#### Hyper-Kähler metrics from isomonodromy

Timothy Moy

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# Hyper-Kähler metrics from isomonodromy

Timothy Moy joint work with Maciej Dunajski, arXiv:2402.14352

> Glasgow 16 April 2024

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#### Slides available at damtp.cam.ac.uk/tjahm2

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Some slightly vague motivating remarks:

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Some slightly vague motivating remarks:

Bridgeland introduced Joyce structures as a geometric structure that should exist on the space *M* of stability conditions of a *CY*<sub>3</sub> triangulated category.

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- Bridgeland introduced Joyce structures as a geometric structure that should exist on the space *M* of stability conditions of a *CY*<sub>3</sub> triangulated category.
- The argument for their existence involves DT invariants... won't talk about this today

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- The argument for their existence involves DT invariants... won't talk about this today
- Joyce structure: complex hyperkähler g metric on X = TM with a homothetic Killing vector field  $(\mathcal{L}_W g = g)$  and some lattice invariance conditions (see Bridgeland-Strachan (2021) for a precise geometric definition)

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### Lax distributions

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In what follows let X be a complex manifold of dimension 4n and TX the holomorphic tangent bundle. For our purposes:

#### Definition (Lax distribution)

A (hyper-Hermitian) Lax distribution is a subbundle of TX

$$L(\lambda) = \operatorname{span}\left\{v_i + \lambda h_i\right\}_{i=1}^{2n}.$$
 (1)

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depending on a spectral parameter  $\lambda \in \mathbb{C}$ , where  $v_i$ ,  $h_i$  are vector fields on X such that  $TX = span\{v_i, h_i\}_{i=1}^{2n}$ 

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depending on a spectral parameter  $\lambda \in \mathbb{C}$ , where  $v_i$ ,  $h_i$  are vector fields on X such that  $TX = span\{v_i, h_i\}_{i=1}^{2n}$ 

Letting  $M = X/span\{v_i\}$  we can think of  $L(\lambda)$  as a family of (not necessarily linear) Ehresmann connections on the bundle  $X \to M$  depending 'linearly' on  $\lambda$ .

### Quaternionic structure

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$$L(\lambda) = \text{span} \left\{ v_i + \lambda h_i \right\}_{i=1}^{2n}$$
 determines a quaternionic structure:

$$I(v_i) = iv_i, \quad J(v_i) = -h_i, \quad K(v_i) = ih_i$$
(2)

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 $L(\lambda)$  Frobenius integrable for each  $\lambda$  is equivalent to integrability of the complex structures I, J, K and the existence of the twistor space  $\mathcal{Z}$ .

### Associated family of hyper-Hermitian metrics

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 $L(\lambda) = \text{span} \{v_i + \lambda h_i\}_{i=1}^{2n}$  also determines a family of holomorphic metrics satisfying the *hyper-Hermitian condition*  $I^*g = J^*g = K^*g = g$ :

$$g = \sum_{i,j=1}^{2n} e_{ij} h^i \odot v^i$$
(3)

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each corresponding to a non-degenerate skew matrix  $e_{ij}$  of holomorphic functions.

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n = 1 case well studied (e.g. Penrose (1976)): conformal class of holomorphic metrics on 4-dimensional X and  $L(\lambda)$  is the twistor distribution. Frobenius integrability  $\iff g$  has anti-self-dual Weyl tensor

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n = 1 case well studied (e.g. Penrose (1976)): conformal class of holomorphic metrics on 4-dimensional X and  $L(\lambda)$  is the twistor distribution. Frobenius integrability  $\iff g$  has anti-self-dual Weyl tensor n > 1 these are almost Grassmannian (or paraconformal)

geometries. See (Bailey-Eastwood (1991))

## Complex hyper-Kähler metrics

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#### Definition (Complex hyper-Kähler)

A complex hyper-Kähler structure is a holomorphic metric g and triple of holomorphic endomorphisms I, J, K of TX satisfying the quaternion relations such that g is Hermitian for each and  $\nabla I = \nabla J = \nabla K = 0$ .

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When is there a hyper-Kähler metric in the class?

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When is there a hyper-Kähler metric in the class? Conformal case: (Mason-Newman (1989))

$$[\mathbf{v}_1 + \lambda \mathbf{h}_1, \mathbf{v}_2 + \lambda \mathbf{h}_2] = \mathbf{0}$$

and the flows preserved a volume form  $vol_X$ .

