## Example Sheet 3

## 1. Inertia tensor

(a) Prove that the principal moments of inertia, $\left(I_{1}, I_{2}, I_{3}\right)$, are real and non-negative.
(b) During the lectures we outlined the proof of the Parallel Axis Theorem, which is a statement that the inertia tensor about a point $P$, which is displaced by $\mathbf{c}$ from the centre of mass $C$, is related to the inertia tensor about $C$ by

$$
I_{a b}^{P}=I_{a b}^{C}+M\left(c^{2} \delta_{a b}-c_{a} c_{b}\right),
$$

where $M$ is the total mass of the body. Complete the proof of the theorem. (It will be helpful to choose the origin to be at the centre of mass.)

## 2. Euler angles

The rotation matrix that relates the body axes $\left(\mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}\right)$ to the space axes $\left(\mathbf{e}_{x}, \mathbf{e}_{y}, \mathbf{e}_{z}\right)$ is given in terms of the Euler angles $(\theta(t), \phi(t), \psi(t))$ by

$$
\begin{aligned}
\mathrm{R} & =\mathrm{R}_{3}(\psi) \mathrm{R}_{\mathrm{n}}(\theta) \mathrm{R}_{z}(\phi) \\
& =\left(\begin{array}{ccc}
\cos \psi & \sin \psi & 0 \\
-\sin \psi & \cos \psi & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & \cos \theta & \sin \theta \\
0 & -\sin \theta & \cos \theta
\end{array}\right)\left(\begin{array}{ccc}
\cos \phi & \sin \phi & 0 \\
-\sin \phi & \cos \phi & 0 \\
0 & 0 & 1
\end{array}\right) .
\end{aligned}
$$

(a) Calculate the matrix $A=\dot{R} R^{\top}$ and verify that it is antisymmetric. By identifying $A$ as the antisymmetric matrix associated with the angular velocity vector in the body frame, deduce that

$$
\left(\begin{array}{c}
\omega_{1} \\
\omega_{2} \\
\omega_{3}
\end{array}\right)=\left(\begin{array}{c}
\dot{\theta} \cos \psi+\dot{\phi} \sin \theta \sin \psi \\
-\dot{\theta} \sin \psi+\dot{\phi} \sin \theta \cos \psi \\
\dot{\psi}+\dot{\phi} \cos \theta
\end{array}\right)
$$

(b) * [optional extra] Calculate the matrix $B=R^{\top} \dot{R}$ and verify that it is antisymmetric. By identifying B as the antisymmetric matrix associated with the angular velocity vector in the space frame, deduce that

$$
\left(\begin{array}{c}
\omega_{x} \\
\omega_{y} \\
\omega_{z}
\end{array}\right)=\left(\begin{array}{c}
\dot{\theta} \cos \phi+\dot{\psi} \sin \theta \sin \phi \\
\dot{\theta} \sin \phi-\dot{\psi} \sin \theta \cos \phi \\
\dot{\phi}+\dot{\psi} \cos \theta
\end{array}\right) .
$$

## 3. Free symmetric top

(a) Consider the torque-free motion of a round plate. Show that, in the body frame, the angular velocity vector $\boldsymbol{\omega}$ precesses around the body axis $\mathbf{e}_{3}$ perpendicular to the plate with constant angular frequency equal to $\omega_{3}$.
(b) The physicist Richard Feynman tells the following story:
> "I was in the cafeteria and some guy, fooling around, throws a plate in the air. As the plate went up in the air I saw it wobble, and I noticed the red medallion of Cornell on the plate going around. It was pretty obvious to me that the medallion went around faster than the wobbling.
> I had nothing to do, so I start figuring out the motion of the rotating plate. I discover that when the angle is very slight, the medallion rotates twice as fast as the wobble rate - two to one. It came out of a complicated equation! I went on to work out equations for wobbles. Then I thought about how the electron orbits start to move in relativity. Then there's the Dirac equation in electrodynamics. And then quantum electrodynamics. And before I knew it....the whole business that I got the Nobel prize for came from that piddling around with the wobbling plate."

[Here the 'wobble' is associated with precession of the top about $\mathbf{e}_{z}$ (motion in $\phi$ ), not nutation (motion in $\theta$ ).]
Feynman was right about quantum electrodynamics. But what about the plate?
[You could try two alternative methods. First, by using the expression for $\omega_{3}$ in terms of Euler's angles together with the result of part (a). Second, by writing down the Lagrangian of the top and deriving the equation of motion for $\theta$.]
(c) Consider a uniform symmetric ellipsoid of mass $M$ with semi-axes $a=b \neq c$ (see Example 2.5(e)). Find the ratio of the semi-axes for which $\dot{\phi}$, the angular frequency of precession of the top about the angular momentum $\mathbf{L}$, equals $\omega_{3} /(5 \cos \theta)$. Deduce further that $\dot{\psi}=\frac{4}{5} \omega_{3}$. What is the relationship between $\dot{\phi}$ and $\dot{\psi}$ for small values of $\theta$ ? Compare with the result obtained in part (b).

