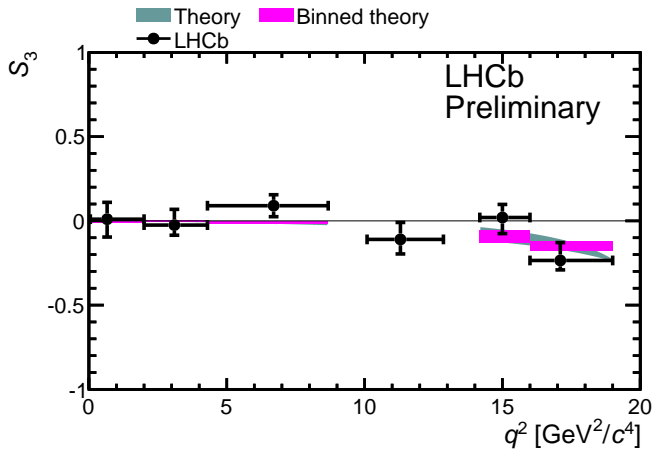
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ fraction of longitudinal polarisation of the $K^{*0}$



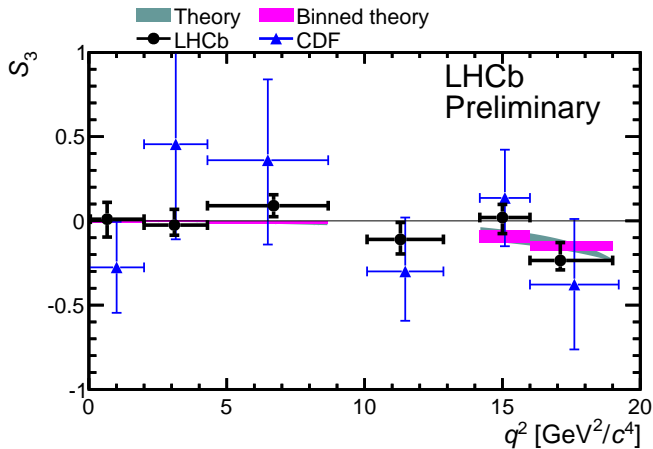
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ fraction of longitudinal polarisation of the $K^{*0}$



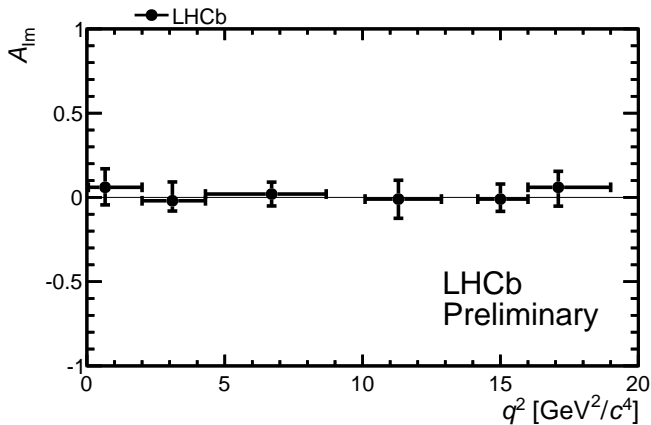
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ transverse asymmetry



# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ transverse asymmetry



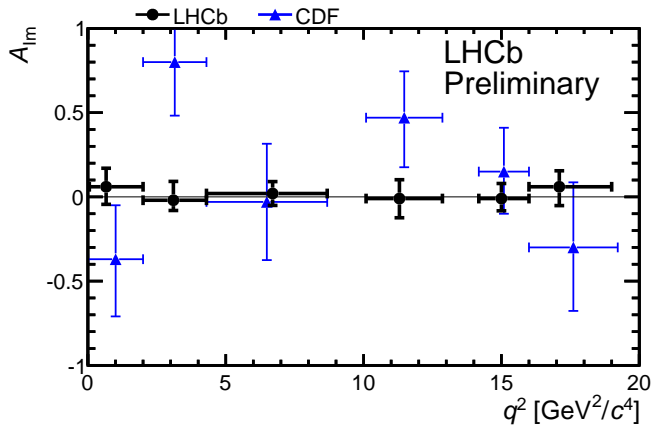
# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ T-odd asymmetry



$A_{lm}$  expected to be  $\mathcal{O}(10^{-3})$  in the SM



# $B_d \rightarrow K^{*0} \mu^+ \mu^-$ T-odd asymmetry



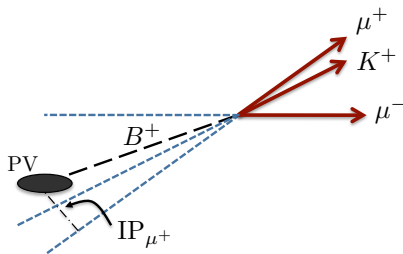
# Why don't we go to the kinematic limit of $q^2 = (m_B - m_X)^2$ ?

- In our  $B \rightarrow X \mu^+ \mu^-$  analyses we typically limit the range of  $q^2$  to be less than the kinematically allowed region (e.g.  $B_d \rightarrow K^{*0} \mu^+ \mu^-$ ,  $0.05 < q^2 < 19.0 \text{ GeV}^2/c^4$ ). Why?

- Using  $B^+ \rightarrow K^+ \mu^+ \mu^-$  as an example, at the end-point of the  $q^2$  spectrum,  $q^2 = (m_B - m_{K^+})^2$ , the  $K^+$  ends up being soft and collinear with the  $B_d$ .

→  $K^+$  has little or no impact parameter.

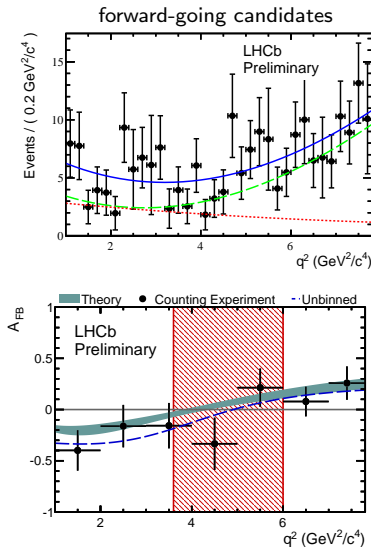
- Experimentally it can be difficult to understand the efficiency at this point.



# Zero-crossing point of the forward-backward asymmetry

- In the SM,  $A_{FB}$  varies with  $q^2$  and changes sign at a well defined point where leading uncertainties from the  $B \rightarrow K^{*0}$  form-factors cancel.
- Estimate zero-crossing point by fitting forward- and backward-going events separately.
- Gives:  $q_0^2 = 4.9^{+1.3}_{-1.1} \text{ GeV}^2/c^4$   
(LHCb preliminary)

c.f. SM predictions in the range  
 $3.9 - 4.3 \text{ GeV}^2/c^4$ .