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When is there a hyper-Kähler metric in the class? Conformal case: (Mason-Newman (1989))

 $[v_1 + \lambda h_1, v_2 + \lambda h_2] = 0$ 

and the flows preserved a volume form  $\operatorname{vol}_X$ . General case (sufficient condition): Integrability of  $L(\lambda) \forall \lambda$  and the existence of  $\omega$ , a symplectic form on  $M := X/\operatorname{span}\{v_i\}$  with

$$\mathcal{L}_{h_i}(\pi^*\omega)=0, \quad i=1,...,2n_i$$

### Plebaśnki's heavenly equation

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There is a hyperkähler metric in the class if and only if there exists coordinates  $(x^i, y^i)$  such that

$$v_{i} = \frac{\partial}{\partial y^{i}}$$

$$h_{i} = \frac{\partial}{\partial x_{i}} + \sum_{j,k=1}^{2n} \eta^{jk} \frac{\partial^{2} \Theta}{\partial y^{i} \partial y^{j}} \frac{\partial}{\partial y^{k}}$$
(5)

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where  $\Theta(x_i, y_i)$  satisfies (a higher dimensional version of) Plebański's second heavenly equation

$$\frac{\partial^2 \Theta}{\partial y^i \partial x^j} - \frac{\partial^2 \Theta}{\partial y^j \partial x^i} - \sum_{k,l=1}^{2n} \eta^{kl} \frac{\partial^2 \Theta}{\partial y^j \partial y^k} \frac{\partial^2 \Theta}{\partial y^j \partial y^l} = 0.$$
(6)

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We will consider particular spaces X parametrising ODE on which will live  $L(\lambda) \subseteq TX$ .

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We will consider particular spaces X parametrising ODE on which will live  $L(\lambda) \subseteq TX$ .

The general solutions of linear ODE with meromorphic coefficients may have branching behaviour near poles  $x_0, ..., x_M$  of the coefficients. The fundamental group of the punctured space  $\mathbb{CP}^1 \setminus \{x_0, ..., x_M\}$  then has a linear representation on the space of solutions called the *monodromy*.

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Given a family of ODE depending on some parameters, a *isomonodromic flow* is a family of deformations of the parameters continuous with the identity which preserves the monodromy.

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Given a family of ODE depending on some parameters, a *isomonodromic flow* is a family of deformations of the parameters continuous with the identity which preserves the monodromy.

The situation for irregular singularities is more complicated... (want to also preserve Stokes' data)

# Equation for isomonodromic flows of 2nd order linear ODE

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For the purposes of our problem the output of the general theory is the following theorem:

#### Theorem (Schlesinger / Jimbo-Miwa-Ueno(1981))

A one-parameter family Q(t, x) of deformations of the potential for the ODE

$$\frac{d^2y}{dx^2} = Q(x)y$$

varying holomorphically with the parameter t has constant (generalised) monodromy if and only if

$$\frac{\partial Q(x,t)}{\partial t} = 2Q(x)\frac{\partial A(t,x)}{\partial x} + \frac{\partial Q(t,x)}{\partial x}A(t,x) - \frac{1}{2}\frac{\partial^3 A(t,x)}{\partial x^3}$$
  
for some  $A(t,x)$ .

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### Geometry of isomonodromic flows

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If potentials are parametrised by a space X with coordinates  $w^i$ , then a one-parameter family of deformations corresponds to a vector field U satisfying the Schlesinger equation:

$$U(Q(x, w^{i})) = 2Q\frac{\partial A}{\partial x} + \frac{\partial Q}{\partial x}A - \frac{1}{2}\frac{\partial^{3}A}{\partial x^{3}}$$
(7)  
for some  $A(x, w^{i})$ .

