

String Theory: Example Sheet 3

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1. Show that an infinitesimal $SL(2, \mathbb{C})$ transformation, acting as the Möbius group on $z \in \mathbb{C} \cup \infty$, takes the form

$$z \mapsto z' = \alpha + \beta z + \gamma z^2$$

where α , β and γ are constants. Hence show that the volume element of $SL(2, \mathbb{C})$ may be written as

$$dV_{SL(2, \mathbb{C})} = \frac{d^2 z_1 d^2 z_2 d^2 z_3}{|z_1 - z_2|^2 |z_2 - z_3|^2 |z_3 - z_1|^2}$$

in terms the coordinates z_1, z_2, z_3 of any three distinct points on the Riemann sphere.

2. The ghost field $c(z)$ has mode expansion $c(z) = \sum_{n \in \mathbb{Z}} c_n z^{-n+1}$. Show that the ghost vacuum $|0\rangle$ is defined by $c_n |0\rangle = 0$ for $n > 1$. Show that the bc CFT has correlation function

$$\langle c(z_1) c(z_2) c(z_3) \rangle_{bc} = K(z_1 - z_2)(z_2 - z_3)(z_3 - z_1)$$

for some constant K that you should find. [*You may assume that $\langle 0 | c_{-1} c_0 c_1 | 0 \rangle = 1$.] Hence show that*

$$\prod_{i=1}^n d^2 z_i \left\langle \prod_{i=1}^3 c(z_i) \bar{c}(\bar{z}_i) \right\rangle_{bc, \bar{b}\bar{c}} / dV_{SL(2, \mathbb{C})} = \prod_{i=4}^n d^2 z_i.$$

What is the significance of this result?

3. Show that the OPE between two plane wave operators $:e^{ik \cdot X}:$ takes the form

$$:e^{ik_1 \cdot X}:(z_1, \bar{z}_1) :e^{ik_2 \cdot X}:(z_2, \bar{z}_2) \sim |z_1 - z_2|^{\alpha' k_1 \cdot k_2} :e^{ik_1 \cdot X}(z_1, \bar{z}_1) e^{ik_2 \cdot X}(z_2, \bar{z}_2):.$$

[*You may find it useful to write the exponential operator as*

$$:e^{ik \cdot X}:(z, \bar{z}) = \sum_{n=0}^{\infty} \frac{i^n}{n!} k_{\mu_1} \cdots k_{\mu_n} :X^{\mu_1} \cdots X^{\mu_n}:(z, \bar{z})$$

and first consider the OPE between $X^\mu(z_1, \bar{z}_1)$ and $:e^{ik \cdot X}:(z_2, \bar{z}_2)$.]

4a. The scattering amplitude for m closed string tachyons is given by

$$\mathcal{A}^{(m)}(k_1, \dots, k_m) = \frac{g_s^{m-2}}{\text{Vol}(SL(2, \mathbb{C}))} \int \prod_{i=1}^m d^2 z_i \langle \hat{V}(z_1, k_1) \cdots \hat{V}(z_m, k_m) \rangle,$$

where $\hat{V}(z, k) = e^{ik \cdot X(z, \bar{z})}$ and the correlation function is computed using the gauge fixed free Polyakov action

$$S_{\text{Poly}} = \frac{1}{2\pi\alpha'} \int \partial X^\mu \bar{\partial} X_\mu d^2 z.$$