- On-going analyses for:

- $B^+ \rightarrow K^+\mu^+\mu^-$  angular analysis ( $F_H$  and  $A_{FB}$ ) and differential branching fraction (with  $\sim 1000$  candidates).
- $B_d \rightarrow K_S^0\mu^+\mu^-$  differential branching fraction.
- $B^+ \rightarrow K^{*+}\mu^+\mu^-$  differential branching fraction.
- $\Lambda_b \rightarrow \Lambda^{(*)}\mu^+\mu^-$  differential branching fractions.
- $\mathcal{CP}$  asymmetries.
- $B_d \rightarrow K^{*0}e^-e^-$  at low- $q^2$ .

and many more ...

$b \rightarrow d\ell^+\ell^-$  processes (?)

# What can we do with $b \rightarrow d\mu^+\mu^-$ processes?

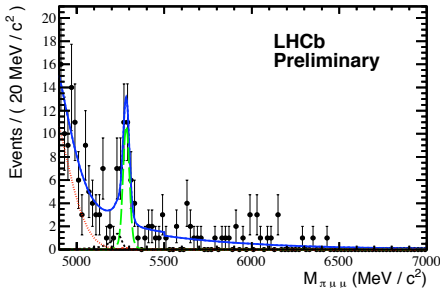
- Could also envisage models with FCNC outside the CKM structure:  
i.e.  $b \rightarrow d$  versus  $b \rightarrow s$  not suppressed by  $|V_{td}/V_{ts}|^2$ .
- With a  $\gtrsim 1 \text{ fb}^{-1}$  data set LHCb can have sensitivity to some of the Cabibbo suppressed modes.

- Very rare decay with SM branching fraction of  $(2 \pm 0.2) \times 10^{-8}$ .  
Hai-Zhen et al.  
[Commun. Theor. Phys. 50].
- Previous best limit from Belle:  
 $\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-) < 6.9 \times 10^{-8}$   
[PRD 78 (2008)].
- LHCb observes  $25.3^{+6.7}_{-6.4}$  candidates in  $1 \text{ fb}^{-1}$ .

$$\mathcal{B}(B^+ \rightarrow \pi^+ \mu^+ \mu^-) = (2.4 \pm 0.6 \pm 0.2) \times 10^{-8}$$

LHCb Preliminary

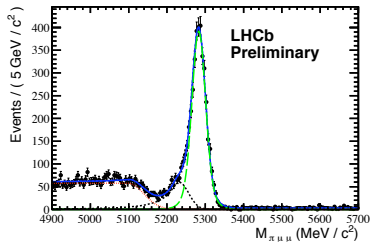
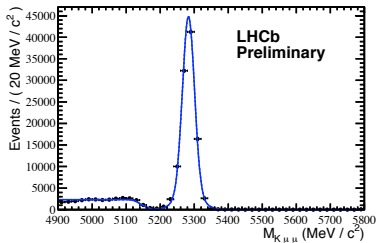
→ Observation of  $B^+ \rightarrow \pi^+ \mu^+ \mu^-$  at  $5.2\sigma$  (rarest observed  $B$  decay).



# $B^+ \rightarrow \pi^+ \mu^+ \mu^-$ (a few details)

- Branching fraction is normalised with respect to high statistics  $B^+ \rightarrow J/\psi K^+$  sample.
- Multivariate selection is used to reject combinatorial background.
- Main challenges are to remove combinatorial background and background from  $B^+ \rightarrow K^+ \mu^+ \mu^-$  with  $K \rightarrow \pi$  mis-id and to account for partially reconstructed decays.

NB Can use  $B^+ \rightarrow J/\psi \pi^+$  to understand the fitting and take shapes where possible from the data.



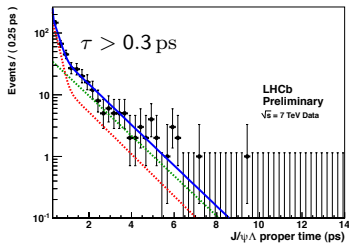
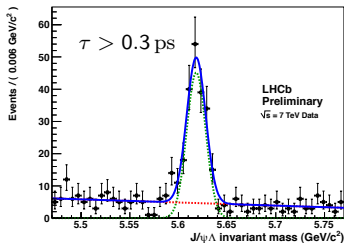
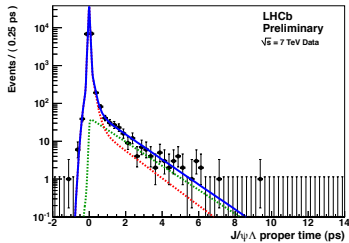
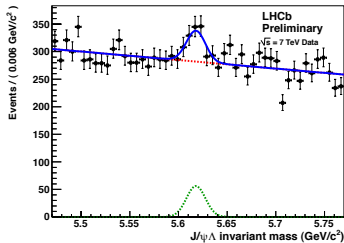


# Meson and Baryon lifetimes

- Preliminary measurement have been made with  $36 \text{ pb}^{-1}$  of integrated luminosity, using lifetime unbiased selection and trigger.

**NB**  $B_s$  lifetime measured with  $B_s \rightarrow J/\psi \phi$  candidates using a single exponential fit (i.e. ignoring  $\Delta\Gamma_s$ ).

# Example: $\Lambda_b$ lifetime



- In  $36 \text{ pb}^{-1}$  LHCb measures (Preliminary)

$$\tau(B^+ \rightarrow J/\psi K^+) = 1.689 \pm 0.022 \pm 0.047 \text{ ps},$$

$$\tau(B^0 \rightarrow J/\psi K^{*0}) = 1.512 \pm 0.032 \pm 0.042 \text{ ps},$$

$$\tau(B^0 \rightarrow J/\psi K_S^0) = 1.558 \pm 0.056 \pm 0.022 \text{ ps},$$

$$\tau^{\text{single}}(B_s^0 \rightarrow J/\psi \phi) = 1.447 \pm 0.064 \pm 0.056 \text{ ps},$$

$$\tau(\Lambda_b \rightarrow J/\psi \Lambda) = 1.353 \pm 0.108 \pm 0.035 \text{ ps}.$$

which are compatible with but not quite competitive with world averages, e.g.  $\tau_{\Lambda_b^0} = 1.391^{+0.038}_{-0.037}$  (PDG).

- Systematic uncertainty is dominated by understanding of the event reconstruction efficiency with proper time. Observe a reduced efficiency for tracking with high-IP tracks.