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(7)

for some  $A(x, w^i)$ . Suppose we have another:

$$V(Q(x,w^{i})) = 2Q\frac{\partial B}{\partial x} + \frac{\partial Q}{\partial x}B - \frac{1}{2}\frac{\partial^{3}B}{\partial x^{3}}.$$
 (8)

Proposition (Lie bracket of isomonodromic flows)

$$[U, V](Q(x, w^{i})) = 2Q\frac{\partial C}{\partial x} + \frac{\partial Q}{\partial x}C - \frac{1}{2}\frac{\partial^{3} C}{\partial x^{3}}$$

where

$$C = U(B) - V(A) - \left(A\frac{\partial B}{\partial x} - B\frac{\partial A}{\partial x}\right)$$

### Deformed cubic oscillator I

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The following example is due to Bridgeland-Masoero (2022):

$$Q(x) = \frac{Q_0(x)}{\lambda^2} + \frac{Q_1(x)}{\lambda} + Q_2(x)$$

where

$$Q_0(x) = x^3 + ax + b$$
 (9)  
 $Q_1(x) = \frac{p}{x-q} + r$  (10)

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where  $p^2 = q^3 + aq + b$ .

### Deformed cubic oscillator I

where  $p^2 = q^3 + aq + b$ .

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where

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 (9)  
 $Q_1(x) = \frac{p}{x - q} + r$  (10)

Loosely, we pick  $Q_2(x)$  to be the simplest function so that the ODE, written as a first order system, has no singularity at q after a gauge transformation. Specifically

$$Q_2(x) = \frac{3}{4(x-q)^2} + \frac{r}{2p(x-q)} + \frac{r^2}{4p^2}.$$
 (11)

### Deformed cubic oscillator II

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$$Q(x) = \frac{Q_0(x)}{\lambda^2} + \frac{Q_1(x)}{\lambda} + Q_2(x)$$
(12)

$$Q_0(x) = x^3 + ax + b, \quad Q_1(x) = \frac{p}{x - q} + r$$
 (13)

$$Q_2(x) = \frac{3}{4(x-q)^2} + \frac{r}{2p(x-q)} + \frac{r^2}{4p^2}.$$
 (14)

The ODE is therefore specified by a point on a manifold X with local coordinates (a, b, q, r). The isomonodromic flows are of the right number and have the right form to define a Lax pair:

$$U = -\frac{\partial}{\partial r} + \lambda \left( \frac{\partial}{\partial b} + \frac{r}{2p^2} \frac{\partial}{\partial r} \right)$$
(15)  
$$V = -2p \frac{\partial}{\partial q} + \lambda \left( \frac{\partial}{\partial a} - \frac{r}{p} \frac{\partial}{\partial q} - \frac{r(3q^2r + ar - qp)}{2p^3} \frac{\partial}{\partial r} \right)$$

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There is a hyper-Kähler metric on X in the conformal class:

$$g^{\omega} = \left(\frac{r(3q^2r + ar - 2qp)}{2p^3}da - \frac{r}{2p^2}db - \frac{q}{2p}dq + dr\right) \odot da$$
$$-\left(\frac{r}{2p^2}da + \frac{1}{2p}dq\right) \odot db.$$

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$$-\left(\frac{r}{2p^2}da + \frac{1}{2p}dq\right) \odot db.$$

with homothetic Killing vector

$$W = \frac{4a}{5}\frac{\partial}{\partial a} + \frac{6b}{5}\frac{\partial}{\partial b} + \frac{2q}{5}\frac{\partial}{\partial q} + \frac{r}{5}\frac{\partial}{\partial r}.$$

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Recall Joyce structures were defined on the total space of the tangent bundle  $TM \rightarrow M$ . *M* a space of stability conditions. How to see this for the cubic oscillator?

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Recall Joyce structures were defined on the total space of the tangent bundle  $TM \rightarrow M$ . M a space of stability conditions. How to see this for the cubic oscillator? Bridgeland-Smith (2013) realise spaces of meromorphic quadratic differentials with fixed pole orders as spaces of stability conditions.

The choices of

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$$Q_0(x) = x^3 + ax + b$$

parametrised by (a, b) correspond to quadratic differentials  $Q_0(x)dx^2$  on  $\mathbb{CP}^1$  with a single pole of order 7 up to Möbius transformation.

The choices of

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 $Q_0(x) = x^3 + ax + b$ 

parametrised by (a, b) correspond to quadratic differentials  $Q_0(x)dx^2$  on  $\mathbb{CP}^1$  with a single pole of order 7 up to Möbius transformation.

So X fibres over M = (a, b) a space of stability conditions.

