## The Foundations of Infinite-Dimensional Spectral Computations:

Can we compute spectral properties?
Matthew Colbrook
University of Cambridge


Spectral measure of magnetic graphene, computed to high precision (see log scale) using a method of this talk.

## Collaborators for papers in this talk

Anders Hansen
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Bogdan Roman (Cambridge)


Andrew Horning (Cornell)


Alex Townsend (Cornell)


## The infinite-dimensional problem

In discrete setting, operator acting on $\ell^{2}(\mathbb{N})$ :

$$
A=\left(\begin{array}{cccc}
a_{11} & a_{12} & a_{13} & \cdots \\
a_{21} & a_{22} & a_{23} & \cdots \\
a_{31} & a_{32} & a_{33} & \cdots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right), \quad(A x)_{j}=\sum_{k \in \mathbb{N}} a_{j k} x_{k} .
$$

In cts setting, deal with differential operators, integral operators etc.

| Finite Case | Infinite Case |
| :--- | :--- | :--- |
| Eigenvalues $\Rightarrow$ | Spectrum |
|  | $\operatorname{Sp}(A)=\{z \in \mathbb{C}: A-z$ not bounded invertible $\}$ |
| Eigenvectors $\Rightarrow$ | Spectral Measure (normal case) |

Goal: compute spectral properties of the operator from matrix elements, PDE coefficients, or other suitable information.

MUCH harder and more subtle than finite dimensions!

Many applications: quantum mechanics, chemistry, matter physics, statistical mechanics, optics, number theory, PDEs, mathematics of information etc.

Mathematicians and physicists contributing to computational spectral theory form a vast set including:
D. Arnold (Minnesota), W. Arveson (Berkeley), A. Böttcher (Chemnitz), W. Dahmen (South Carolina), E. B. Davies (King's College London), P. Deift (NYU), L. Demanet (MIT), C. Fefferman (Princeton), G. Golub (Stanford), A. Iserles (Cambridge), W. Schlag (Yale), E. Schrödinger (DIAS), J. Schwinger (Harvard), N. Trefethen (Oxford), V. Varadarajan (UCLA), S. Varadhan (NYU), J. von Neumann (IAS), M. Zworski (Berkeley),...

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"Most operators that arise in practice are not presented in a representation in which they are diagonalized, and it is often very hard to locate even a single point in the spectrum... Thus, one often has to settle for numerical approximations [to the spectrum], and this raises the question of how to implement the methods of finite dimensional numerical linear algebra to compute the spectra of infinite dimensional operators. Unfortunately, there is a dearth of literature on this basic problem and, so far as we have been able to tell, there are no proven techniques."

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Partial answer: can compute spectra of general bounded operators on $\ell^{2}(\mathbb{N})$ (in Hausdorff metric) using three successive limits ${ }^{1}$

$$
\lim _{n_{3} \rightarrow \infty} \lim _{n_{2} \rightarrow \infty} \lim _{n_{1} \rightarrow \infty} \Gamma_{n_{3}, n_{2}, n_{1}}(A)=\operatorname{Sp}(A) .
$$

Turns out this is sharp! Hence impossible from numerical point of view.

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## Can we do this for general classes of operators?

"Most operators that arise in practice are not presented in a representation in which they are diagonalized, and it is often very hard to locate even a single point in the spectrum... Thus, one often has to settle for numerical approximations [to the spectrum], and this raises the question of how to implement the methods of finite dimensional numerical linear algebra to compute the spectra of infinite dimensional operators. Unfortunately, there is a dearth of literature on this basic problem and, so far as we have been able to tell, there are no proven techniques."

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Turns out this is sharp! Hence impossible from numerical point of view. Q: What assumptions do we need to make it easier? What about other spectral problems?
${ }^{1}$ Hansen. JAMS (2011)

## Example: Bounded Diagonal Operators (Very Easy)

$$
A=\left(\begin{array}{llll}
a_{1} & & & \\
& a_{2} & & \\
& & a_{3} & \\
& & & \ddots
\end{array}\right)
$$

If $\Gamma_{n}(A)=\left\{a_{1}, \ldots, a_{n}\right\}$ then $\Gamma_{n}(A) \rightarrow \operatorname{Sp}(A)$ in Hausdorff metric.
Also have $\Gamma_{n}(A) \subset \operatorname{Sp}(A)$.
This is optimal from a foundations point of view.

## Example: Compact Operators (Still Easy)

$$
A=\left(\begin{array}{cccc}
a_{11} & a_{12} & a_{13} & \cdots \\
a_{21} & a_{22} & a_{23} & \cdots \\
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\end{array}\right), \text { compact }
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If $\Gamma_{n}(A)=\operatorname{Sp}\left(P_{n} A P_{n}\right)$, then $\Gamma_{n}(A) \rightarrow \operatorname{Sp}(A)$ in Hausdorff metric.
Known for decades.