## 4. Free asymmetric top (1)

(a) Throw a book in the air. (Secure it with an elastic band first!) If the principal moments of inertia are $I_{3}>I_{2}>I_{1}$, convince yourself that the book can rotate in a stable manner about the principal axes $\mathbf{e}_{1}$ and $\mathbf{e}_{3}$, but not about $\mathbf{e}_{2}$.
(b) Use Euler's equations to show that the energy $E$ and the squared angular momentum $|\mathbf{L}|^{2}$ of a free asymmetric top are conserved. Suppose that the initial conditions are such that

$$
|\mathbf{L}|^{2}=2 E I_{2}
$$

with the initial angular velocity $\boldsymbol{\omega}$ perpendicular to the intermediate principal axis $\mathbf{e}_{2}$. Show that $\boldsymbol{\omega}$ will ultimately end up parallel to $\mathbf{e}_{2}$. What is the characteristic timescale required to reach this steady state?

## 5. Free asymmetric top (2)

A rigid lamina (i.e. a two-dimensional object) has principal moments of inertia about the centre of mass given by

$$
I_{1}=\left(\mu^{2}-1\right), \quad I_{2}=\left(\mu^{2}+1\right), \quad I_{3}=2 \mu^{2}
$$

where $\mu>1$. Write down Euler's equations for the lamina moving freely in space. Show that the component of the angular velocity in the plane of the lamina (i.e. $\sqrt{\omega_{1}^{2}+\omega_{2}^{2}}$ ) is constant in time.

Choose the initial angular velocity to be $\boldsymbol{\omega}=\mu N \mathbf{e}_{1}+N \mathbf{e}_{3}$. Define $\tan \alpha=\omega_{2} / \omega_{1}$, which is the angle the component of $\boldsymbol{\omega}$ in the plane of the lamina makes with $\mathbf{e}_{1}$. Show that it satisfies

$$
\ddot{\alpha}+N^{2} \cos \alpha \sin \alpha=0
$$

and deduce that, at time $t$,

$$
\boldsymbol{\omega}=\mu N \operatorname{sech}(N t) \mathbf{e}_{1}+\mu N \tanh (N t) \mathbf{e}_{2}+N \operatorname{sech}(N t) \mathbf{e}_{3} .
$$

## 6. Lagrange top

Consider a heavy symmetric top of mass $M$, fixed at the point $P$ which is a distance $l$ from the centre of mass. The principal moments of inertia about $P$ are $\left(I_{1}, I_{1}, I_{3}\right)$ and the Euler angles are defined as in the lectures. The top is spun with initial conditions $\dot{\phi}=0$ and $\theta=\theta_{0}$. Show that $\theta$ obeys the equation of motion

$$
I_{1} \ddot{\theta}=-\frac{d V_{\mathrm{eff}}}{d \theta},
$$

where

$$
V_{\mathrm{eff}}(\theta)=\frac{I_{3}^{2} \omega_{3}^{2}}{2 I_{1}} \frac{\left(\cos \theta-\cos \theta_{0}\right)^{2}}{\sin ^{2} \theta}+M g l \cos \theta
$$

Suppose that the top is spinning very fast so that

$$
I_{3} \omega_{3} \gg \sqrt{M g l I_{1}} .
$$

Show that the minimum of $V_{\text {eff }}(\theta)$ is close to $\theta_{0}$. Use this fact to deduce that the top nutates with angular frequency

$$
\Omega \approx \frac{I_{3}}{I_{1}} \omega_{3}
$$

and sketch the subsequent motion.

## 7. Lagrange top in Hamiltonian formalism

The Lagrangian for the heavy symmetric top is

$$
L=\frac{1}{2} I_{1}\left(\dot{\theta}^{2}+\dot{\phi}^{2} \sin ^{2} \theta\right)+\frac{1}{2} I_{3}(\dot{\psi}+\dot{\phi} \cos \theta)^{2}-M g l \cos \theta .
$$

Find the conjugate momenta $p_{\theta}, p_{\phi}$ and $p_{\psi}$ and the Hamiltonian $H\left(\theta, \phi, \psi, p_{\theta}, p_{\phi}, p_{\psi}\right)$. Derive Hamilton's equations.
8. Hamilton's equations

A system with two degrees of freedom $x$ and $y$ has the Lagrangian

$$
L=x \dot{y}+y \dot{x}^{2}+\dot{x} \dot{y} .
$$

Derive Lagrange's equations. Obtain the Hamiltonian $H\left(x, y, p_{x}, p_{y}\right)$. Derive Hamilton's equations and show that they are equivalent to Lagrange's equations.

Please send any comments and corrections to gio10@cam.ac.uk