$\phi_s$ ,  $\Delta\Gamma_s$  and  $\Delta m_s$

- Measure  $\Delta m_s$  in flavour-tagged time-dependent analysis of  $B_s \rightarrow D_s \pi$  events, where  $D_s \rightarrow \phi \pi$  and  $D_s \rightarrow K^+ K^- \pi$ .

$$\mathcal{P}_t(t, q | \sigma_t, \eta) \propto \left\{ \Gamma_s e^{-\Gamma_s t} \frac{1}{2} \left[ \cosh \left( \frac{\Delta \Gamma_s}{2} t \right) + q [1 - 2\omega(\eta)] \cos(\Delta m_s t) \right] \theta(t) \right\} \\ \otimes G(t, S_{\sigma_t} \sigma_t) \epsilon(t) \epsilon_s.$$

- Effective tagging efficiency,

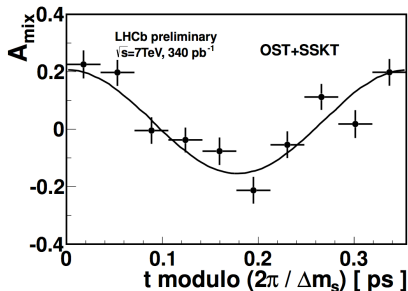
$$\epsilon_{\text{eff}, SS} = 1.3 \pm 0.4\% \text{ and}$$

$$\epsilon_{\text{eff}, OS} = 2.1 \pm 0.2\%.$$

- In  $0.34 \text{ fb}^{-1}$ :

$$\Delta m_s = 17.725 \pm 0.041 \pm 0.026 \text{ ps}^{-1} \\ (\text{LHCb Preliminary})$$

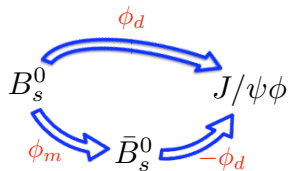
- Dominant systematic is from z-scale of the detector.



- Interference between mixing and decay gives rise to a  $\mathcal{CP}$  phase:  $\phi_s = \phi_m - 2\phi_d$ .

$$\phi_s^{\text{SM}} = -2\arg\left(\frac{V_{ts}V_{tb}^*}{V_{cs}V_{cb}^*}\right) = 0.036 \pm 0.002 \text{ rad}$$

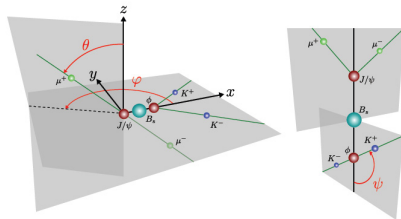
Charles et al. [[Phys. Rev. D84 \(2011\) 033005](#)]



# Measuring $\phi_s$ with $B_s \rightarrow J/\psi \phi$ and $B_s \rightarrow J/\psi f_0$

$B_s \rightarrow J/\psi \phi$ :

- Mixture of  $\mathcal{CP}$ -odd and  $\mathcal{CP}$ -even final state.
- Need to perform a time dependent angular analysis to separate  $\mathcal{CP}$  states and measure  $\phi_s$ .



$$S(\vec{\lambda}, t, \vec{\Omega}) = \underbrace{\epsilon(t, \vec{\Omega})}_{\text{acceptance}} \times \left( \underbrace{\frac{1+qD}{2} s(\vec{\lambda}, t, \vec{\Omega}) + \frac{1-qD}{2} \bar{s}(\vec{\lambda}, t, \vec{\Omega})}_{\text{flavour tagging}} \right) \otimes \underbrace{R_t}_{\text{time resolution}}$$

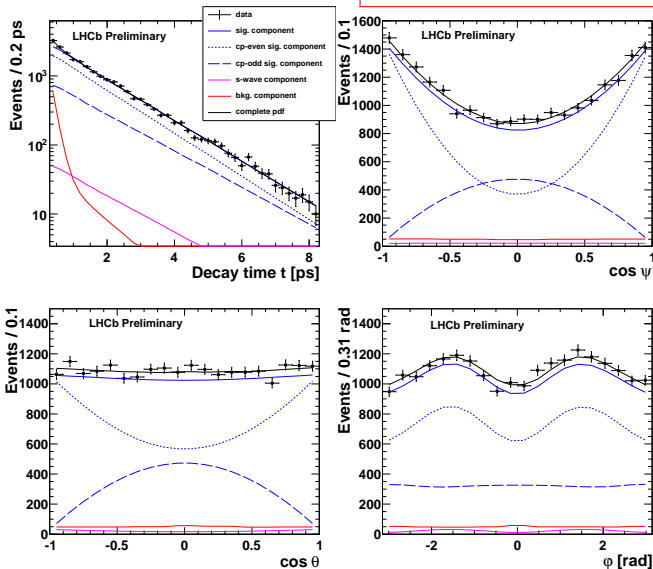
$$\vec{\lambda} = (\Gamma_s, \Delta\Gamma_s, \Delta m_s, \phi_s, |A_0|^2, |A_\perp|^2, \delta_\parallel, \delta_\perp, |A_S|^2, \delta_S)$$

$B_s \rightarrow J/\psi \pi^+ \pi^-$ :

- Region of  $\pi^+ \pi^-$  mass around the  $f_0$  gives a  $\mathcal{CP}$ -odd final state. Can measure  $\phi_s$  from fits to the  $B_s$  and  $\bar{B}_s$  lifetime.



21k signal candidates with  $\tau > 0.3$  ps



- $B_s \rightarrow J/\psi \phi$ :

$$\phi_s = 0.00 \pm 0.10 \pm 0.03 \text{ rad}$$

$$\Delta\Gamma_s = 0.12 \pm 0.02 \pm 0.01 \text{ ps}^{-1}$$

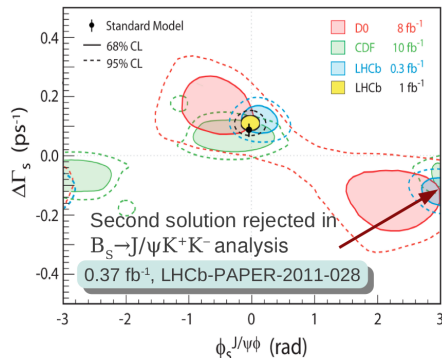
(Preliminary)

[LHCb-CONF-2012-002]

- $B_s \rightarrow J/\psi \pi^+ \pi^-$ :

$$\phi_s = -0.02 \pm 0.17 \pm 0.02 \text{ rad}$$

[LHCb-PAPER-2012-006]



Combination  $\phi_s = 0.00 \pm 0.08 \pm 0.03 \text{ rad}$  (LHCb Preliminary)

- Two-fold ambiguity  $\Delta\Gamma_s \rightarrow -\Delta\Gamma_s$  and  $\phi_s \rightarrow \phi_s + \pi$  resolved using  $B_s \rightarrow J/\psi K^+ K^-$  and  $S$ -wave interference. [LHCb-PAPER-2011-028]

- Can also constrain  $\Delta\Gamma_s$  and  $\phi_s$  using effective lifetime measurements of:

- $B_s \rightarrow J/\psi f_0$  (CP-odd)
- $B_s \rightarrow K^+ K^-$  (CP-even)

with an untagged analysis.

$$\tau_{K^+ K^-} = 1.44 \pm 0.10 \pm 0.01 \text{ ps}$$

LHCb, PLB 707 (2012)