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Known for decades.
Q: Can we gain error control as before?
No! No algorithm can gain error control on the whole class, even for self-adjoint compact operators.

What about Jacobi operators?

$$
A=\left(\begin{array}{llll}
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b_{1} & a_{2} & b_{2} & \\
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What about sparse normal operators? Surely this is much harder?!
New result: Large class $\Omega$ (covering arguably most applications and including sparse normal) such that we can compute $\Gamma_{n}(A) \rightarrow \operatorname{Sp}(A)$ and $E_{n}(A) \downarrow 0$ for $A \in \Omega$ with

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Paradox: Easier problem than compact operators!

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Paradox: Easier problem than compact operators!
Can also be extended to many non-self-adjoint Schrödinger operators, singular potentials, more general partial differential operators etc.

## Foundations of Infinite-Dimensional Spectral Computations

How: Deal with operators directly, instead of previous 'truncate-then-solve'
$\Rightarrow$ Compute many spectral properties for the first time.

Framework: Classify problems in a computational hierarchy (the SCI hierarchy) measuring their difficulty and the optimality of algorithms. ${ }^{2}$
$\Rightarrow$ Algorithms that realise the boundaries of what computers can achieve.

NB: This holds regardless of model of computation - measure the intrinsic difficulty of these problems.

Common theme in examples of this talk: use the resolvent $(A-z)^{-1}$

[^1]
## Structure of the Hierarchy



## Example 1: Computing spectra with error control

Computing the resolvent norm $\gamma(z):=\left\|(A-z)^{-1}\right\|^{-1}$ $\Downarrow$

First algorithm that computes spectra of a very general class of operators. Also does so with (rigorous provable) error control.

For simplicity, assume $A$ acts on $I^{2}(\mathbb{N})$.

## Definition (Dispersion: off-diagonal decay)

Dispersion of $A$ bounded by function $f: \mathbb{N} \rightarrow \mathbb{N}$ and null sequence $\left\{c_{n}\right\}$ if

$$
\max \left\{\left\|\left(I-P_{f(n)}\right) A P_{n}\right\|,\left\|P_{n} A\left(I-P_{f(n)}\right)\right\|\right\} \leq c_{n}
$$



## Definition (Controlled growth of the resolvent: well-conditioned)

Continuous increasing function $g:[0, \infty) \rightarrow[0, \infty)$ with $g(x) \leq x$. Controlled growth of the resolvent by $g$ if

$$
g(\operatorname{dist}(z, \operatorname{Sp}(A))) \leq\left\|(A-z)^{-1}\right\|^{-1} \quad \forall z \in \mathbb{C}
$$

- Measures conditioning of the problem through

$$
\left\{z \in \mathbb{C}:\left\|(A-z)^{-1}\right\|^{-1} \leq \epsilon\right\}=: \operatorname{Sp}_{\epsilon}(A)=\bigcup_{\|B\| \leq \epsilon} \operatorname{Sp}(A+B) .
$$

- Normal operators ( $A$ commutes with $A^{*}$ ) well-conditioned with

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## Theorem (C., Roman, Hansen. PRL (2019))

Know $f, g \Rightarrow$ can compute $\operatorname{Sp}(A)$ with error control.

Step 1: Approximate locally via smallest singular value:

$$
\gamma_{n}(z):=\min \left\{\sigma_{1}\left(P_{f(n)}(A-z) P_{n}\right), \sigma_{1}\left(P_{f(n)}\left(A^{*}-\bar{z}\right) P_{n}\right)\right\}+c_{n} .
$$

This converges locally uniformly down to $\gamma(z)=\left\|(A-z)^{-1}\right\|^{-1}$.
Step 2: Bound the distance to the spectrum:

$$
\gamma(z) \leq \operatorname{dist}(z, \operatorname{Sp}(A)) \leq g^{-1}(\gamma(z)) \leq g^{-1}\left(\gamma_{n}(z)\right)
$$

Step 3: Find (almost) local minimisers and output $\Gamma_{n}(A)$ with

$$
\Gamma_{n}(A) \rightarrow \operatorname{Sp}(A), \quad \operatorname{dist}(z, \operatorname{Sp}(A)) \leq g^{-1}\left(\gamma_{n}(z)\right), \sup _{z \in \Gamma_{n}(A)} g^{-1}\left(\gamma_{n}(z)\right) \rightarrow 0
$$

NB: Without either $f$ or $g$, constructing an algorithm convergent to the spectrum is impossible (even dropping error control).

Example: quartic potential on $L^{2}(\mathbb{R})$ using a Hermite basis


## New exemplar of spectral computation

Method is:

- Local and parallelisable.
- Convergent for first time. (e.g. no spectral pollution)
- Explicitly bounds the error:

$$
\text { Error } \leq a_{n} \downarrow 0
$$

- Optimal from foundations point of view.