By expressing the correlation function as a Gaussian integral, show that the amplitude is

$$\mathcal{A}^{(m)} \sim \frac{g_s^{m-2}}{\text{Vol}(\text{SL}(2, \mathbb{C}))} \delta^{26} \left(\sum_i k_i \right) \int \prod_{i=1}^m d^2 z_i \prod_{j<l} |z_j - z_l|^{\alpha' k_j \cdot k_l}.$$

b. Show that this integral is invariant under the worldsheet $\text{SL}(2, \mathbb{C})$ transformation

$$z_i \rightarrow \frac{az_i + b}{cz_i + d}$$

only when the momenta are on-shell, *i.e.* $k_i^2 = 4/\alpha'$ for each i . Hence explain why the 4-tachyon amplitude can be reduced to the integral

$$\mathcal{A}^{(4)} \sim g_s^2 \delta^{26} \left(\sum_i k_i \right) \int d^2 z |z|^{\alpha' k_2 \cdot k_3} |1 - z|^{\alpha' k_3 \cdot k_4} \quad (\dagger)$$

c. By evaluating the remaining integral in terms of Γ -functions, show that string theory reproduces the Virasoro-Shapiro amplitude

$$\mathcal{A}^{(4)} \sim g_s^2 \delta^{26} \left(\sum_i k_i \right) \frac{\Gamma(-1 - \alpha' s/4) \Gamma(-1 - \alpha' t/4) \Gamma(-1 - \alpha' u/4)}{\Gamma(2 + \alpha' s/4) \Gamma(2 + \alpha' t/4) \Gamma(2 + \alpha' u/4)},$$

where s, t, u are the usual Mandelstam variables.

5. Explain why the limit $s \rightarrow \infty$, with t fixed corresponds to small angle scattering at high energy. Show that in this limit the Virasoro-Shapiro amplitude exhibits so-called Regge behaviour,

$$\mathcal{A}^{(4)} \rightarrow g_s^2 \delta^{26} \left(\sum_i p_i \right) \frac{\Gamma(-1 - \alpha' t/4)}{\Gamma(2 + \alpha' t/4)} s^{2 + \alpha' t/2}.$$

Show that this result agrees with the saddle point approximation to the integral (\dagger) above.

6a. Write down the vertex operator for a massless closed string state with polarization $\zeta_{\mu\nu}$ and momentum k_μ . What are the restrictions on k_μ and $\zeta_{\mu\nu}$?

b. Consider the scattering of a massless closed string mode with momentum k_1 and two tachyons with momenta k_2 and k_3 . Show that $k_1 \cdot k_2 = k_1 \cdot k_3 = 0$, while $k_2 \cdot k_3 = -4/\alpha'$.

c. Show that the 3-point scattering amplitude for these particles is given by

$$\mathcal{A}^{(3)} \sim \frac{g_s}{\text{Vol}(\text{SL}(2, \mathbb{C}))} \delta^{26} \left(\sum_i k_i \right) \int \prod_{i=1}^3 d^2 z_i \frac{1}{|z_{23}|^4} \zeta_{\mu\nu} \left(\frac{k_2^\mu}{z_{12}} + \frac{k_3^\mu}{z_{13}} \right) \left(\frac{k_2^\nu}{\bar{z}_{12}} + \frac{k_3^\nu}{\bar{z}_{13}} \right)$$

where $z_{ij} = z_i - z_j$.

d. Show that the residual $\text{SL}(2; \mathbb{C})$ gauge redundancy allows us to simplify this to

$$\mathcal{A}^{(3)} \sim g_s \delta^{(26)} \left(\sum_i k_i \right) \zeta_{\mu\nu} (k_2^\mu - k_3^\mu) (k_2^\nu - k_3^\nu). \quad (\ddagger)$$

e. Consider the 26-dimensional theory

$$S[G, \Phi] = -\frac{1}{2\kappa^2} \int \left(R(G) + \frac{1}{2} G^{\mu\nu} \partial_\mu \Phi \partial_\nu \Phi + \frac{1}{2} M^2 \Phi^2 \right) \sqrt{-G} d^{26}x$$

of gravity coupled to a real scalar field Φ of mass $M^2 = -4/\alpha'$. Show that the tree-level amplitude for a single graviton of polarization $\zeta_{\mu\nu}$ and momentum k_1 to be emitted from a scalar of momentum k_2 agrees with the result (‡). Why does this amplitude not receive stringy corrections?