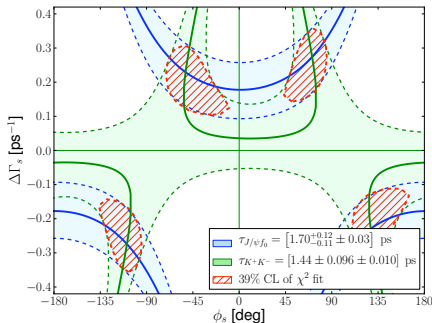
$$\tau_{J/\psi f_0} = 1.70^{+0.12}_{-0.11} \pm 0.03 \text{ ps}$$

CDF, PRD 84 (2011)

- LHCb update for  $B_s \rightarrow K^+ K^-$  with  $1 \text{ fb}^{-1}$ , using a lifetime unbiased trigger and offline selection:

$$1.468 \pm 0.046(\text{stat}) \text{ ps}$$

R. Fleischer et. al. [arXiv:1109.5115]



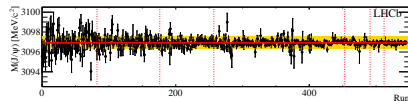
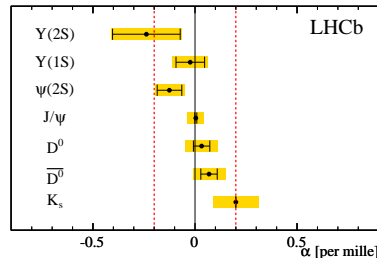
LHCb Preliminary

[LHCb-CONF-2012-001]

# Meson and Baryon masses

- LHCb profits from large tracking volume + large integral field.
- With a dedicated alignment programme can reach  $2 \times 10^{-4}$  on momentum scale (calibrated using  $J/\psi \rightarrow \mu^+ \mu^-$  decays).

**NB** Momentum scale becomes dominant systematic for many high mass states (can also look at decays with small q-value).



- With  $35 \text{ fb}^{-1}$  measure:

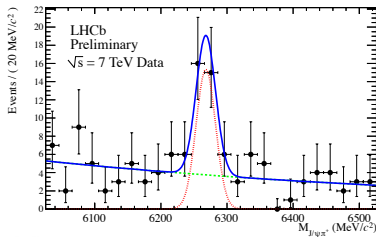
Quantity	LHCb measurement	Best previous measurement	PDG fit
$M(B^+)$	$5279.38 \pm 0.35$	$5279.10 \pm 0.55$ [4]	$5279.17 \pm 0.29$
$M(B^0)$	$5279.58 \pm 0.32$	$5279.63 \pm 0.62$ [4]	$5279.50 \pm 0.30$
$M(B_s^0)$	$5366.90 \pm 0.36$	$5366.01 \pm 0.80$ [4]	$5366.3 \pm 0.6$
$M(\Lambda_b^0)$	$5619.19 \pm 0.76$	$5619.7 \pm 1.7$ [4]	—
$M(B^0) - M(B^+)$	$0.20 \pm 0.20$	$0.33 \pm 0.06$ [15]	$0.33 \pm 0.06$
$M(B_s^0) - M(B^+)$	$87.52 \pm 0.32$	—	—
$M(\Lambda_b^0) - M(B^+)$	$339.81 \pm 0.72$	—	—

[4] CDF, [PRL 96 (2006)] , [15] BaBar, [PRD 78 (2008)]

- World best measurement of  $B^+$ ,  $B_d$ ,  $B_s$  and  $\Lambda_b^0$  masses.

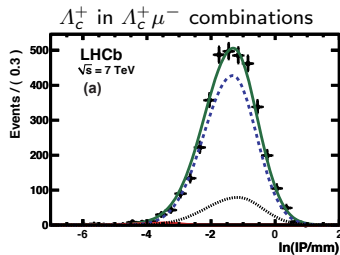
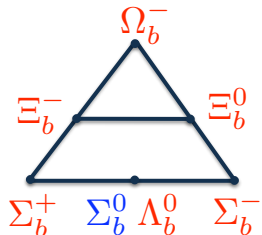
- $B_c^+$  mass measured using  $B_c^+ \rightarrow J/\psi \pi^+$  decays in  $35 \text{ pb}^{-1}$  of integrated luminosity.
- Preliminary result:  
 $M(B_c^+) = 6268.0 \pm 4.0 \pm 0.6 \text{ MeV}/c^2$

c.f. PDG (2010) average of  
 $6277 \pm 6 \text{ MeV}/c^2$

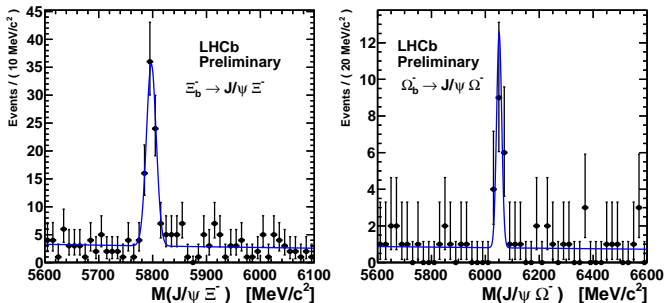


# Beautiful baryons

- Seven ground state b-baryons:
  - $\Sigma_b^0$  unobserved.
  - $\Xi_b^0$  recently observed at CDF.
- Large b-Baryon production at the LHC:  
 $\sigma(pp \rightarrow b\bar{b}X) = 75.3 \pm 5.4 \pm 13.0 \mu\text{b}$   
in  $2 < \eta < 6$  [LHCb-PAPER-2010-002].
- Estimate  $f_\Lambda/(f_u + f_d)$  using semileptonic decays ( $\Lambda_c^+ \mu^-$  vs  $D^0 \mu^-$  and  $D^+ \mu^-$ ) to be  $\sim 0.4$  at ( $p_T = 0$ ) and  $\sim 0.25$  ( $p_T = 10 \text{ GeV}/c$ ).  
[LHCb-PAPER-2011-018]







- Using  $0.62 \text{ fb}^{-1}$  measure (LHCb Preliminary):

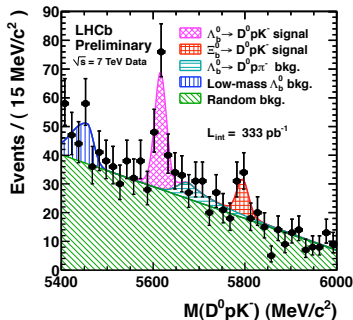
$$M_{\Xi_b^-} = 5796.5 \pm 1.2 \pm 1.2 \text{ MeV}/c^2$$

$$M_{\Omega_b^-} = 6050.3 \pm 4.5 \pm 2.2 \text{ MeV}/c^2$$

$$M_{\Omega_b^-} - M_{\Xi_b^-} = 253.9 \pm 4.8 \pm 1.2 \text{ MeV}/c^2$$

Systematic uncertainty is dominated by the momentum scale.