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Extends to unbounded operators and PDEs with coefficients of locally bounded total variation (e.g. algorithms point sample coefficients).

## Example: Operators in condensed matter physics



Left: Dan Shechtman, Nobel Prize in Chemistry 2011 for discovery of quasicrystal. Right: Diffraction pattern of a quasicrystal.

Magnetic properties of quasicrystal.
Hard problem - no previous method even converges to spectrum.

Example: Operators in condensed matter physics

Finite truncations
Edge states.


Unreliable
Does not converge No error control

## Infinite-dimensional techniques

First convergent computation.


Reliable
Converges
Error control

## Example: Laplacian on Penrose tile



## Example: Laplacian on Penrose tile



## Example 2: Computing spectral measures

Compute the resolvent $(A-z)^{-1} \mathbf{x}(\operatorname{Im}(z)>0)$ with asymptotic error control as $z$ approaches $\mathbb{R}$.
$\Downarrow$
First algorithm that computes spectral measures of general self-adjoint operators - 'diagonalisation.'

Assume for rest of talk that $A$ is self-adjoint.

## Spectral Measures

If $A$ normal, associated projection-valued measure $\mathcal{E}$ s.t.

$$
A \mathbf{x}=\int_{\operatorname{Sp}(A)} \lambda d \mathcal{E}(\lambda) \mathbf{x}, \quad \forall \mathbf{x} \in \mathcal{D}(A)
$$

Allows computation of functional calculus, has interesting physics etc.
Idea: For $z=x+i \epsilon$, use

$$
\frac{(A-z)^{-1}-(A-\bar{z})^{-1}}{2 \pi i}=\frac{1}{\pi} \int_{\operatorname{Sp}(A)} \frac{\epsilon}{(x-\lambda)^{2}+\epsilon^{2}} d \mathcal{E}(\lambda)
$$

Convolution with Poisson kernel: smoothed measure.
Converges weakly to measure as $\epsilon \downarrow 0$ (cf. Stone's formula).

## Example: Magnetic Graphene




## Example: Magnetic Graphene



## Example: Magnetic Graphene



## Theorem (C. Preprint (2019))

Know $f \Rightarrow$ can compute measure in one limit.
This is through a rectangular least squares type problem that computes $(A-z)^{-1} \mathbf{x}$ with (asymptotic) error control. $N(\epsilon)$ chosen adaptively.


## Example: Integral Operator

$$
A u(x)=x u(x)+\int_{-1}^{1} e^{-\left(x^{2}+y^{2}\right)} u(y) d y, \quad x \in[-1,1] .
$$

Discretise using adaptive Chebyshev collocation method.
Look at $\mu_{f}: B \rightarrow\langle\mathcal{E}(B) f, f\rangle$. E.g. $f(x)=\sqrt{3 / 2} x$.


## Example: Integral Operator




Converges like $\mathcal{O}\left(\epsilon \log \left(\epsilon^{-1}\right)\right)$ and need $N \approx 20 / \epsilon$.
$\Rightarrow$ Infeasible to get more than five or six digits!
Q: Can we do better?

## Accelerating convergence

Let $m \in \mathbb{N}, K \in L^{1}(\mathbb{R})$. We say $K$ is an $m$ th order kernel if:
(i) Normalized: $\int_{\mathbb{R}} K(x) d x=1$,
(ii) Zero moments: $K(x) x^{j}$ is integrable and $\int_{\mathbb{R}} K(x) x^{j} d x=0$ for $0<j<m$, and
(iii) Decay at $\pm \infty$ : There is a constant $C_{K}$, independent of $x$, such that

$$
|K(x)| \leq C_{K}(1+|x|)^{-(m+1)}, \quad x \in \mathbb{R}
$$

## Theorem (C., Horning, Townsend. Preprint (2020))

If $K$ is $m$ th order, $K_{\epsilon}(x)=\epsilon^{-1} K\left(x \epsilon^{-1}\right)$ and $\mu_{f}$ locally absolutely continuous near $x_{0}$ with density $\rho_{f}$ then

- Pointwise: If $\rho_{f}$ locally $\mathcal{C}^{n, \alpha}$ near $x_{0}$ then

$$
\left|\left[K_{\epsilon} * \mu_{f}\right]\left(x_{0}\right)-\rho_{f}\left(x_{0}\right)\right|=\mathcal{O}\left(\epsilon^{n+\alpha}\right)+\mathcal{O}\left(\epsilon^{m} \log \left(\epsilon^{-1}\right)\right)
$$

- $L^{p}$ : If $\rho_{f}$ locally $\mathcal{W}^{n, p}$ near $x_{0}(1 \leq p<\infty)$ then

$$
\left\|\left[K_{\epsilon} * \mu_{f}\right]-\rho_{f}\right\|_{L_{\text {loc }}^{p}}=\mathcal{O}\left(\epsilon^{n}\right)+\mathcal{O}\left(\epsilon^{m} \log \left(\epsilon^{-1}\right)\right)
$$