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To get a (local identification)  $TM \rightarrow X$  we note that a point in  $(a, b) \in M$  defines an elliptic curve

$$\Sigma_{(a,b)} = \{(x,y) \in \mathbb{C}^2 \mid y^2 = x^3 + ax + b\} \cup \{\infty\}$$
(16)

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$$\Sigma_{(a,b)} = \{ (x,y) \in \mathbb{C}^2 \mid y^2 = x^3 + ax + b \} \cup \{ \infty \}$$
 (16)

Consider the holomorphic vector bundle *E* of rank two with fibre  $H^1(\Sigma_{(a,b)}, \mathbb{C})$  at (a, b). It has a canonical connection  $\nabla^{GM} : T^*M \otimes \Gamma(E) \to \Gamma(E)$  the

*Gauss-Manin* connection: the flat connection with parallel sections those that take values in the fundamental co-cycles.

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Consider the holomorphic vector bundle E of rank two with fibre  $H^1(\Sigma_{(a,b)}, \mathbb{C})$  at (a, b). It has a canonical connection  $\nabla^{GM} : T^*M \otimes \Gamma(E) \to \Gamma(E)$  the *Gauss-Manin* connection: the flat connection with parallel sections those that take values in the fundamental co-cycles. We also have a canonical section Z with value at (a, b) given by

$$Z_{(a,b)} = [y \, dx] \in H^1(\Sigma_{(a,b)}, \mathbb{C})$$

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To get a (local identification)  $TM \rightarrow X$  we note that a point in  $(a, b) \in M$  defines an elliptic curve

$$\Sigma_{(a,b)} = \{ (x,y) \in \mathbb{C}^2 \mid y^2 = x^3 + ax + b \} \cup \{ \infty \}$$
 (16)

Consider the holomorphic vector bundle E of rank two with fibre  $H^1(\Sigma_{(a,b)}, \mathbb{C})$  at (a, b). It has a canonical connection  $\nabla^{GM} : T^*M \otimes \Gamma(E) \to \Gamma(E)$  the *Gauss-Manin* connection: the flat connection with parallel sections those that take values in the fundamental co-cycles. We also have a canonical section Z with value at (a, b) given by

$$Z_{(a,b)} = [y \, dx] \in H^1(\Sigma_{(a,b)}, \mathbb{C})$$

Then we get an isomorphism  $TM \cong E$  given by

$$v \mapsto \nabla_v^{GM} Z.$$

### Abelian holonomy map

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We claim 
$$(a, b, q, r) \in X$$
 defines a class in  $H^1(\Sigma_{(a,b)}, \mathbb{C}^{\times})$ :

$$Q_1(x) = \frac{p}{x-q} + r \leftrightarrow \varpi = 2\pi i \left(\frac{y+p}{x-q} + r\right) \frac{dx}{2y}$$
(17)

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a meromorphic one-form on  $\Sigma_{(a,b)}$  on the elliptic curve  $\Sigma_{(a,b)}$ with residues integer multiples of  $2\pi i \implies$  integration of  $\varpi$ over homology classes is well-defined after exponentiation.

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a meromorphic one-form on  $\Sigma_{(a,b)}$  on the elliptic curve  $\Sigma_{(a,b)}$ with residues integer multiples of  $2\pi i \implies$  integration of  $\varpi$ over homology classes is well-defined after exponentiation. We have maps over M:

$$X \longleftarrow E^{\times} \underset{\text{exp}}{\longleftarrow} E_{\underset{v \mapsto \nabla_{v}^{GM}}{\longleftarrow} Z} TM$$

Where E,  $E^{\times}$  have fibres  $H^1(\Sigma_{(a,b)}, \mathbb{C})$ ,  $H^1(\Sigma_{(a,b)}, \mathbb{C}^{\times})$  respectively. All this is made rigorous in Bridgeland-Masoero (2022).

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## Generalising the deformed cubic oscillator metric

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This example was the first non-trivial example of a Joyce structure with a description in local coordinates.

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This example was the first non-trivial example of a Joyce structure with a description in local coordinates.

In arXiv:2402.14352 we generalise this construction to produce explicit expressions for a complex hyper-Kähler metric in 4n dimensions from the isomonodromy of ODE with potentials having leading term a polynomial of degree 2n + 1.