- First observation by CDF, [PRL 107 (2011)].
- Search in the channel  $\Xi_b^0 \rightarrow D^0 p K^-$ , with  $0.33 \text{ fb}^{-1}$ .
- Observe hint of a signal with  $2.6\sigma$  significance.



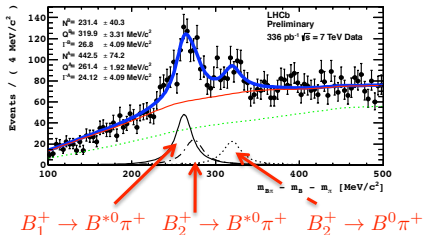
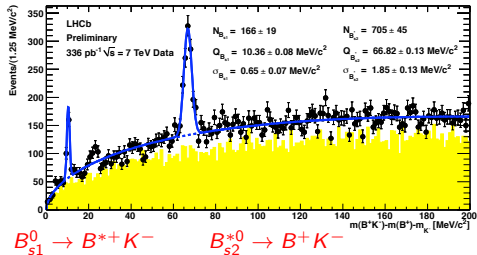
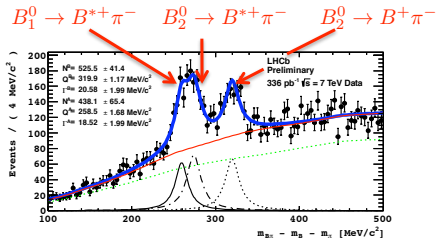
$$\Delta M_{\Xi_b^0} = M_{\Xi_b^0} - M_{\Lambda_b^0} = 181.8 \pm 5.5 \pm 0.5 \text{ MeV}/c^2$$

$$M_{\Xi_b^0} = 5802.0 \pm 5.5 \pm 1.7 \text{ MeV}/c^2$$

$$M_{\Xi_b^0} = 5787.8 \pm 5.0 \pm 1.3 \text{ MeV}/c^2$$

(LHCb Preliminary)

- Perform search for  $B_{(s)}^{**}$  states in  $B^{\pm}K^{\mp}$ ,  $B^{\pm}\pi^{\mp}$  and  $B^0\pi^{\pm}$  channels with  $0.33\text{ fb}^{-1}$ .
- Photons from  $B^*$  are not reconstructed  $\rightarrow$  shifted  $Q$ -values.



- No direct determination of quantum numbers. Matching to expected states from HQET.
- LHCb Preliminary result:

$$\begin{aligned}
 M_{B_{s1}^0} &= (5828.99 \pm 0.08_{\text{stat}} \pm 0.13_{\text{syst}} \pm 0.45_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_{s1}^0} &= (10.36 \pm 0.08_{\text{stat}} \pm 0.13_{\text{syst}}) \text{ MeV}/c^2, \\
 M_{B_{s2}^{*0}} &= (5839.67 \pm 0.13_{\text{stat}} \pm 0.17_{\text{syst}} \pm 0.29_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_{s2}^{*0}} &= (66.82 \pm 0.13_{\text{stat}} \pm 0.17_{\text{syst}}) \text{ MeV}/c^2, \\
 M_{B_1^0} &= (5724.1 \pm 1.7_{\text{stat}} \pm 2.0_{\text{syst}} \pm 0.5_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_1^0} &= (259.6 \pm 1.7_{\text{stat}} \pm 2.0_{\text{syst}}) \text{ MeV}/c^2, \\
 M_{B_1^+} &= (5726.3 \pm 1.9_{\text{stat}} \pm 3.0_{\text{syst}} \pm 0.5_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_1^+} &= (261.4 \pm 1.9_{\text{stat}} \pm 3.0_{\text{syst}}) \text{ MeV}/c^2, \\
 M_{B_2^{*0}} &= (5738.6 \pm 1.2_{\text{stat}} \pm 1.2_{\text{syst}} \pm 0.3_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_2^{*0}} &= (319.9 \pm 1.2_{\text{stat}} \pm 1.3_{\text{syst}}) \text{ MeV}/c^2, \\
 M_{B_2^{*+}} &= (5739.0 \pm 3.3_{\text{stat}} \pm 1.6_{\text{syst}} \pm 0.3_{\text{syst}}^{B \text{ mass}}) \text{ MeV}/c^2, & Q_{B_2^{*+}} &= (319.9 \pm 3.3_{\text{stat}} \pm 1.6_{\text{syst}}) \text{ MeV}/c^2,
 \end{aligned}$$

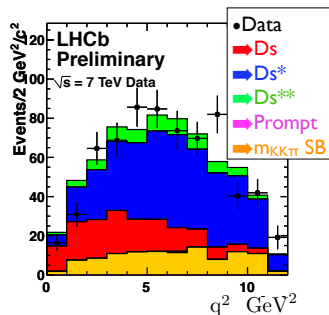
- The masses of  $B_1^+$  and  $B_2^{*+}$  are measured for the first time.

NB Masses of Isospin partners are compatible.

# Semileptonic $b$ -decays

# Semileptonic $b$ -decays at LHCb

- Semileptonic  $B_d$ ,  $B^+$ ,  $B_s$  and  $\Lambda_b$  decays have been extensively used to understand  $b\bar{b}$  production and  $B_d/B^+/B_s/\Lambda_b$  hadronisation fractions in LHCb.
- Ongoing programme to measure:
  1. Exclusive Cabibbo favoured and suppressed decays (including  $B_s$  and  $\Lambda_b$ ), probing form-factors and CKM parameters:
    - e.g.  $V_{ub}$  from  $B_s \rightarrow K\mu\nu$ .
  2. Semileptonic asymmetry,  $A_{sl}$ .
- Can perform neutrino 'reconstruction' (with two-fold using ambiguity) using  $B$  pointing constraints to separate different semileptonic contributions.



LHCb  $c(\text{harm})$

- LHCb is not just a place to study  $B$  decays / production.
- Also have prolific charm production:

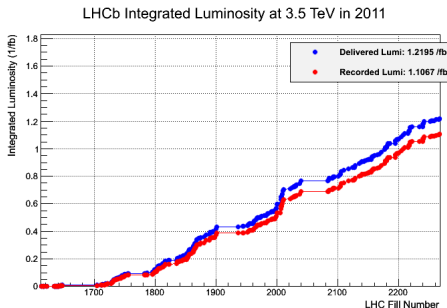
$$\sigma_{c\bar{c}} \sim 20 \times \sigma_{b\bar{b}}.$$

→ Charm mixing and CP violation ( $y_{\mathcal{CP}}$ ,  $A_\Gamma$ ,  $A_{\mathcal{CP}}$ ), charm rare decays, charm semileptonic decays, charm masses and  $D_{(s)}^{**}$  spectroscopy.