## Accelerating convergence with rational kernels

Idea: Replace Poisson kernel with rational kernel

$$
K(x)=\frac{1}{2 \pi i} \sum_{j=1}^{n_{1}} \frac{\alpha_{j}}{x-a_{j}}-\frac{1}{2 \pi i} \sum_{j=1}^{n_{2}} \frac{\beta_{j}}{x-b_{j}}
$$

Can compute convolution with error control using resolvent

$$
\begin{aligned}
& {\left[K_{\epsilon} * \mu_{f}\right](x)} \\
& =\frac{1}{2 \pi i}\left[\sum_{j=1}^{n_{1}} \alpha_{j}\left\langle\left(A-\left(x+\epsilon a_{j}\right)\right)^{-1} f, f\right\rangle-\sum_{j=1}^{n_{2}} \beta_{j}\left\langle\left(A-\left(x+\epsilon b_{j}\right)\right)^{-1} f, f\right\rangle\right] .
\end{aligned}
$$

## Proposition (C., Horning, Townsend. Preprint (2020))

Let $n_{1}=n_{2}=m \in \mathbb{N}$ and fix $a_{j}$ in UHP, $b_{j}$ in LHP. Then there exists unique $\left\{\alpha_{j}, \beta_{j}\right\}$ such that $K$ is an mth order kernel.

## Integral operator revisited




See paper for general differential (even PDEs), integral and lattice operator examples - use sparse spectral methods for discretisation.

## Eigenvalue hunting without spectral pollution

Can extend to compute decompositions (of measures and spectral sets).
Example: Dirac operator.

- Describes the motion of a relativistic spin-1/2 particle.
- Essential spectrum given by $\mathbb{R} \backslash(-1,1) \Rightarrow$ spectral pollution!
- Consider radially symmetric potential, coupled system on half-line:

$$
\mathcal{D}_{V}=\left(\begin{array}{cc}
1+V(r) & -\frac{d}{d r}+\frac{\kappa}{r} \\
\frac{d}{d r}+\frac{\kappa}{r} & -1+V(r)
\end{array}\right) .
$$

- Map to $[-1,1]$ using

$$
x=\frac{r-L}{r+L}
$$

and compute $(A-z)^{-1} \mathbf{x}$ using the ultraspherical method.

## Eigenvalue hunting without spectral pollution

$$
\nu_{f}^{\epsilon}(x):=\epsilon \cdot \operatorname{Im}\left(\left\langle\left(\mathcal{D}_{V}-(x+i \epsilon)\right)^{-1} f, f\right\rangle\right)=\sum_{\lambda \in \Lambda^{\mathrm{P}}\left(\mathcal{D}_{V}\right) \cap\{x\}}\left\langle\mathcal{P}_{\lambda} f, f\right\rangle+\mathcal{O}(\epsilon)
$$



Absolute Error


NB: Previous state-of-the-art achieves a few digits for a few excited states.

## Results in this talk:

- Can compute spectra of a large class of operators with error control. New algorithm is fast, local and parallelisable.
- Can compute spectral measures and spectral decompositions through resolvent. Can be combined with state-of-the-art PDE methods.

Also have foundations for: spectral type (pure point, absolutely continuous, singularly continuous), Lebesgue measure and fractal dimensions of spectra, discrete spectra, essential spectra, eigenvectors + multiplicity, spectral radii, essential numerical ranges, geometric features of spectrum (e.g. capacity), spectral gap problem, ...

Ongoing and future work: foundations of computational PDEs, foundations of (stable) neural networks, and computer-assisted proofs.

Code: high-performance numerical package with resolvent based algorithms for computing spectral measures (written with Andrew Horning). https://github.com/ajhPHROS/SpecSolve

## References for algorithms in this talk

- M.J. Colbrook, B. Roman, and A.C. Hansen. "How to compute spectra with error control." Physical Review Letters 122.25 (2019).
- M.J. Colbrook, A.C. Hansen. "On the infinite-dimensional QR algorithm." Numerische Mathematik 143.1 (2019).
- M.J. Colbrook. "Computing spectral measures and spectral types." Submitted.
- M.J. Colbrook, A. Horning, and A. Townsend. "Computing spectral measures of self-adjoint operators." Submitted.

For further papers in this program, classifications in the SCI hierarchy and numerical code: http://www.damtp.cam.ac.uk/user/mjc249/home.html

If you have further ideas or problems for collaboration, please get in touch!


[^0]:    ${ }^{1}$ Hansen. JAMS (2011)

[^1]:    ${ }^{2}$ Also has deep connections with logic and descriptive set theory: C. Preprint (2019)