### Deformed polynomial oscillator

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The ODE setup is the obvious generalisation:

$$Q(x) = \frac{Q_0(x)}{\lambda^2} + \frac{Q_1(x)}{\lambda} + Q_2(x)$$

#### where

$$Q_0(x) = x^{2n+1} + a_n x^{2n-1} + \dots + a_1 x^n + b_n x^{n-1} + \dots + b_1$$
(18)  
$$Q_1(x) = \sum_{i=1}^n \frac{p_i}{x - q_i} + R(x)$$
(19)

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where  $p_i^2 = q_i^3 + aq_i + b$  and R(x) is the general polynomial of degree at most n - 1 (parametrised  $(v_1, ..., v_n)$ ).

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where  $p_i^2 = q_i^3 + aq_i + b$  and R(x) is the general polynomial of degree at most n - 1 (parametrised  $(v_1, \dots v_n)$ ). Again  $Q_2(x)$  is picked so that there is no singularity at  $q_i$  after a gauge transformation.

## Isomonodromy result

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### Proposition (Dunajski-M,(2023))

The equation with potential specified by (18) and (19) and  $Q_2(x)$  chosen appropriately has 2n linearly independent isomonodromic flows of the form

$$L_i = v_i + \lambda h_i$$

where  $TX = span\{v_i, h_i\}_{i=1}^{2n}$  and the  $v_i$  are vertical for the projection  $X \to M$ .

The proof proceeds by breaking down the Schlesinger equation into manageable subsystems by Laurent expanding at the various poles  $\infty$ ,  $q_1$ , ...,  $q_n$ .

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So we have the Lax distribution and hence a family of metrics. How can we distinguish a hyper-Kähler metric?

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So we have the Lax distribution and hence a family of metrics. How can we distinguish a hyper-Kähler metric? Recall the family of metrics

$$g = e_{ij}h^i \odot v^j \tag{20}$$

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Choose  $e_{ij} = \omega_{ij}$ , the pull-back of the natural symplectic form  $\omega$  on M (affine symplectic fibration).

### Theorem (Dunajski-M(2023))

The metric

$$g^{\omega} = \omega_{ij} h^i \odot v^j$$

is complex hyper-Kähler.

We call it the  $A_{2n}$  complex hyper-Kähler metric.

## Symplectic structure on M

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To see this symplectic structure note  $p \in M$  defines a hyper-elliptic curve

$$\Sigma = \{y^2 = x^{2n+1} + a_n x^{2n-1} + \dots + a_1 x^n + b_n x^{n-1} + \dots + b_1\}$$
(21)

each with cohomology intersection form  $H^1(\Sigma, \mathbb{C}) \times H^1(\Sigma, \mathbb{C}) \to \mathbb{C}$ .

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(21)

each with cohomology intersection form  $H^1(\Sigma, \mathbb{C}) \times H^1(\Sigma, \mathbb{C}) \to \mathbb{C}$ .

Recall the Gauss-Manin connection defines an isomorphism  $T_{(a,b)}M \to H^1(\Sigma, \mathbb{C})$  by

$$v \mapsto \nabla_v^{GM} Z.$$

 $\omega$  is the pull-back of the intersection form by this isomorphism.

## Intriguing geometry: hyper-Lagrangians

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This countable family of hyper-Kähler metrics has some nice properties:

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This countable family of hyper-Kähler metrics has some nice properties: Recall the Plebański potential  $\Theta$ . Such that we may write

$$h_{i} = \frac{\partial}{\partial x_{i}} + \sum_{j,k=1}^{2n} \eta^{jk} \frac{\partial^{2}\Theta}{\partial y^{i} \partial y^{j}} \frac{\partial}{\partial y^{k}}$$
(22)

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The  $A_{2n}$  metrics  $g^{\omega}$  admit foliations by submanifolds which are Lagrangian for the symplectic forms  $\Omega_I, \Omega_J, \Omega_K$ . We call such a foliation a hyper-Lagrangian foliation.

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#### Proposition (Projectable hyper-Lagrangian foliation)

Given a hyper-Lagrangian foliation which pushes down to a Lagrangian foliation of M,  $\Theta$  can be taken to be at most quadratic in half of the coordinates  $y^i$ . When n = 1 such a foliation implies the heavenly equation linearises in an appropriate sense.

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### Many open questions:

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### Many open questions:

Example rather contrived without link to Joyce structures... Generally, which isomonodromy problems are set up correctly to have an analogous complex hyper-Kähler metric on the space X parametrising potentials?

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- Calculations can presumably be adapted to quadratic differentials with any fixed number of poles with fixed orders... Does some confluence phenomenon manifest on the level of the metric?

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