# Outlook and Summary

- Current analyses are based on 2010 dataset ( $36 \text{ pb}^{-1}$ ), Summer 2011 dataset ( $0.33 \text{ fb}^{-1}$ ) or full 2011 dataset ( $1 \text{ fb}^{-1}$ ).
- Updates of many analyses to the full 2011 dataset are ongoing.
- In 2012 expect to record an integrated luminosity of  $\sim 1.5 \text{ fb}^{-1}$  at  $\sqrt{s} = 8 \text{ TeV}$ .



# Summary

Excellent performance from LHC and LHCb in 2011.

Large number of new heavy flavour results presented this year. More to come at the Spring & Summer conferences.

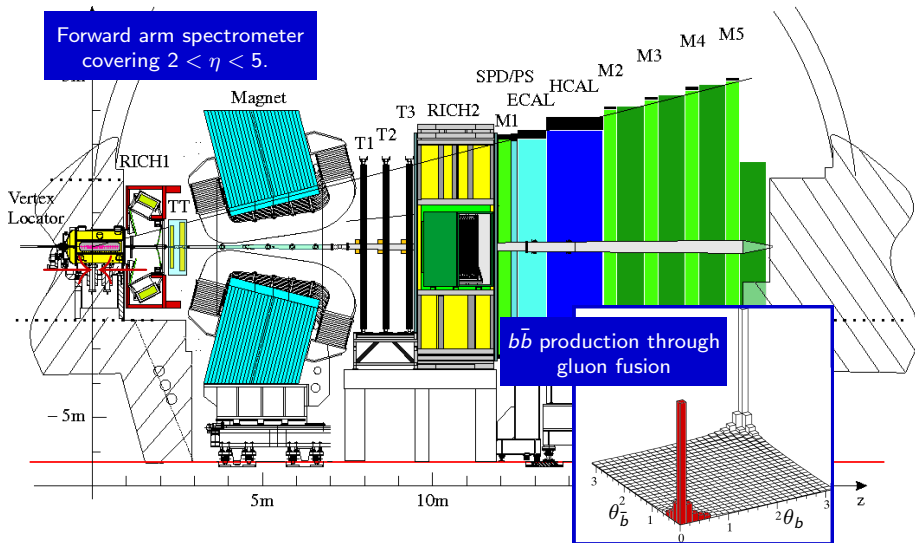
Feedback from the theory community is very welcome!





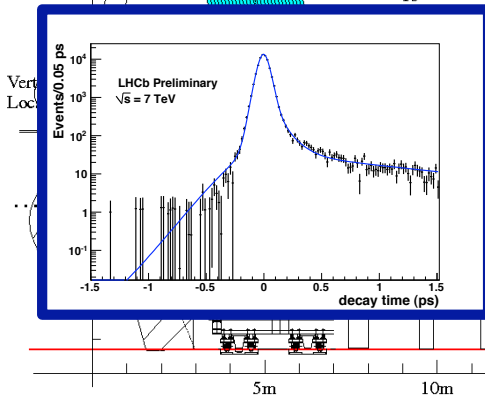
# BACKUP

# The LHCb detector

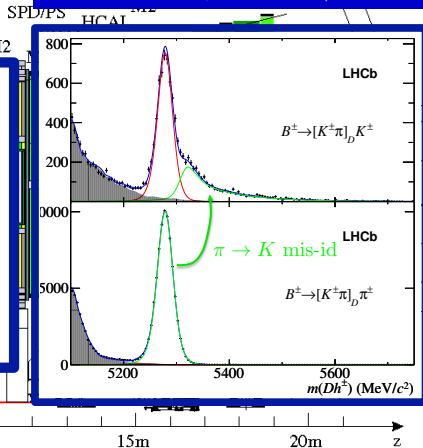


# The LHCb detector

$r - \phi$  vertex detector that approaches within 8 mm of the beam line + large boost  $\rightarrow \sigma_\tau \sim 45$  fs

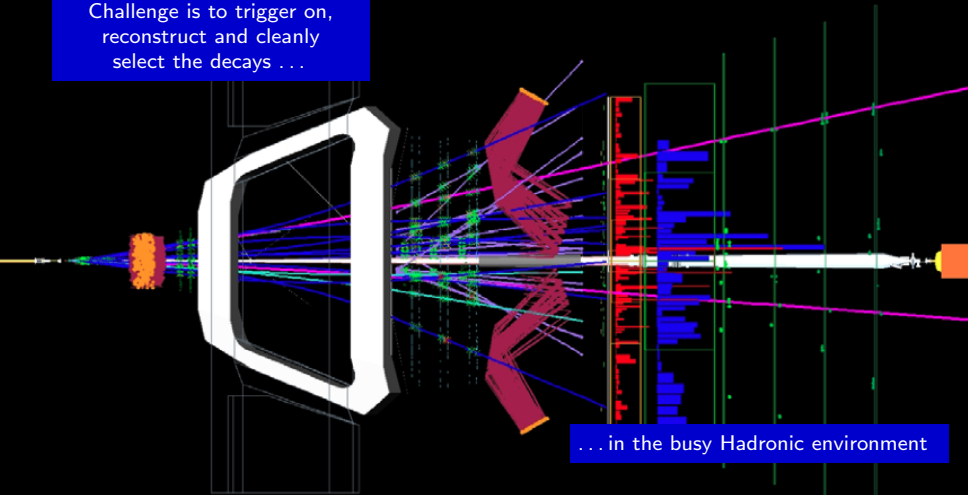


RICH detectors  $\rightarrow K/\pi/p$  separation from  $2 < p < 100$  GeV/c



# Needle in a haystack?

Challenge is to trigger on,  
reconstruct and cleanly  
select the decays ...



... in the busy Hadronic environment



# Exotics: X(3872) and X(4140)

X3872 [LHCb-PAPER-2011-034]:

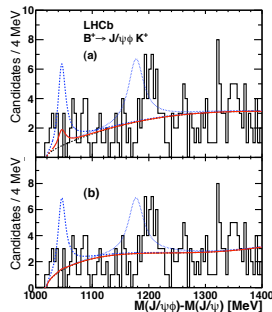
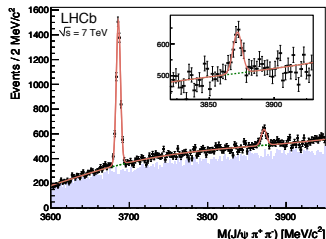
- Measure mass of prompt X(3872):

$$M_{X(3872)} = 3871.95 \pm 0.48 \pm 0.12 \text{ MeV}/c^2$$

- Not yet at the level of Belle [PRD 84 (2011)].

X4140 [LHCb-PAPER-2011-033]:

- Search for X(4140) in  $B^+ \rightarrow J/\psi \phi K^+$  decays in  $0.37 \text{ fb}^{-1}$ .
- Do not confirm excesses seen by CDF [arXiv:1101.6058].



$$B_d \rightarrow K^{*0} \gamma$$

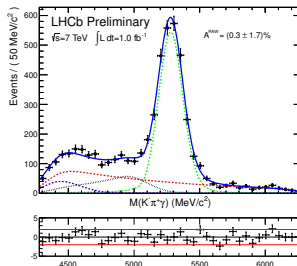
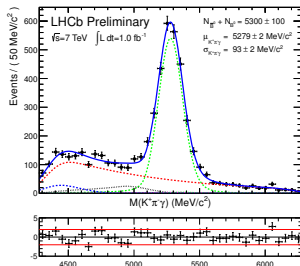
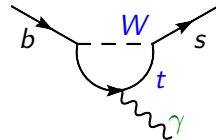
- SM prediction for CP asymmetry:

$$A_{CP} = -0.006 \pm 0.004 \text{ [arXiv:0406055]}$$

- Previous best measurement from *BABAR*:



$$A_{CP} = -0.016 \pm 0.022 \pm 0.007 \text{ [PRL 103 (2009)]}$$

$$\begin{aligned} A_{CP}(B_d \rightarrow K^{*0}\gamma) &= A_{CP}^{\text{RAW}}(B_d \rightarrow K^{*0}\gamma) - A_{\text{prod.}}(B_d) - A_{\text{det.}}(K\pi) \\ &= 0.008 \pm 0.017(\text{stat}) \pm 0.009(\text{syst}) \text{ [Preliminary]} \end{aligned}$$



- Also collecting large samples of  $B_s$  decays. Observe 240  $B_s \rightarrow \phi\gamma$  candidates in  $0.37 \text{ fb}^{-1}$  [LHCb-PAPER-2011-042].

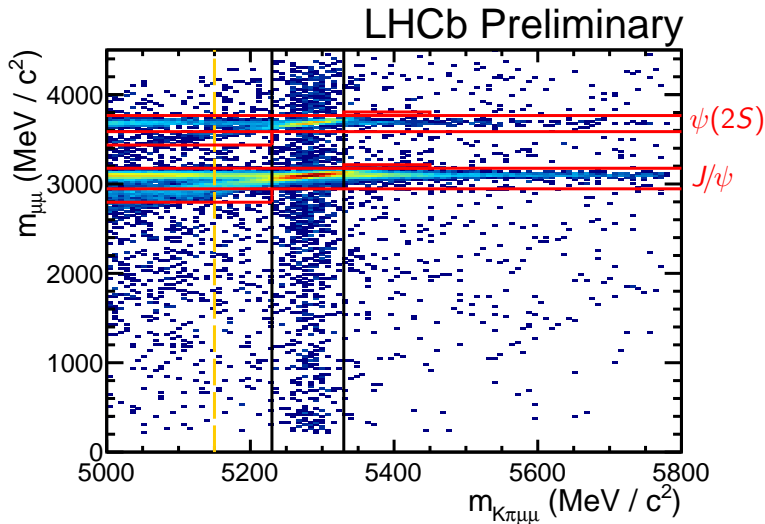
# Sensitivity to NP through rare decays

	Operator $\mathcal{O}_i$	$B \rightarrow K^{*0}\gamma$	$B \rightarrow K^{*0}\mu^+\mu^-$	$B \rightarrow \mu^+\mu^-$
	$\mathcal{O}_7 \sim m_b(\bar{s}_L\sigma_{\mu\nu}b_R)F_{\mu\nu}$	✓	✓	
	$\mathcal{O}_9 \sim (\bar{s}b)_{V-A}(\bar{\ell}\ell)_V$		✓	
	$\mathcal{O}_{10} \sim (\bar{s}b)_{V-A}(\bar{\ell}\ell)_A$		✓	✓
	$\mathcal{O}_{S,P} \sim (\bar{s}b)_{S+P}(\bar{\ell}\ell)_{S,P}$			✓

In the SM:

- $\mathcal{C}_{S,P} \propto m_\ell m_b / m_W^2 \sim 0$ .
- Helicity flipped operators ( $\mathcal{C}'_i \mathcal{O}'_i$ ) suppressed by  $m_s / m_b$ .

$B_d \rightarrow K^{*0} \mu^+ \mu^-$   $m_{\mu^+ \mu^-}$  versus  $m_{K^+ \pi^- \mu^+ \mu^-}$



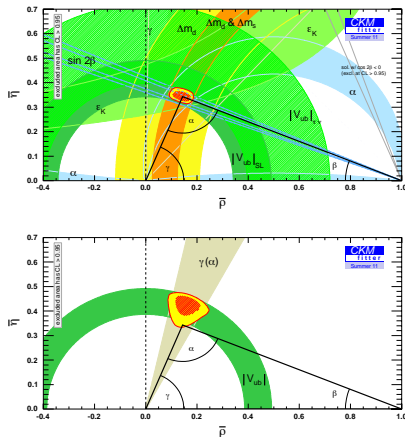
Partially reconstructed backgrounds

# Towards a measurement of $\gamma$

# CKM picture

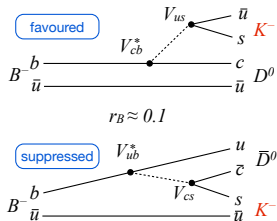
- CKM picture seems to describe nature remarkably well.

But Have just been discussing the possible effects from NP to loop-order processes. The picture from tree-level is much less complete ...

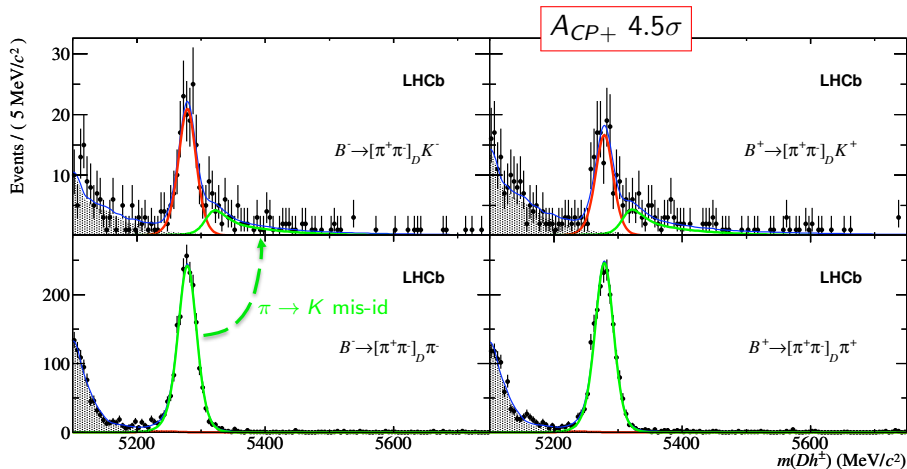


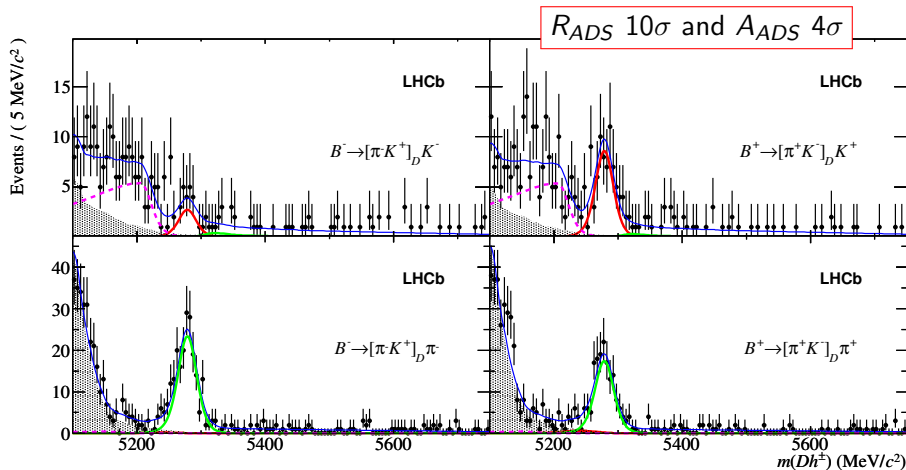
# How can we access $\gamma$

- Access CKM phase  $\gamma$  through interference of  $b \rightarrow u$  and  $b \rightarrow c$  transitions in decays with a common final state.
- One way is through  $B^\pm \rightarrow DK^\pm$  decays, where the  $D^0$  and  $\bar{D}^0$  decay to common final states:
  - $D^0, \bar{D}^0 \rightarrow \pi^+\pi^-$  or  $K^+K^-$  (GLW)  
see e.g. [PLB 265 (1991)]
  - $D^0, \bar{D}^0 \rightarrow K^+\pi^-$  (ADS)  
[PRL 78 (1997)]
- ADS mode is experimentally challenging:
  - It's a fully hadronic  $B$  decay with an effective branching fraction of  $2 \times 10^{-7}$ !









First observation of the ADS mode

- What do  $A_{ADS}$  and  $A_{CP+}$  tell us about  $\gamma$ ?

$$A_{CP+} = \frac{2r_B \sin \delta_B \sin \gamma}{1 + r_B^2 \pm 2r_B \cos \delta_B \cos \gamma} = 0.145 \pm 0.032 \pm 0.010,$$

$$A_{ADS} = \frac{2r_B r_D \sin(\delta_B + \delta_D) \sin \gamma}{r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B + \delta_D) \cos \gamma} = -0.52 \pm 0.15 \pm 0.02,$$

$$R_{CP+} = 1 + r_B^2 + 2r_B \cos \delta_B \cos \gamma = 1.007 \pm 0.038 \pm 0.012,$$

$$R_{ADS} = \frac{r_B^2 + r_D^2 + 2r_B r_D \cos(\delta_B + \delta_D) \cos \gamma}{1 + r_B^2 r_D^2 \cos(\delta_B - \delta_D) \cos \gamma} = 0.0152 \pm 0.0020 \pm 0.0004$$

(LHCb Preliminary)

- Combine with GGSZ modes ( $B^- \rightarrow DK^-$ ,  $D \rightarrow K_S^0 \pi^+ \pi^-$ ) to estimate  $\gamma$ .

# Anatomy of the $B_d \rightarrow K^{*0} \mu^+ \mu^-$ decay

- Anatomy of the decay was explained in detail by Christoph yesterday. Described by 12 angular terms:

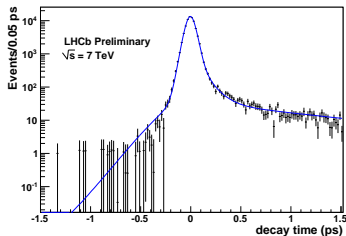
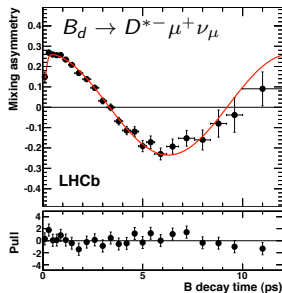
$$\frac{d^4\Gamma}{d\cos\theta_\ell d\cos\theta_K d\phi dq^2} \propto \left[ J_1^s + J_1^c + (J_2^s + J_2^c) \cos 2\theta_\ell + J_3 \sin^2 \theta_\ell \cos 2\phi + \right. \\ J_4 \sin 2\theta_\ell \cos \phi + J_5 \sin \theta_\ell \cos \phi + \\ J_6 \cos \theta_\ell + J_7 \sin \theta_\ell \sin \phi + J_8 \sin 2\theta_\ell \sin \phi + \\ \left. J_9 \sin^2 \theta_\ell \sin 2\phi \right]$$

- Some possibility to access  $S_4$ ,  $S_5$ ,  $S_7$  and  $S_8$  using other folding techniques. See A. Bharucha & W. Reece [[arXiv:1002.4310](https://arxiv.org/abs/1002.4310)].

# Opposite side flavour tagging and proper time resolution

- Opposite side tagging ( $\mu$ ,  $K$ ,  $e$  and vertex charge) calibrated using  $B_d \rightarrow D^{*-} \mu^+ \nu_\mu$ ,  $B^+ \rightarrow J/\psi K^+$  and  $B_d \rightarrow J/\psi K^{*0}$  events.
- Performance is typically  $\varepsilon \mathcal{D}^2 \sim 2 - 3\%$ .
- Work ongoing to improve tagging performance and to include same-sign tagging.
- Proper time resolution measured in data using prompt events ( $\tau = 0$ ):  
 $\sigma_\tau \sim 45 \text{ fs}$  (c.f.  $\sim 350 \text{ fs}$  period of  $B_s$ ).
- In practice per-event resolutions and mistag probabilities are used in time dependent + tagged analyses.

[LHCb-PAPER-2011-027]



$$\mathcal{P}_t(t, q | \sigma_t, \eta) \propto \left\{ \Gamma_s e^{-\Gamma_s t} \frac{1}{2} \left[ \cosh \left( \frac{\Delta \Gamma_s}{2} t \right) + q [1 - 2\omega(\eta)] \cos(\Delta m_s t) \right] \theta(t) \right\} \\ \otimes G(t, S_{\sigma_t} \sigma_t) \epsilon(t) \epsilon_s.$$

- Measure  $\Delta m_s$  in flavour-tagged time-dependent analysis of  $B_s \rightarrow D_s^+ \pi^+$  and  $B_s \rightarrow D_s^+ \pi^+ \pi^- \pi^+$  events, where  $D_s^+ \rightarrow \phi \pi^-$  and  $D_s^+ \rightarrow K^+ K^- \pi^-$ ,
- Effective tagging efficiency,  $\epsilon_{\text{eff}} = 3.8 \pm 2.1\%$ .
- In  $36 \text{ pb}^{-1}$ :

$$\Delta m_s = 17.63 \pm 0.11 \pm 0.02 \text{ ps}^{-1}$$

- Dominant systematic is from z-scale of the detector.

